

Simulating dynamics and topological phases of cold fermionic gases

Lei Wang
ETH Zurich

Collaborators:

ETH Zürich

Ping Nang Ma
Ilya Zintchenko
Alexey Soluyanov
Matthias Troyer



Sebastiano Pilati

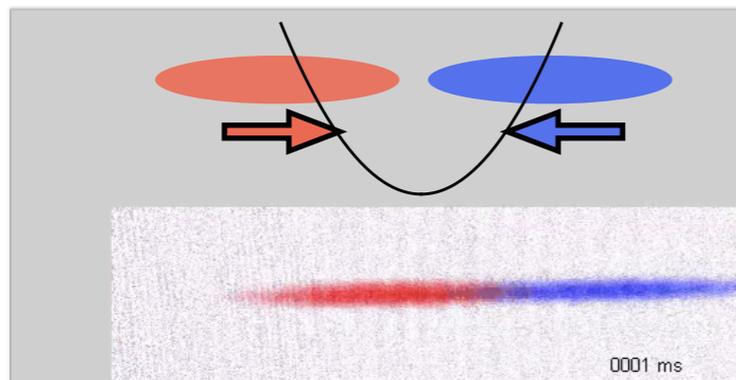


Xi Dai

Numerical simulations for **dynamics** of **fermionic** atoms in **high** dimensions

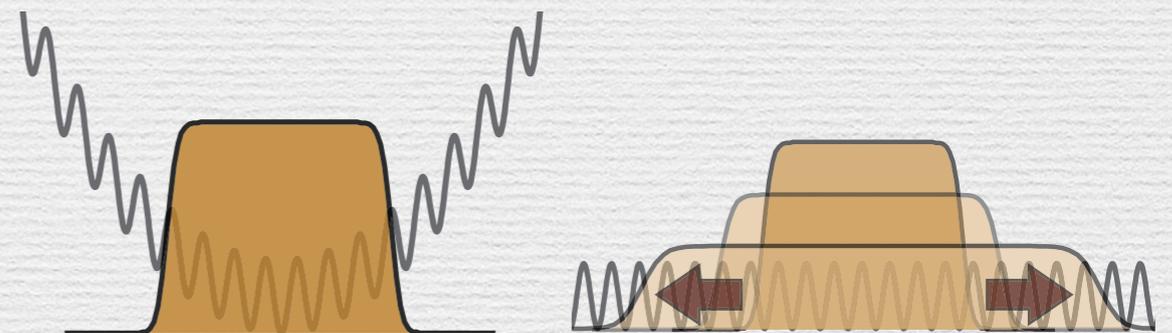
- ☛ It is difficult
- ☛ But, there were experiments...

Collision in a 3D trap



Sommer *et al*, 2011

Expansion in 2D optical lattice

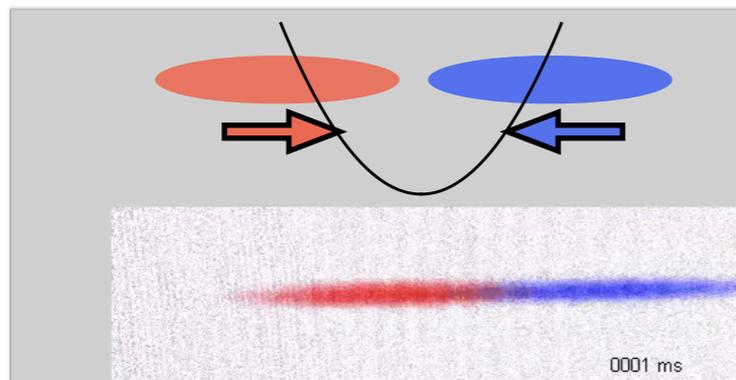


Schneider *et al*, 2012

Numerical simulations for **dynamics** of **fermionic** atoms in **high** dimensions

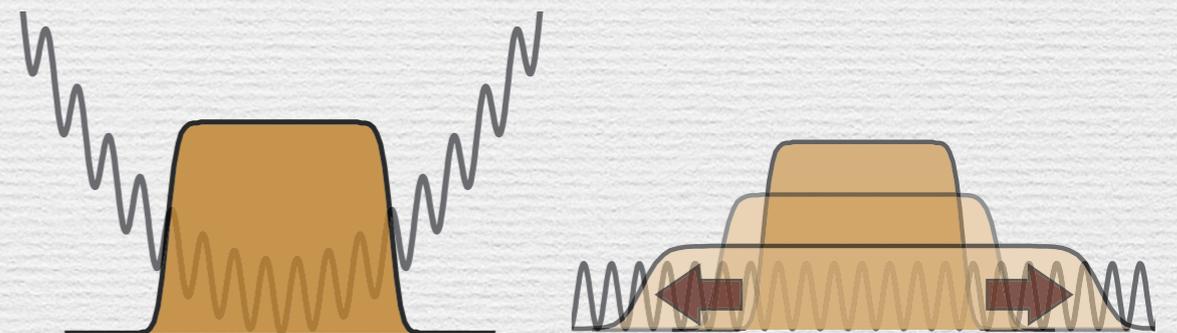
- It is difficult
- But, there were experiments...

Collision in a 3D trap



Sommer et al, 2011

Expansion in 2D optical lattice



Schneider et al, 2012

Density functional theory

Hohenberg and Kohn 1964

$$\Psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N)$$

$$\mathbb{R}^{3N} \mapsto \mathbb{C}$$



$$\rho(\mathbf{r})$$

$$\mathbb{R}^3 \mapsto \mathbb{R}$$

- Hohenberg-Kohn theorem: **All** properties of the system are **completely** determined by the ground state density.
- **Exact** ground state density and energy can be obtained by minimizing a **universal** density functional.

Density functional theory

Hohenberg and Kohn 1964

$$\Psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N)$$

$$\mathbb{R}^{3N} \mapsto \mathbb{C}$$

$$\rho(\mathbf{r})$$

$$\mathbb{R}^3 \mapsto \mathbb{R}$$

- Hohenberg-Kohn theorem: **All** properties of the system are **completely** determined by the ground state density.
- **Exact** ground state density and energy can be obtained by minimizing a **universal** density functional.
- In practice, obtain many-particle density from an auxiliary noninteracting system

Kohn and Sham 1965

$$\left(-\frac{\hbar^2 \nabla^2}{2m} + V_{\text{ext}} + V_{\text{H}}[\rho] + V_{\text{XC}}[\rho] \right) \psi_j = \varepsilon_j \psi_j$$

Time-dependent DFT

Runge and Gross, 1984

- **Time-dependent density** also plays a central role for non-equilibrium systems
- In practice, it is obtained from

$$i\frac{\partial}{\partial t}\psi_j(\mathbf{r}, t) = \left[-\frac{\hbar^2\nabla^2}{2m} + V_{\text{ext}}(\mathbf{r}, t) + V_{\text{H}}(\mathbf{r}, t) + V_{\text{xc}}[\rho(\mathbf{r}', t')](\mathbf{r}, t) \right] \psi_j(\mathbf{r}, t)$$

Time-dependent DFT

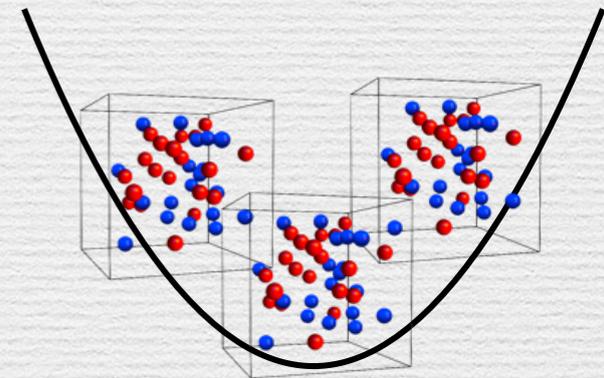
Runge and Gross, 1984

- **Time-dependent density** also plays a central role for non-equilibrium systems

- In practice, it is obtained from

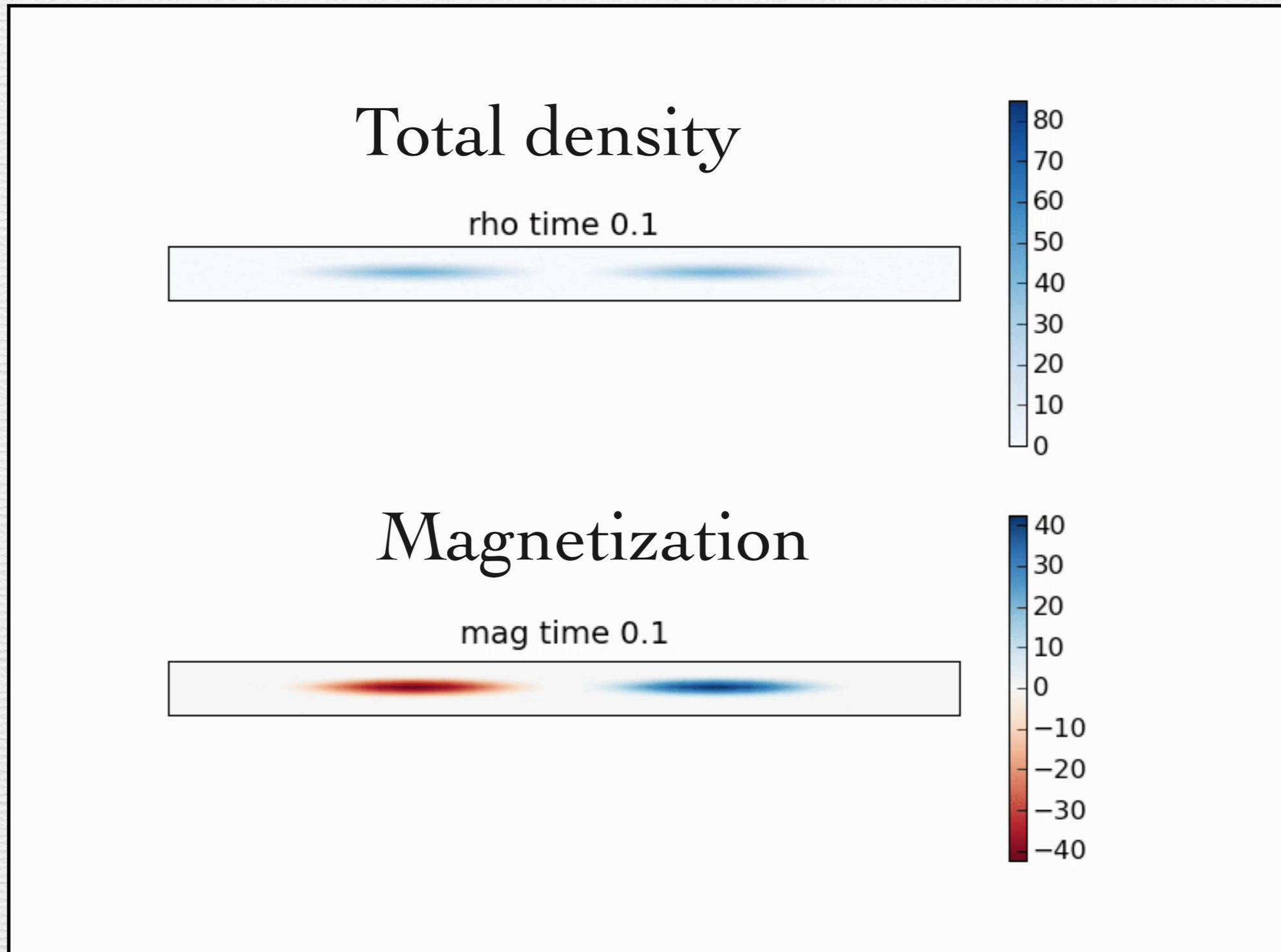
$$i\frac{\partial}{\partial t}\psi_j(\mathbf{r}, t) = \left[-\frac{\hbar^2\nabla^2}{2m} + V_{\text{ext}}(\mathbf{r}, t) + V_{\text{H}}(\mathbf{r}, t) + V_{\text{xc}}[\rho(\mathbf{r}', t')](\mathbf{r}, t) \right] \psi_j(\mathbf{r}, t)$$

- We use the **adiabatic** local-density approximation

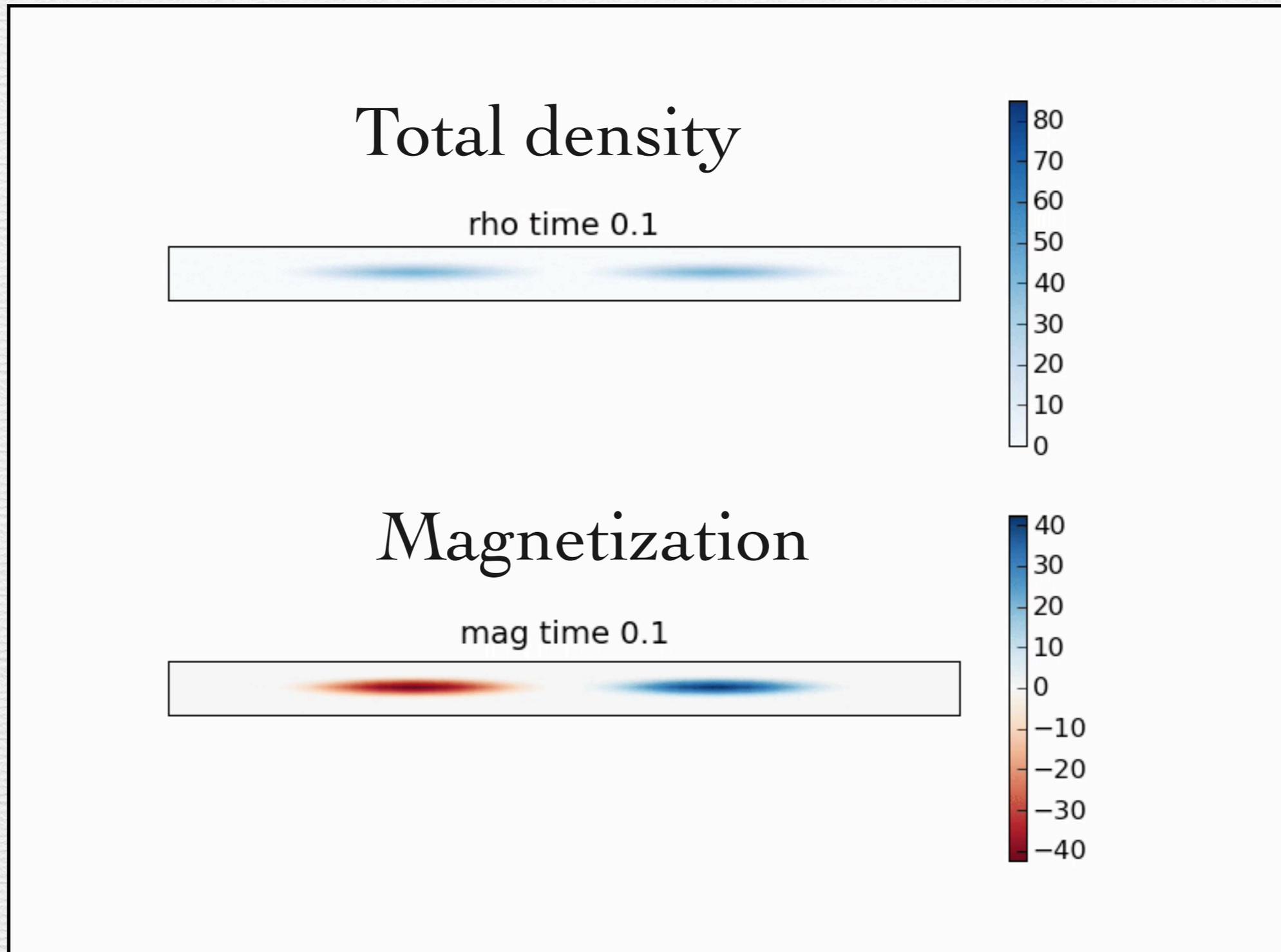


- **V_{xc}** from diffusion Monte-Carlo simulation of uniform atomic gases [Pilati *et al* 2010](#), [Ping Nang Ma *et al* 2012](#)

Simulation of cloud collisions



Simulation of cloud collisions



Deep optical lattices: Time-dependent Gutzwiller method

Schiro and Fabrizio, 2010

$$H = - \sum_{i,j,\sigma} t_{ij} \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_{i,\sigma} v_i^{\text{ext}} \hat{n}_{i\sigma} + U \sum_i \hat{n}_{i\uparrow} \hat{n}_{i\downarrow}$$

Deep optical lattices: Time-dependent Gutzwiller method

Schiro and Fabrizio, 2010

$$H = - \sum_{i,j,\sigma} t_{ij} \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_{i,\sigma} v_i^{\text{ext}} \hat{n}_{i\sigma} + U \sum_i \hat{n}_{i\uparrow} \hat{n}_{i\downarrow}$$

$$|\Psi(t)\rangle = \prod_i \hat{\mathcal{P}}_i |\text{Slater}\rangle$$

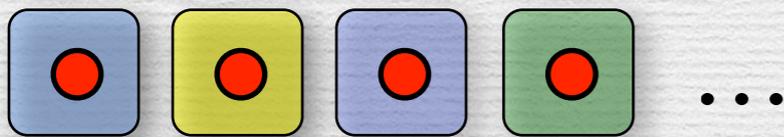
Deep optical lattices: Time-dependent Gutzwiller method

Schiro and Fabrizio, 2010

$$H = - \sum_{i,j,\sigma} t_{ij} \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_{i,\sigma} v_i^{\text{ext}} \hat{n}_{i\sigma} + U \sum_i \hat{n}_{i\uparrow} \hat{n}_{i\downarrow}$$

$$|\Psi(t)\rangle = \prod_i \hat{\mathcal{P}}_i |\text{Slater}\rangle$$

$$H_B = \sum_i \hat{b}_i^\dagger \mathcal{H}_{iB} \hat{b}_i$$



$$\hat{b}_i = \begin{pmatrix} e \\ \uparrow \\ \downarrow \\ d \end{pmatrix}_i$$

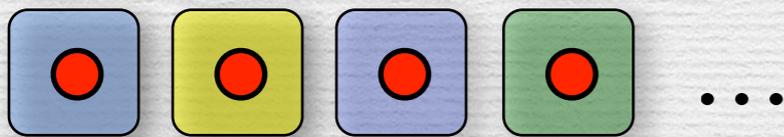
Deep optical lattices: Time-dependent Gutzwiller method

Schiro and Fabrizio, 2010

$$H = - \sum_{i,j,\sigma} t_{ij} \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_{i,\sigma} v_i^{\text{ext}} \hat{n}_{i\sigma} + U \sum_i \hat{n}_{i\uparrow} \hat{n}_{i\downarrow}$$

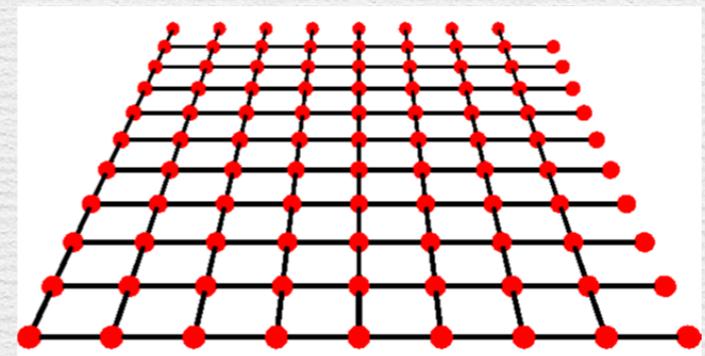
$$|\Psi(t)\rangle = \prod_i \hat{\mathcal{P}}_i |\text{Slater}\rangle$$

$$H_B = \sum_i \hat{b}_i^\dagger \mathcal{H}_{iB} \hat{b}_i$$



$$\hat{b}_i = \begin{pmatrix} e \\ \uparrow \\ \downarrow \\ d \end{pmatrix}_i$$

$$H_F = - \sum_{i,j,\sigma} t_{ij}^* \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_{i,\sigma} v_i^{\text{ext}} \hat{n}_{i\sigma}$$



$$t_{ij}^* = z_i^* t_{ij} z_j$$

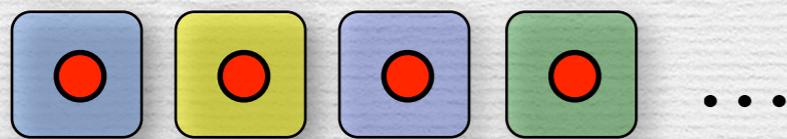
Deep optical lattices: Time-dependent Gutzwiller method

Schiro and Fabrizio, 2010

$$H = - \sum_{i,j,\sigma} t_{ij} \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_{i,\sigma} v_i^{\text{ext}} \hat{n}_{i\sigma} + U \sum_i \hat{n}_{i\uparrow} \hat{n}_{i\downarrow}$$

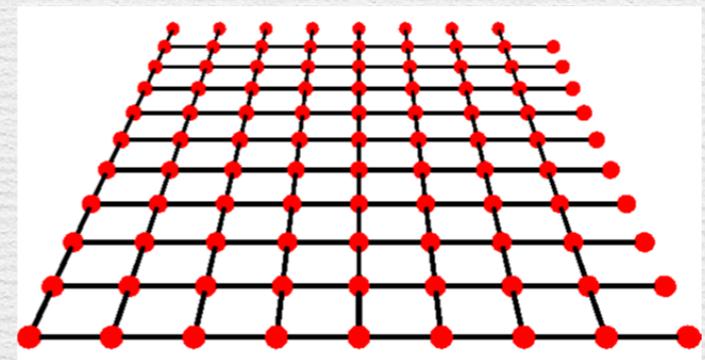
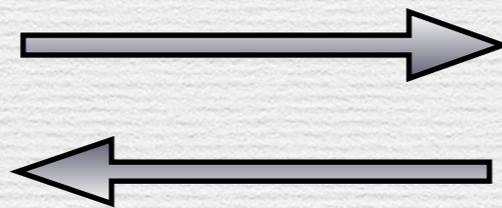
$$|\Psi(t)\rangle = \prod_i \hat{\mathcal{P}}_i |\text{Slater}\rangle$$

$$H_B = \sum_i \hat{b}_i^\dagger \mathcal{H}_{iB} \hat{b}_i$$



$$\hat{b}_i = \begin{pmatrix} e \\ \uparrow \\ \downarrow \\ d \end{pmatrix}_i$$

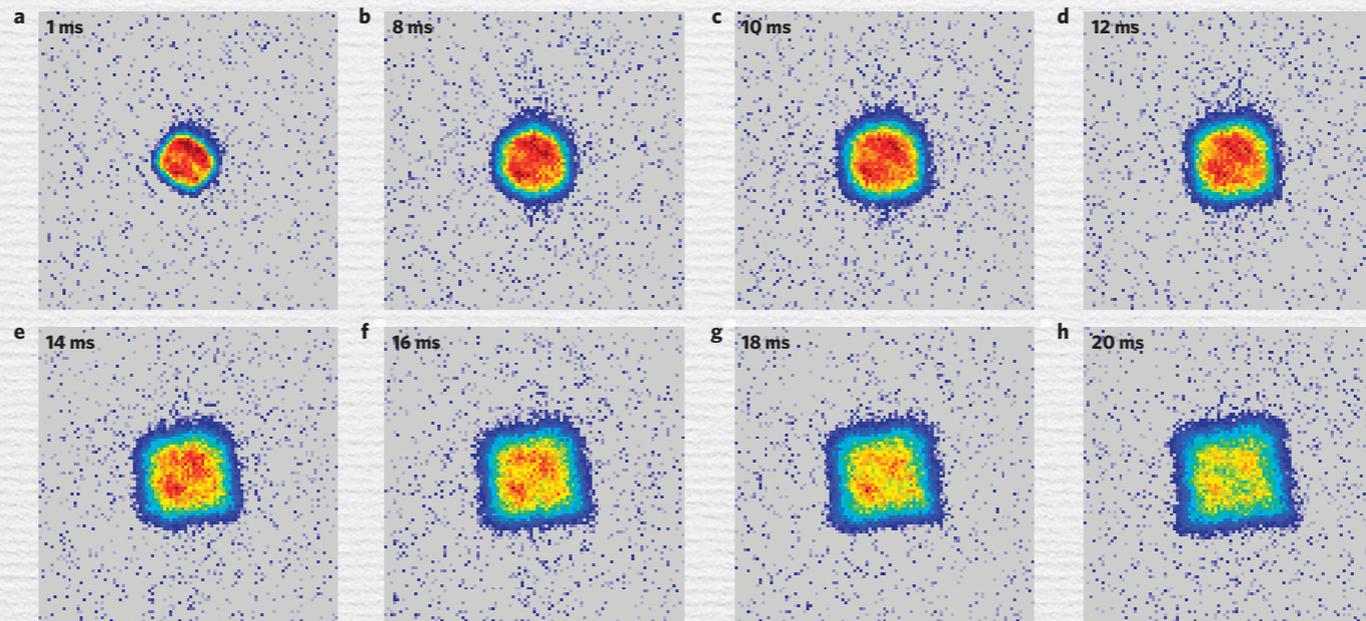
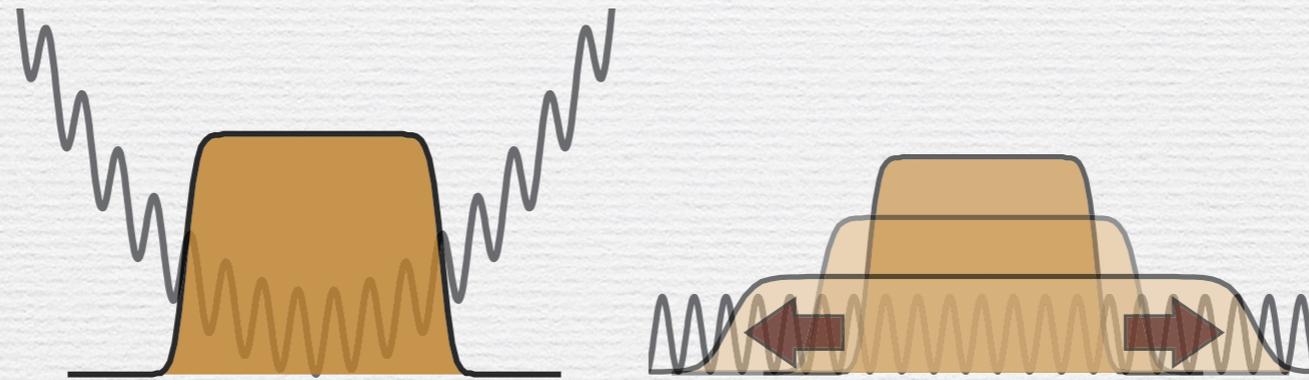
$$H_F = - \sum_{i,j,\sigma} t_{ij}^* \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_{i,\sigma} v_i^{\text{ext}} \hat{n}_{i\sigma}$$



$$t_{ij}^* = z_i^* t_{ij} z_j$$

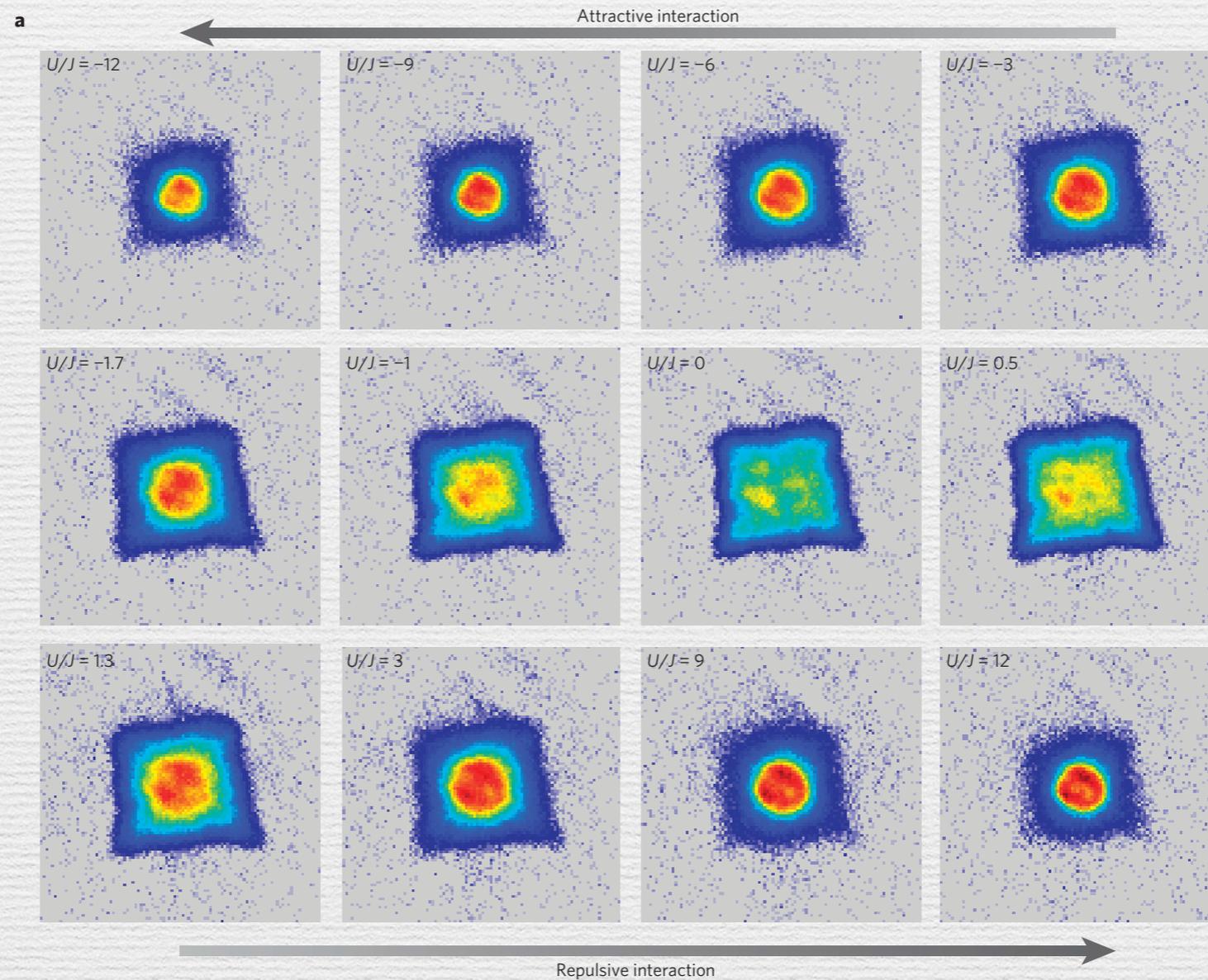
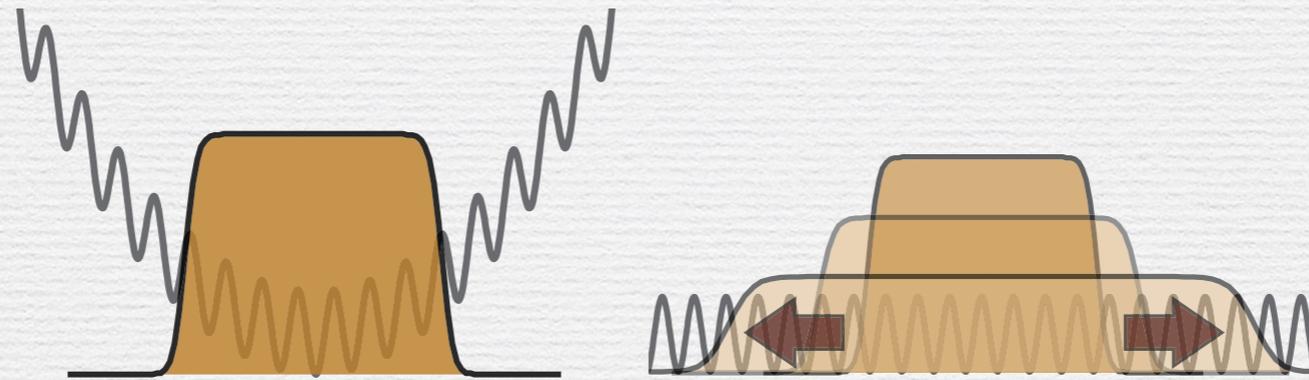
Simulation of cloud expansion

Schneider *et al*, 2012



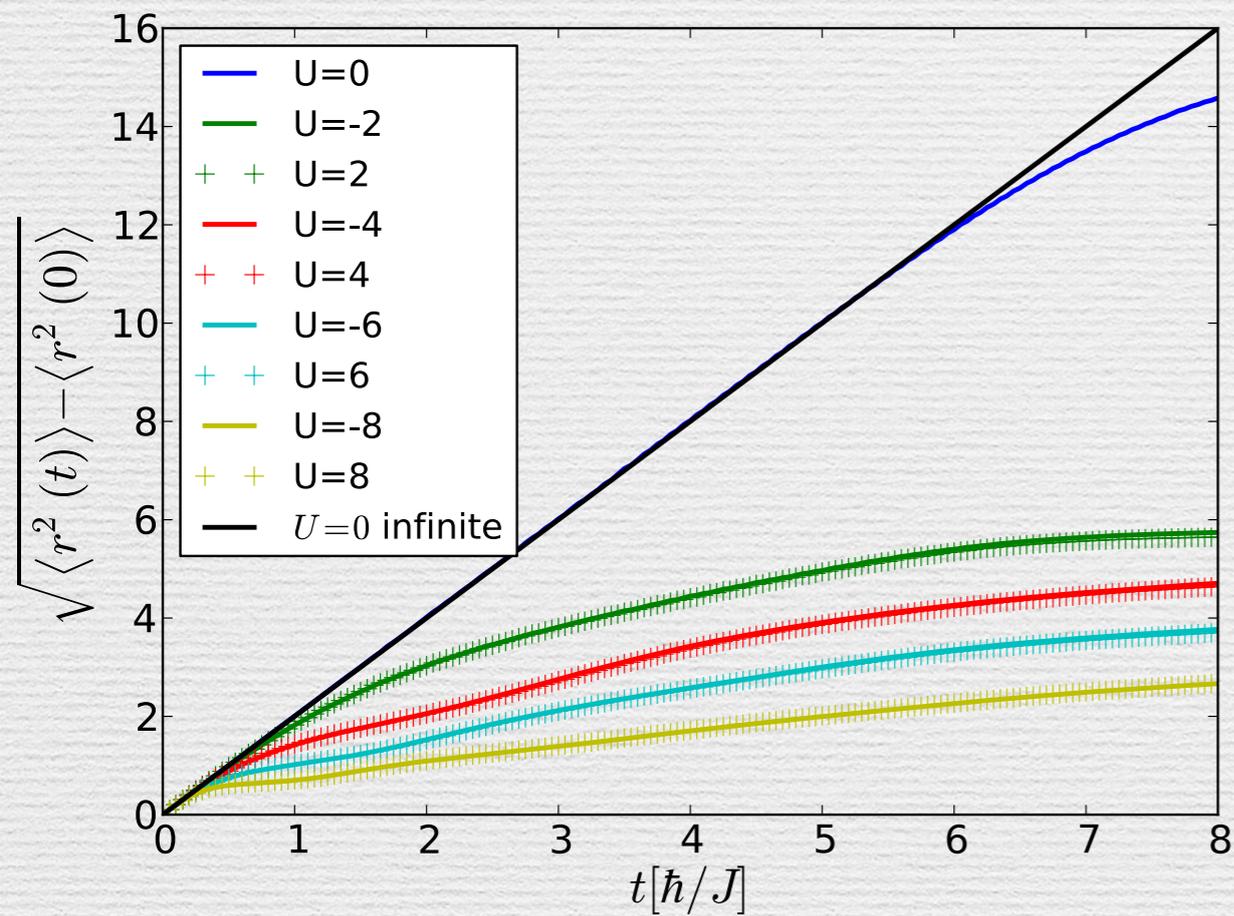
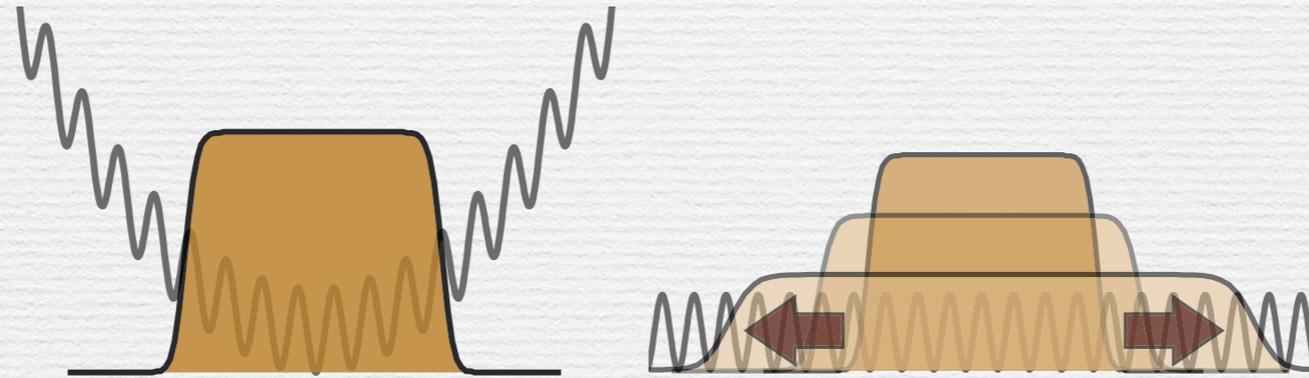
Simulation of cloud expansion

Schneider *et al*, 2012



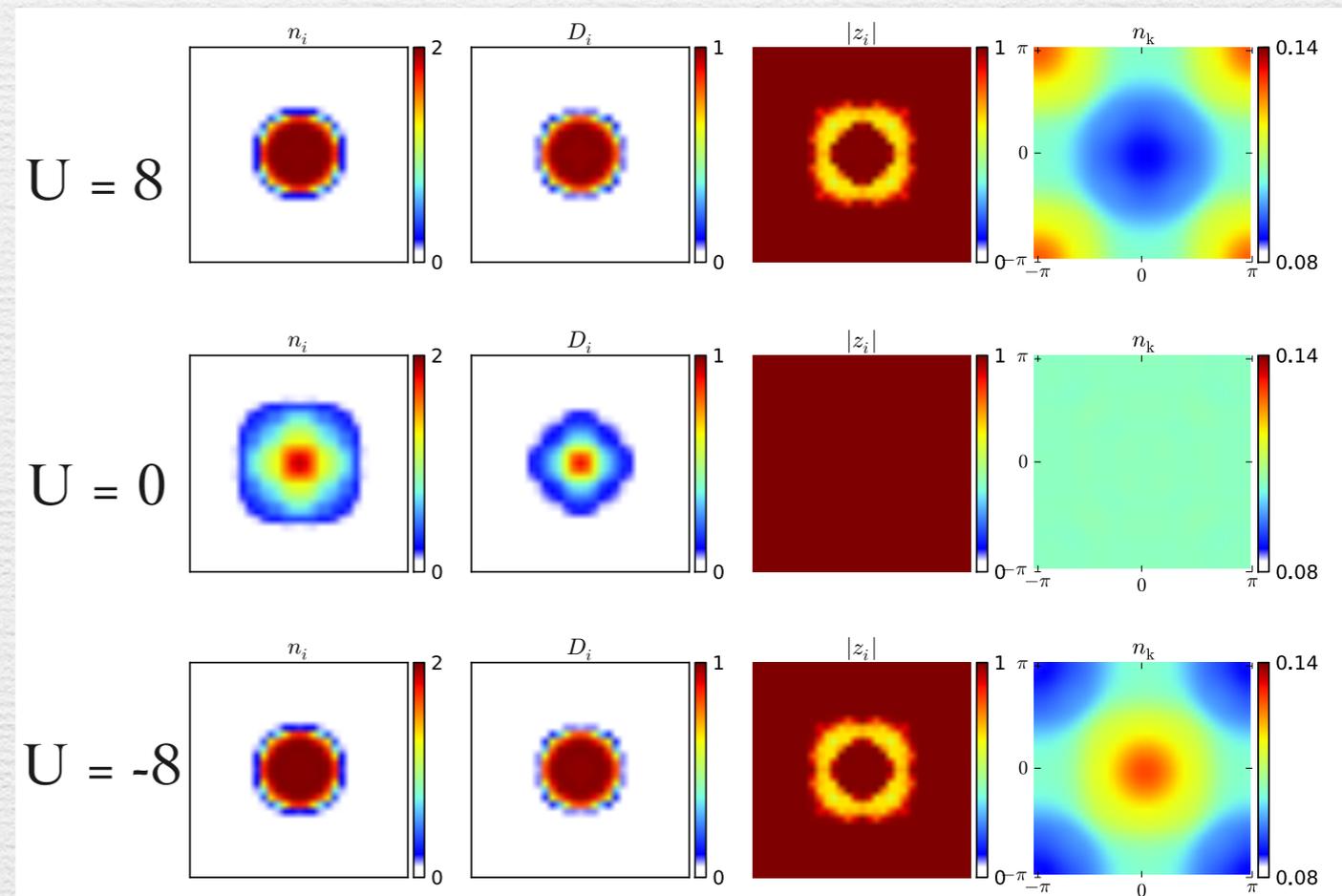
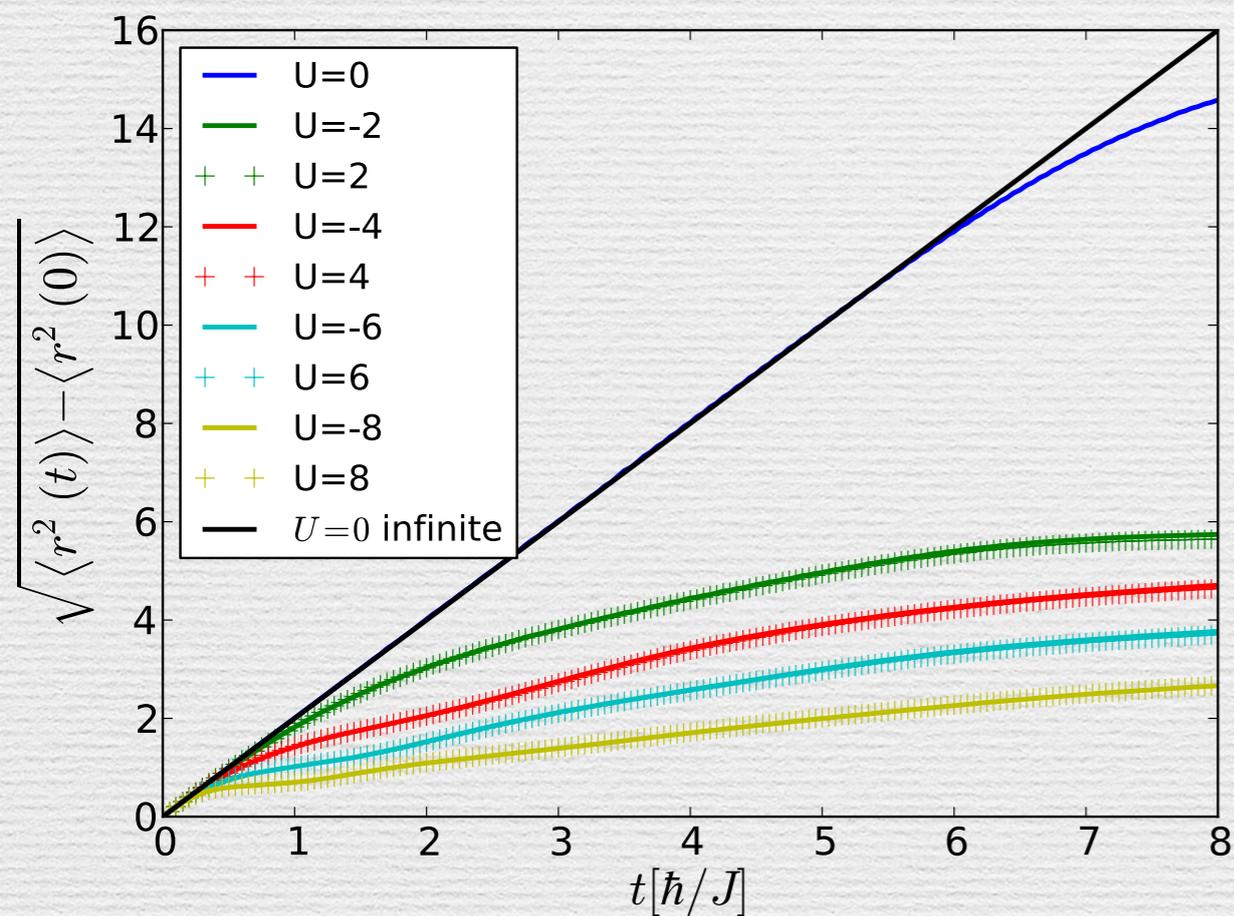
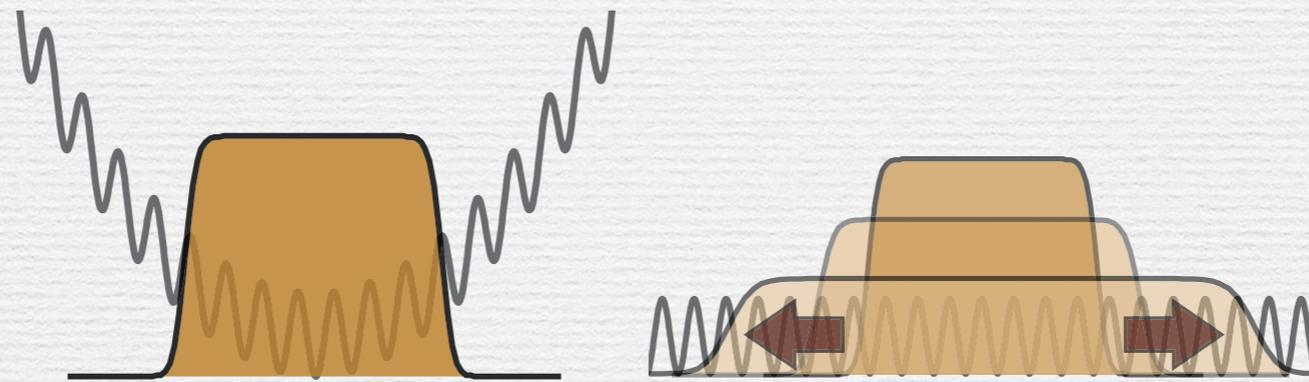
Simulation of cloud expansion

Schneider *et al*, 2012

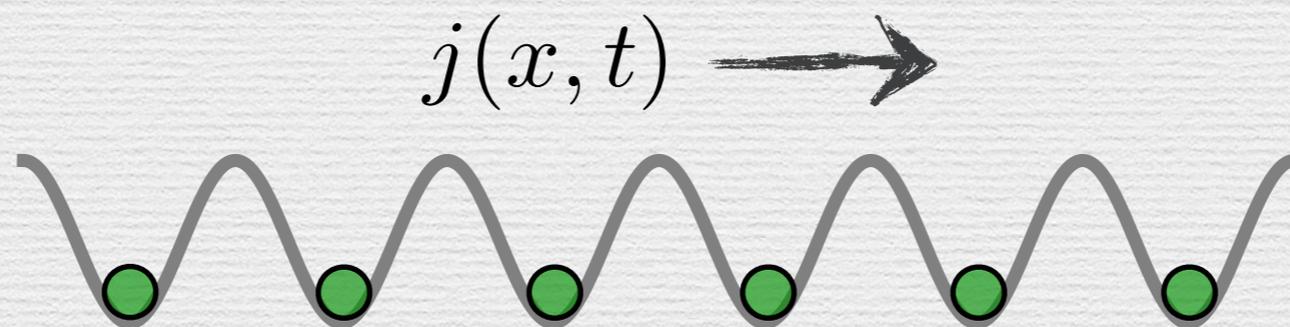


Simulation of cloud expansion

Schneider *et al*, 2012



Topological charge pumping of cold atoms



Pumps



A **pump** is a device that moves fluids, or sometimes slurries, by mechanical action.

Pumps



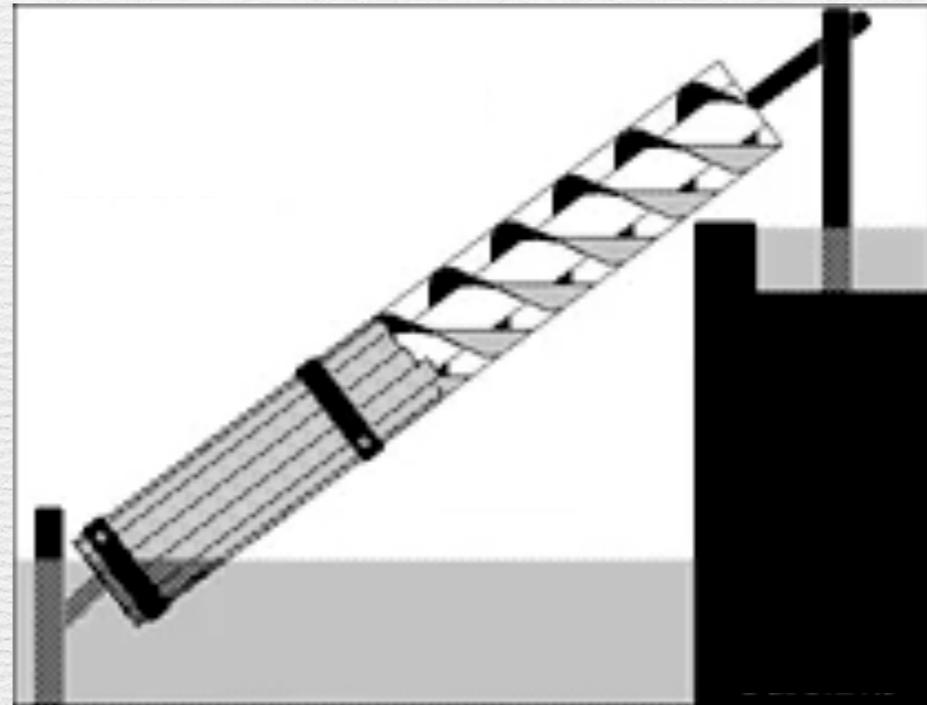
A **pump** is a device that moves fluids, or sometimes slurries, by mechanical action.



Pumps



A **pump** is a device that moves fluids, or sometimes slurries, by mechanical action.

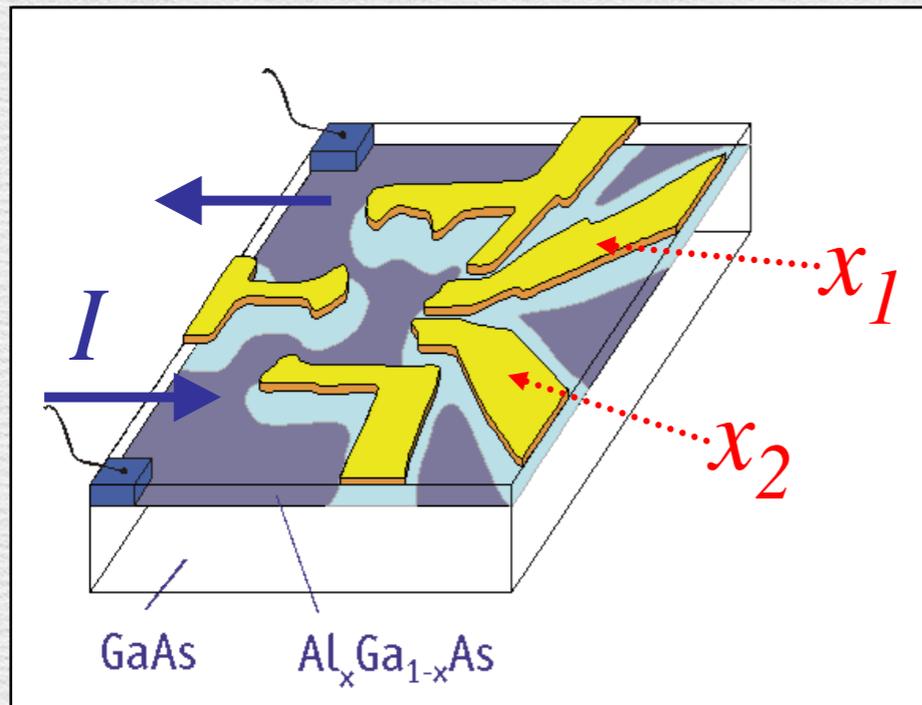


Archimedes' screw ~250 BC

Pumps

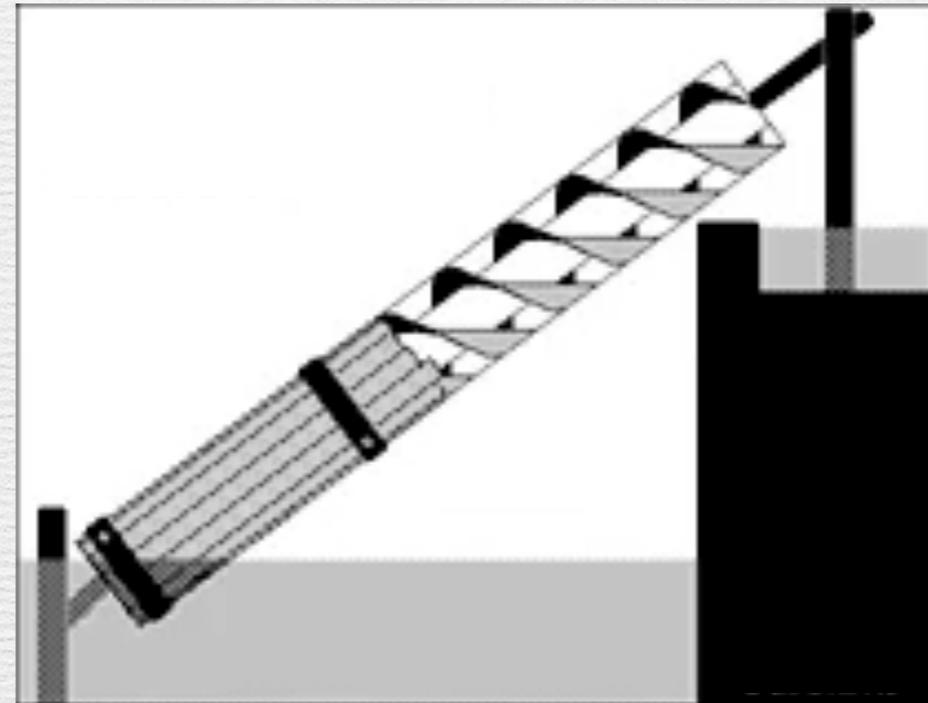


A **pump** is a device that moves fluids, or sometimes slurries, by mechanical action.



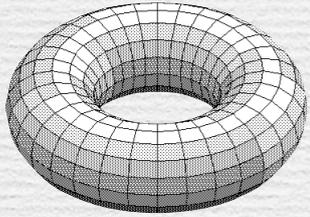
Switkes *et al* 1999

Buttiker, Brouwer, Zhou, Spivak, Altshuler ...



Archimedes' screw ~250 BC

Topological pump

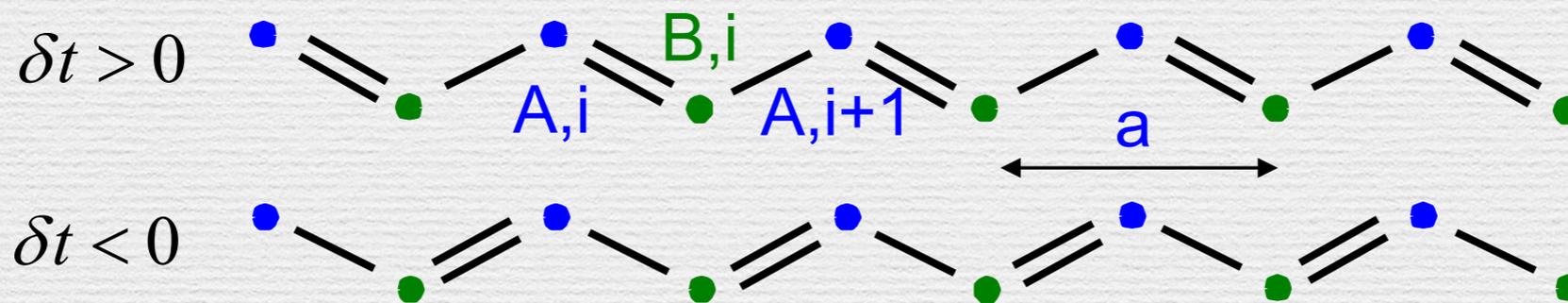


A **topological pump** transfers quantized charge in each pumping cycle.

Thouless, Niu, 1980s

$$H = \sum_i (t + \delta t) c_{A_i}^\dagger c_{B_i} + (t - \delta t) c_{A_{i+1}}^\dagger c_{B_i} + H.c.$$

Su, Schrieffer, Heeger, 1979

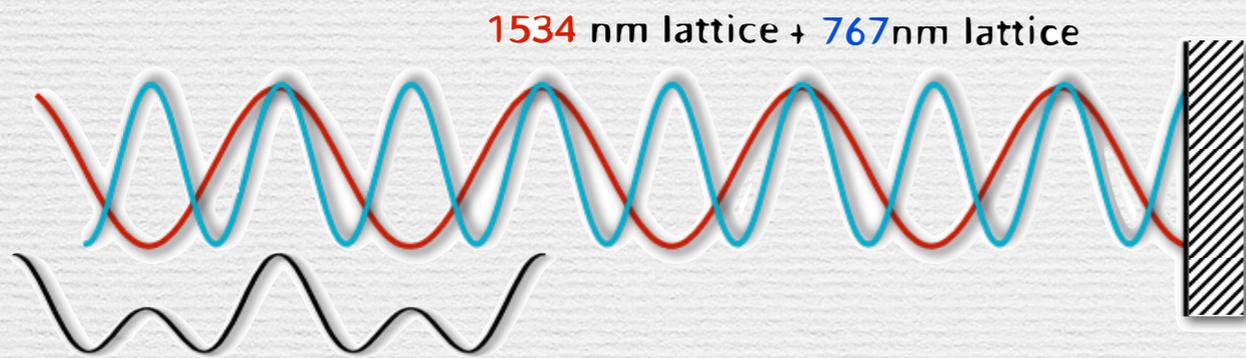


- Current flows in an insulating state
- No dissipation!
- Dynamical analog of quantum Hall effect

Experimental progresses

Optical Superlattice

Fölling *et al*, Atala *et al*

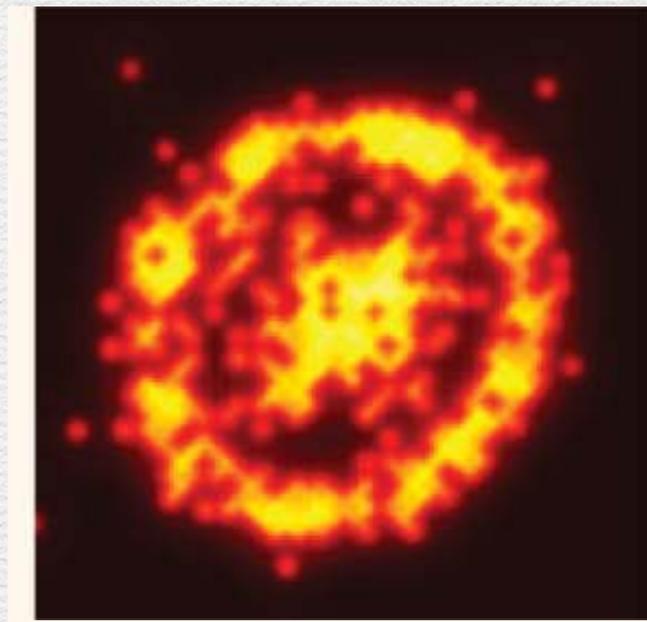


$$V_{\text{OL}}(x) = V_1 \cos^2\left(\frac{2\pi x}{d}\right) + V_2 \cos^2\left(\frac{\pi x}{d} - \varphi\right)$$

Full (independent) dynamical control over V_1 , V_2 and φ

in-situ imaging

Gemelke, *et al*, Sherson *et al*, Bakr *et al*

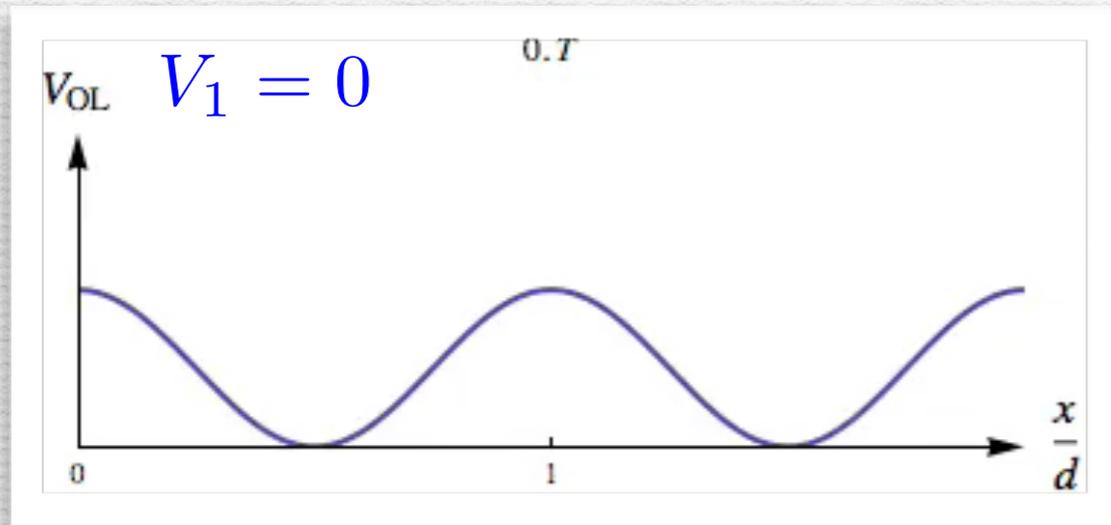


1D pumping lattices

$$V_{\text{OL}}(x, t) = V_1 \cos^2 \left(\frac{2\pi x}{d} \right) + V_2 \cos^2 \left(\frac{\pi x}{d} - \frac{\pi t}{T} \right)$$

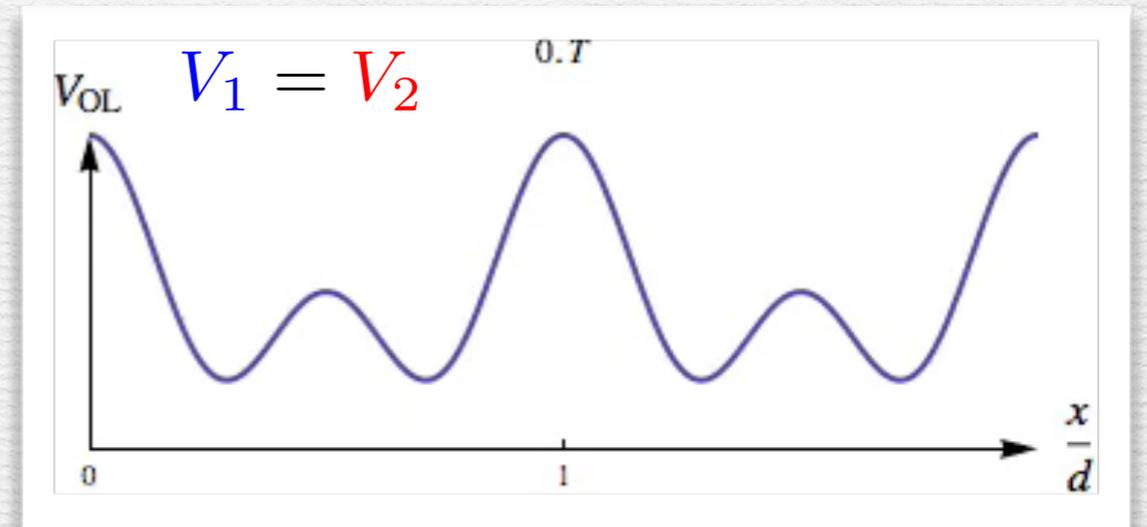
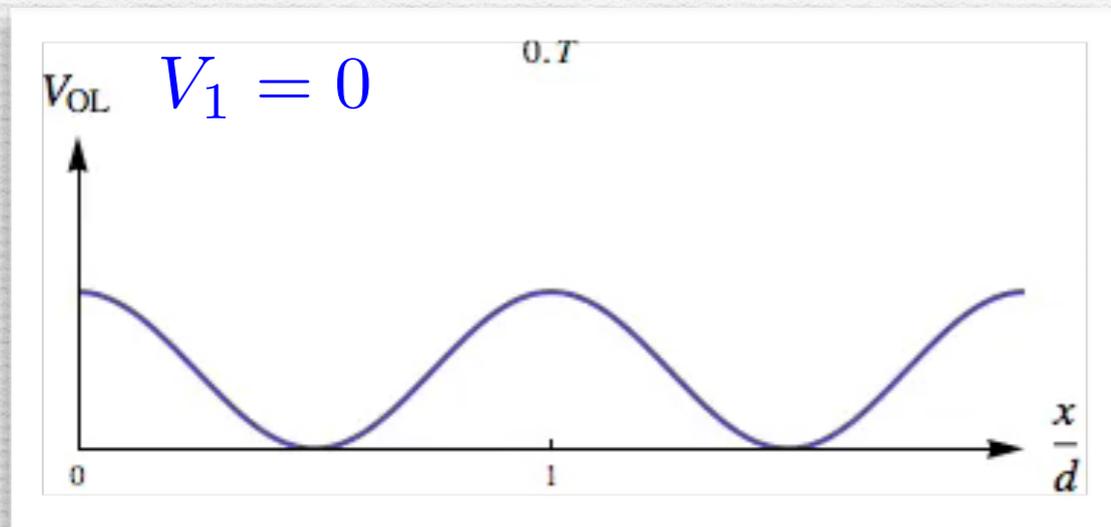
1D pumping lattices

$$V_{\text{OL}}(x, t) = V_1 \cos^2 \left(\frac{2\pi x}{d} \right) + V_2 \cos^2 \left(\frac{\pi x}{d} - \frac{\pi t}{T} \right)$$



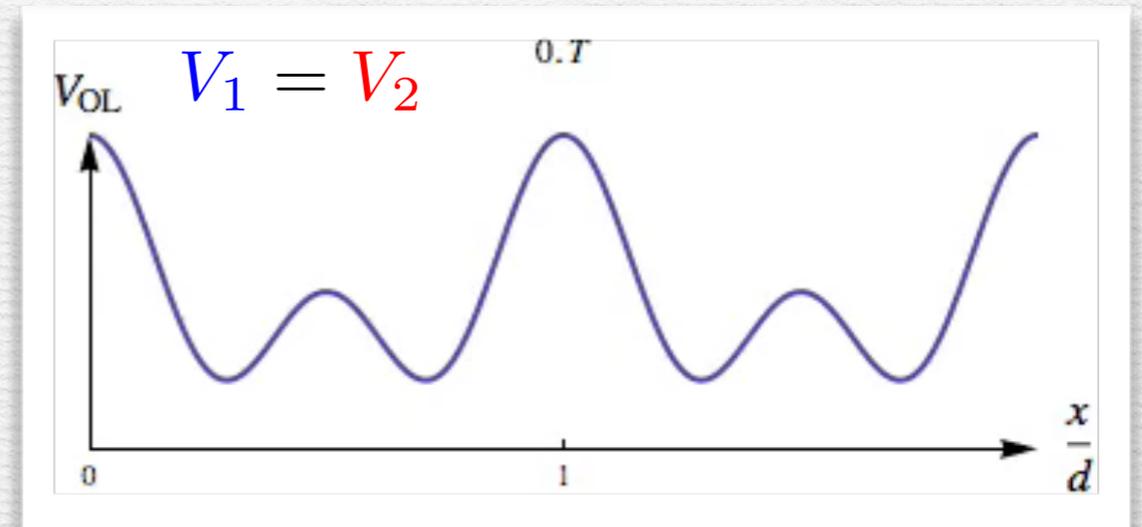
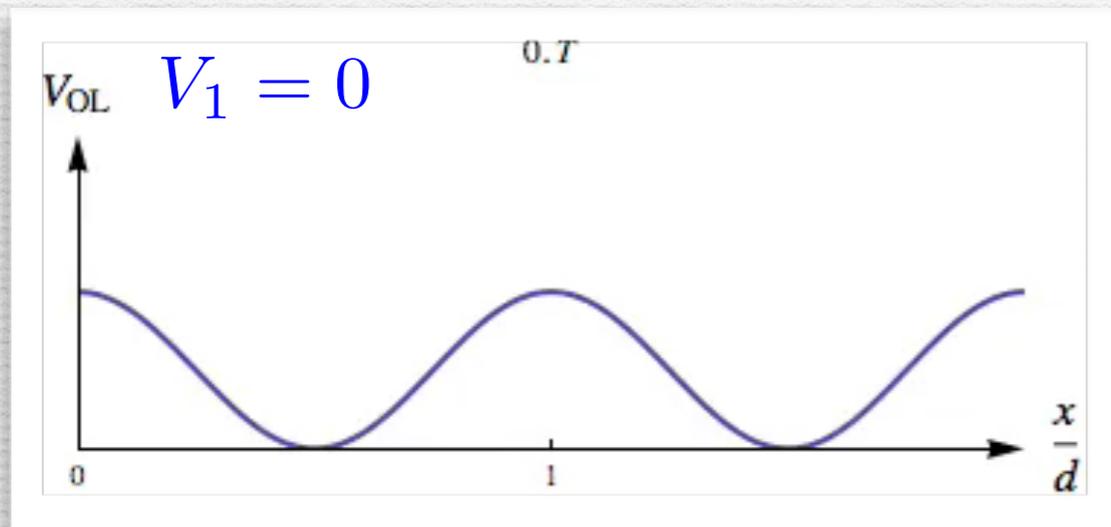
1D pumping lattices

$$V_{\text{OL}}(x, t) = V_1 \cos^2\left(\frac{2\pi x}{d}\right) + V_2 \cos^2\left(\frac{\pi x}{d} - \frac{\pi t}{T}\right)$$



1D pumping lattices

$$V_{\text{OL}}(x, t) = V_1 \cos^2\left(\frac{2\pi x}{d}\right) + V_2 \cos^2\left(\frac{\pi x}{d} - \frac{\pi t}{T}\right)$$



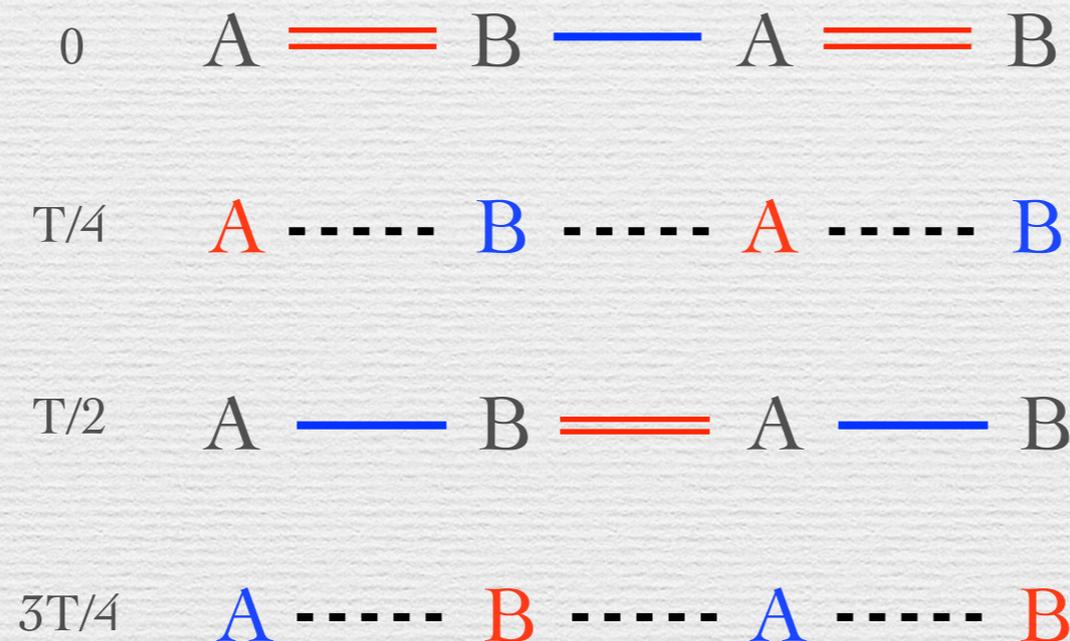
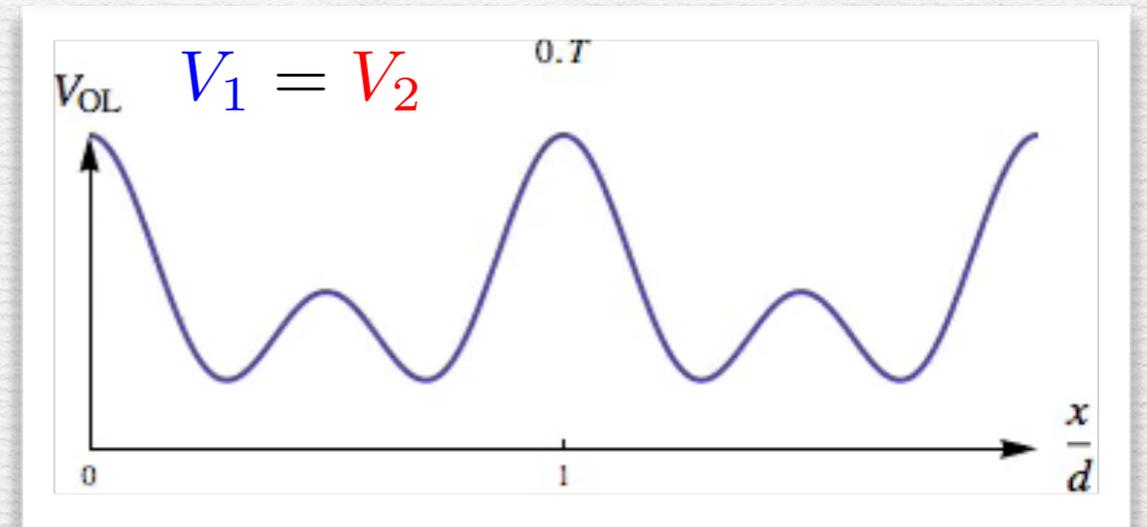
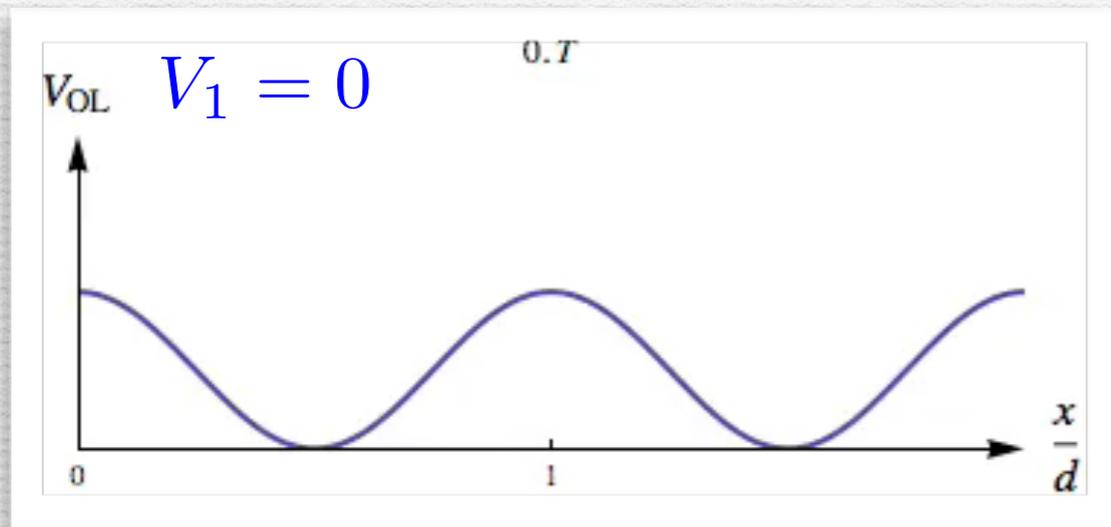
0 A == B ——— A == B

Su, Schrieffer, Heeger, 1979

T/2 A ——— B == A ——— B

1D pumping lattices

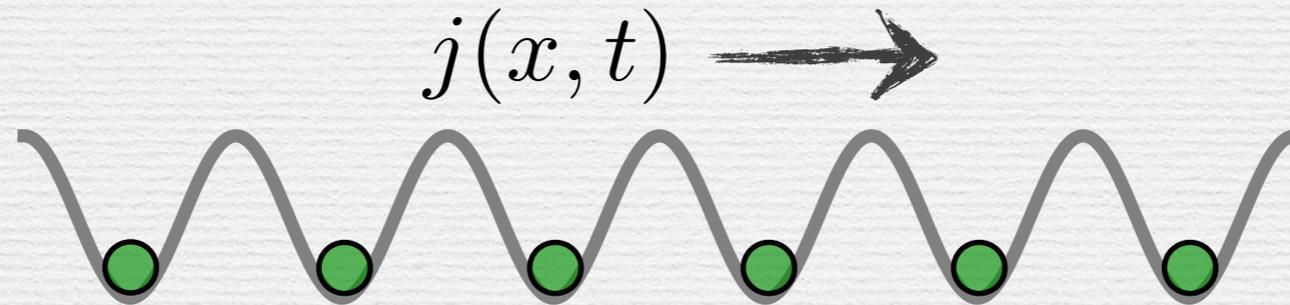
$$V_{\text{OL}}(x, t) = V_1 \cos^2\left(\frac{2\pi x}{d}\right) + V_2 \cos^2\left(\frac{\pi x}{d} - \frac{\pi t}{T}\right)$$



Su, Schrieffer, Heeger, 1979

Rice, Mele, 1982

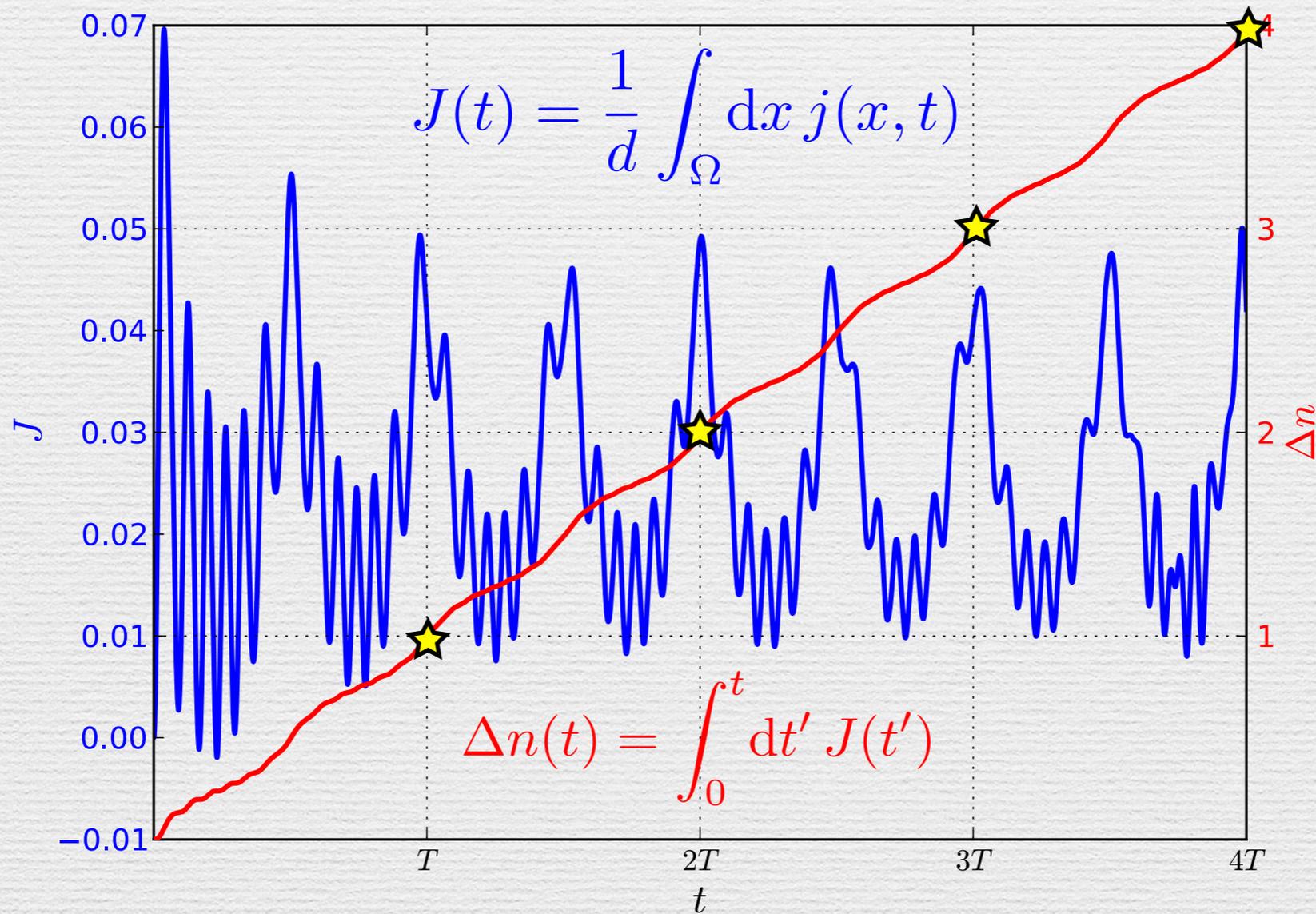
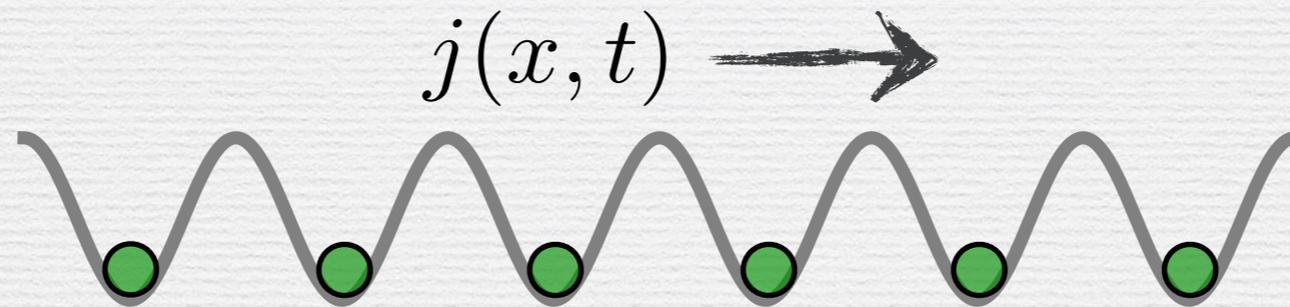
Quantum dynamics



$$H(x, t) = -\frac{\hbar^2}{2m} \nabla^2 + V_{\text{OL}}(x, t)$$

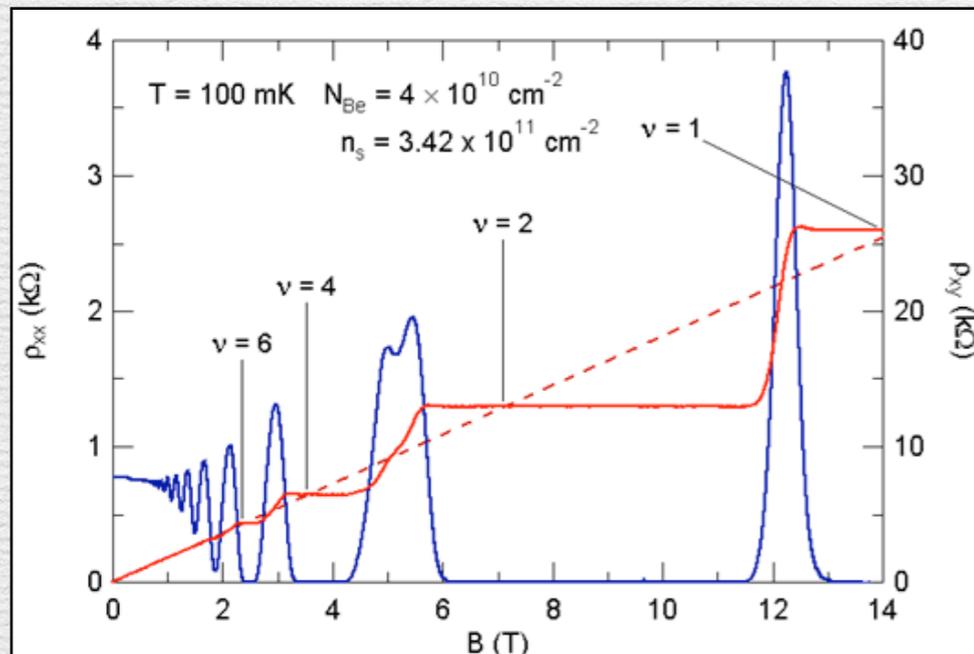
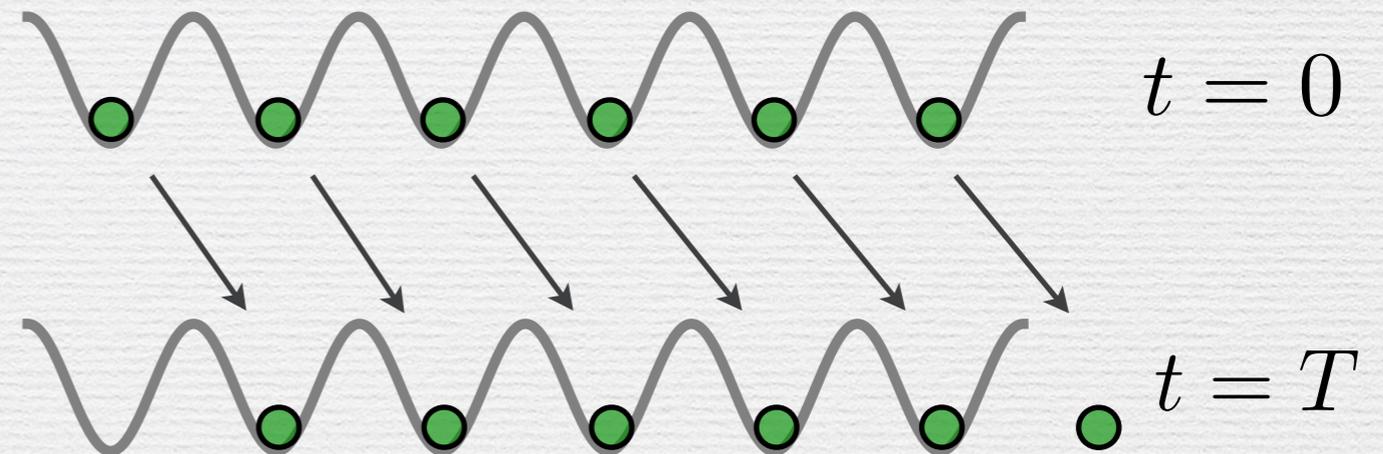
$$i \frac{\partial}{\partial t} |\Psi\rangle = H(x, t) |\Psi\rangle$$

Quantum dynamics



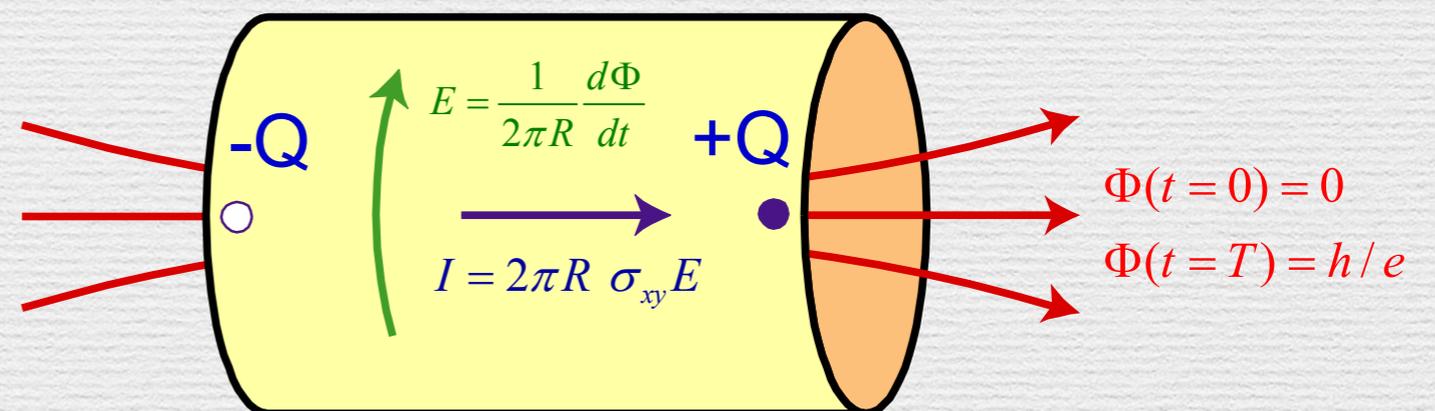
1D pump and 2D QHE

$$H(k_x, t) = H(k_x, t + T)$$



Von Klitzing *et al*, 1980

Adiabatically thread a quantum of magnetic flux through cylinder.



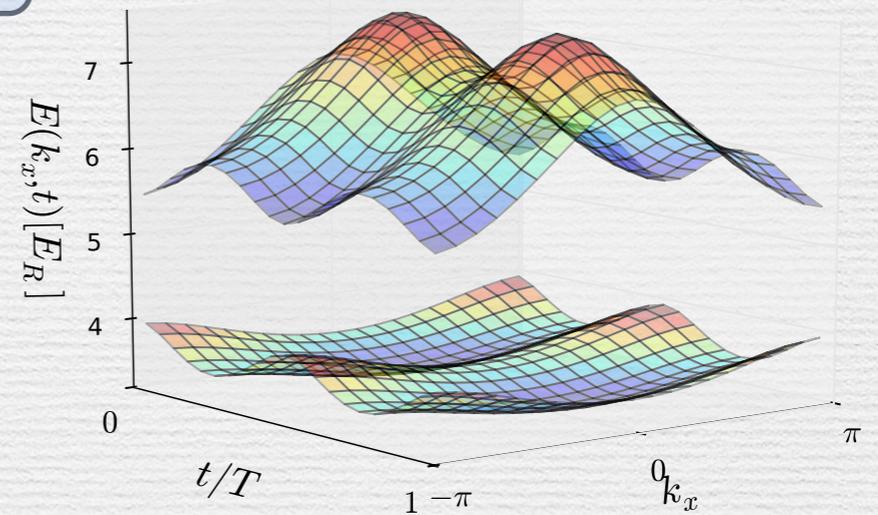
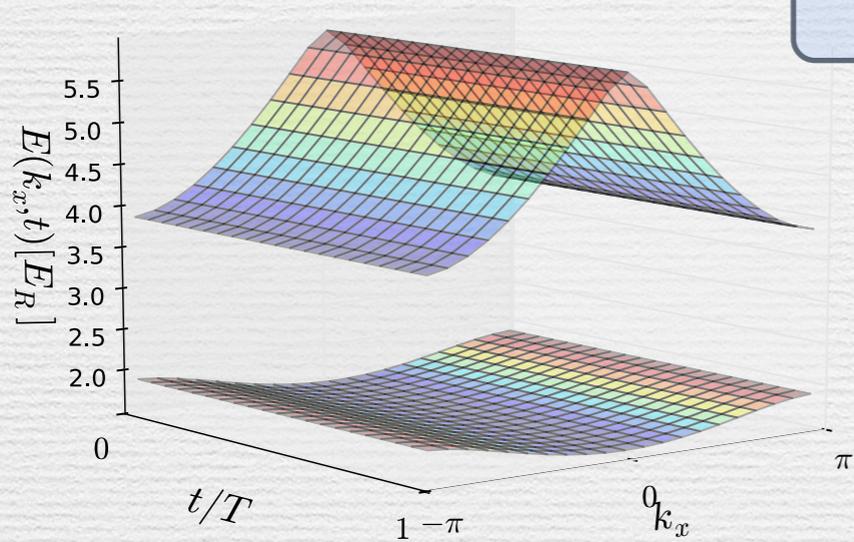
Laughlin, 1981

Gap & Chern number

$$V_1 = 0 \quad V_2 = 4E_R$$

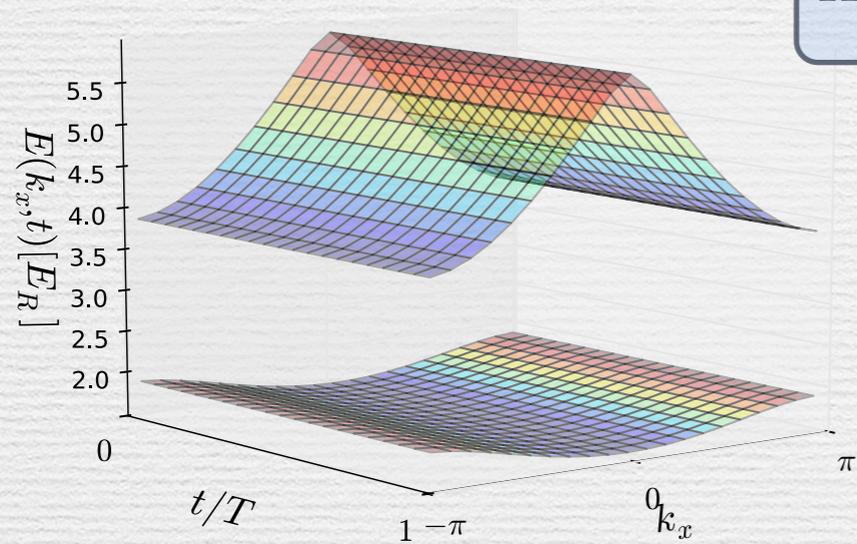
$$H(x, t) = -\frac{\hbar^2}{2m} \nabla^2 + V_{\text{OL}}(x, t)$$

$$V_1 = 4E_R \quad V_2 = 4E_R$$



Gap & Chern number

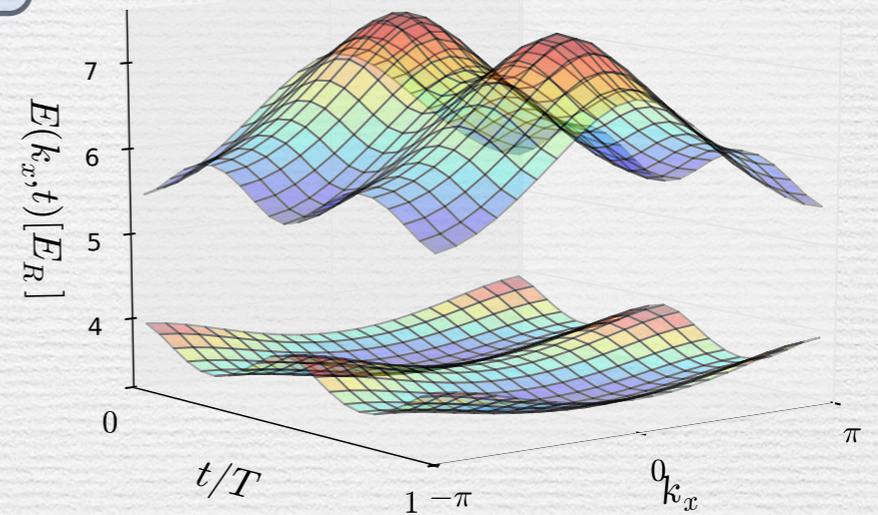
$$V_1 = 0 \quad V_2 = 4E_R$$



$$H(x, t) = -\frac{\hbar^2}{2m} \nabla^2 + V_{\text{OL}}(x, t)$$

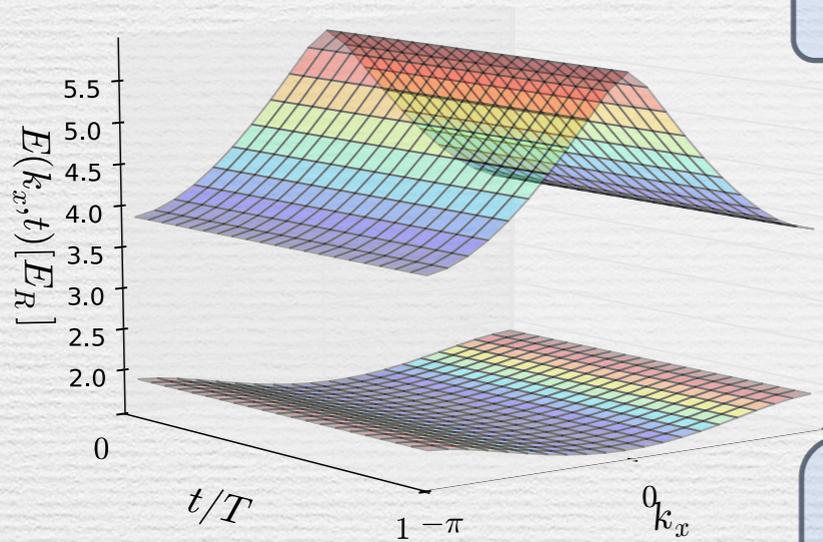


$$V_1 = 4E_R \quad V_2 = 4E_R$$



Gap & Chern number

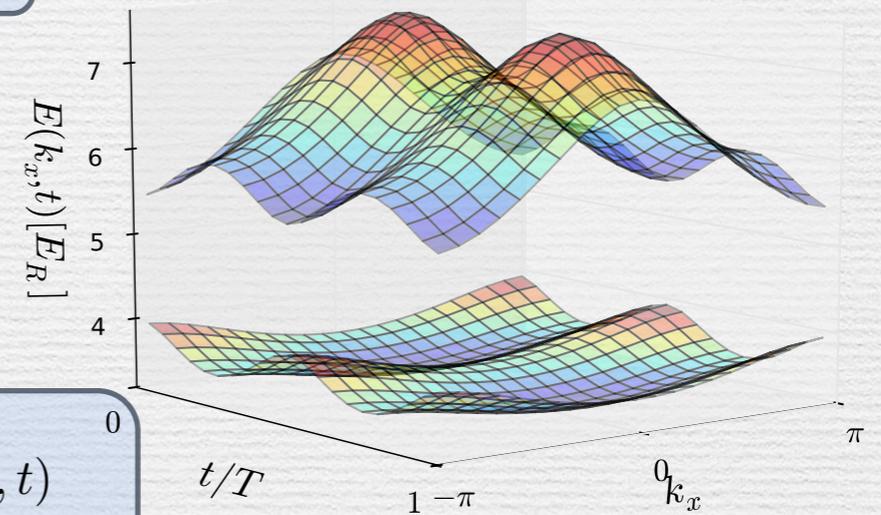
$$V_1 = 0 \quad V_2 = 4E_R$$



$$H(x, t) = -\frac{\hbar^2}{2m} \nabla^2 + V_{\text{OL}}(x, t)$$



$$V_1 = 4E_R \quad V_2 = 4E_R$$



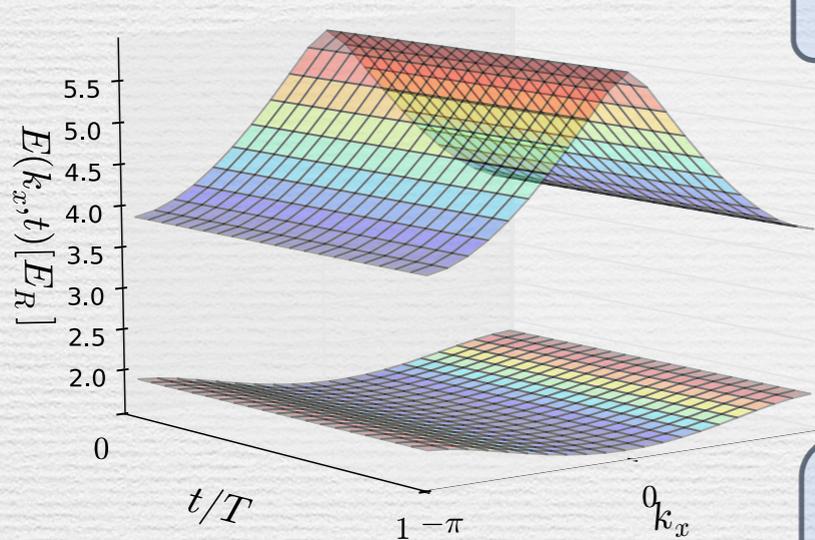
$$\Delta n = \frac{1}{2\pi} \int_0^T dt \int_0^{2\pi} dk_x \mathcal{F}(k_x, t)$$

$$\mathcal{F}(k_x, t) = \partial_t A_{k_x} - \partial_{k_x} A_t$$

$$A_{t(k_x)} = -i \langle \psi_{k_x}(t) | \partial_{t(k_x)} | \psi_{k_x}(t) \rangle$$

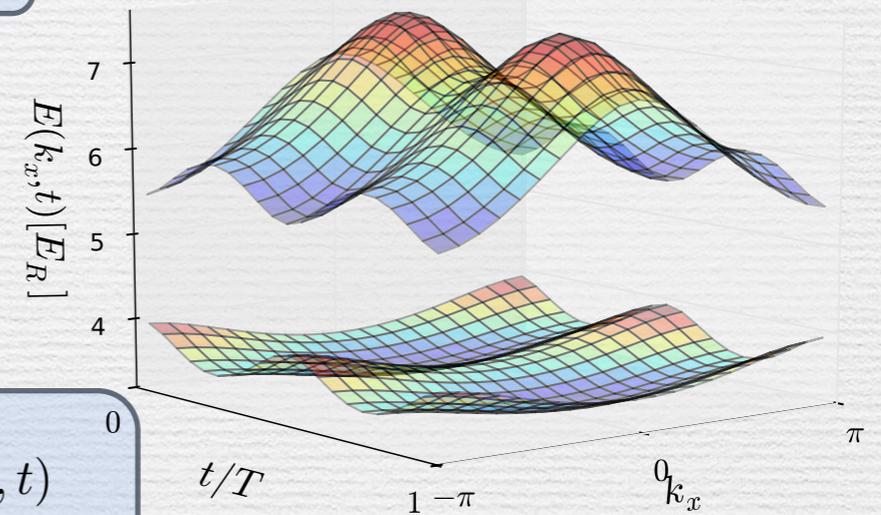
Gap & Chern number

$$V_1 = 0 \quad V_2 = 4E_R$$



$$H(x, t) = -\frac{\hbar^2}{2m} \nabla^2 + V_{OL}(x, t)$$

$$V_1 = 4E_R \quad V_2 = 4E_R$$

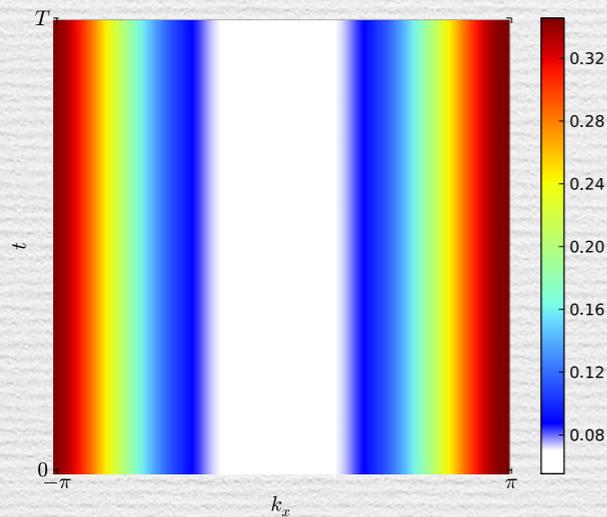


$$\Delta n = \frac{1}{2\pi} \int_0^T dt \int_0^{2\pi} dk_x \mathcal{F}(k_x, t)$$

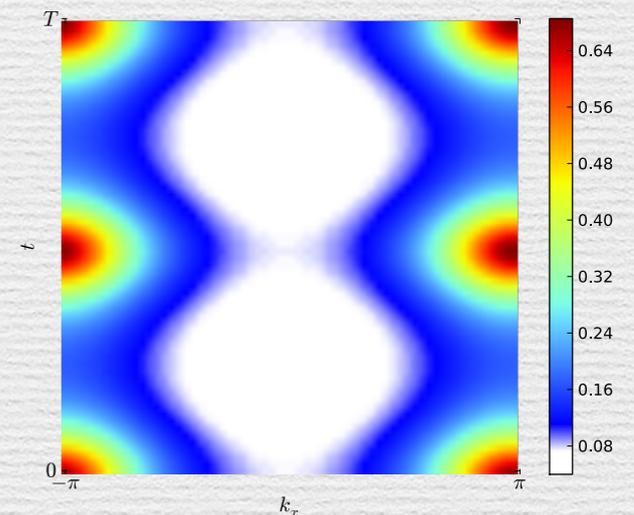
$$\mathcal{F}(k_x, t) = \partial_t A_{k_x} - \partial_{k_x} A_t$$

$$A_t(k_x) = -i \langle \psi_{k_x}(t) | \partial_{t(k_x)} | \psi_{k_x}(t) \rangle$$

$$\mathcal{F}(k_x, t)$$

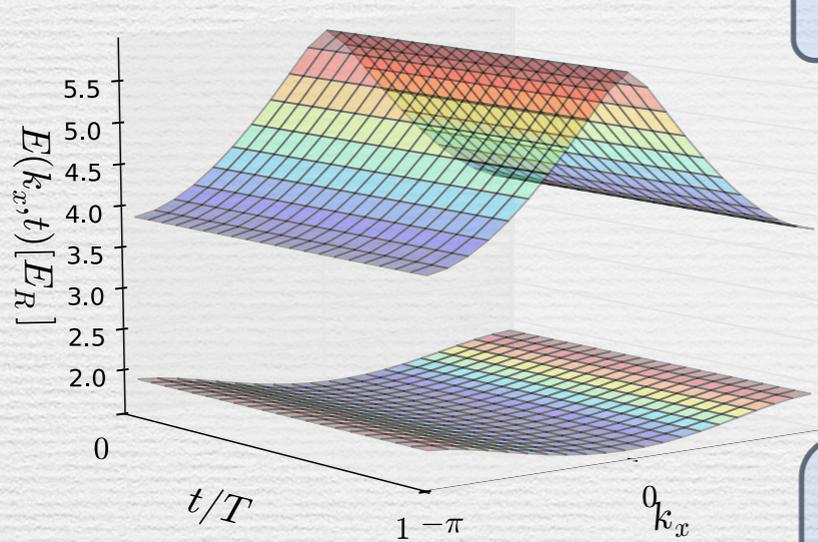


$$\mathcal{F}(k_x, t)$$



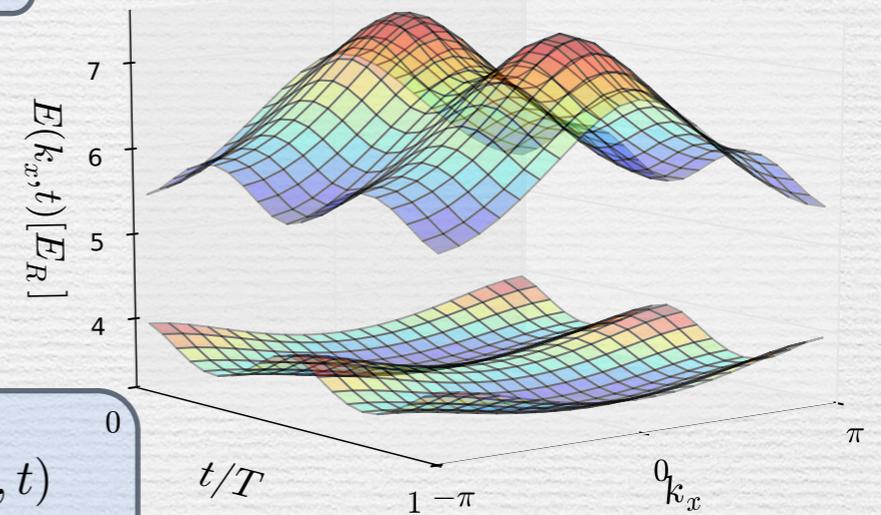
Gap & Chern number

$$V_1 = 0 \quad V_2 = 4E_R$$



$$H(x, t) = -\frac{\hbar^2}{2m} \nabla^2 + V_{OL}(x, t)$$

$$V_1 = 4E_R \quad V_2 = 4E_R$$

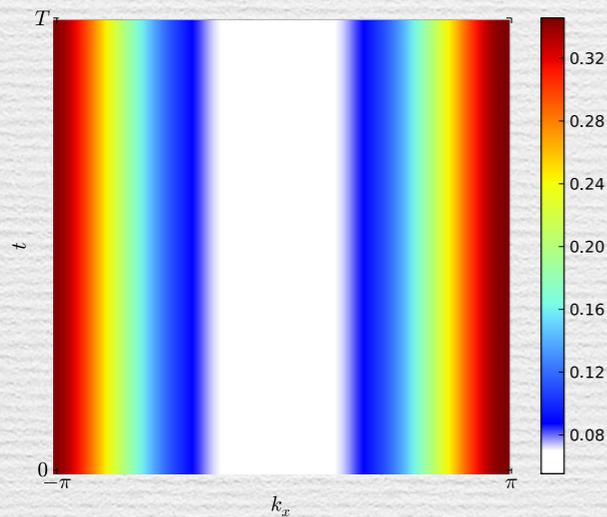


$$\Delta n = \frac{1}{2\pi} \int_0^T dt \int_0^{2\pi} dk_x \mathcal{F}(k_x, t)$$

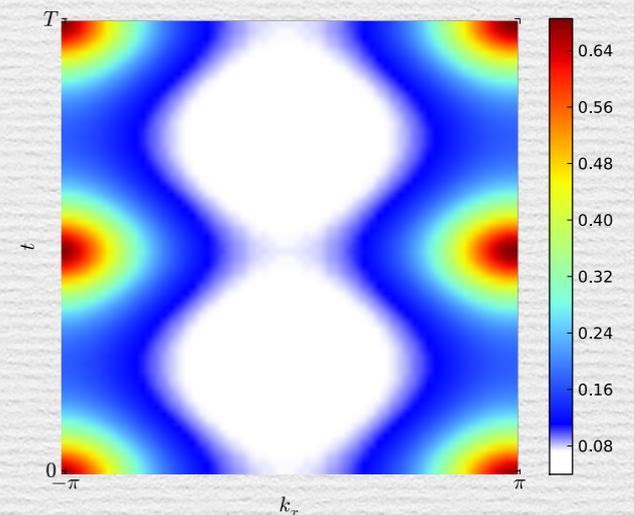
$$\mathcal{F}(k_x, t) = \partial_t A_{k_x} - \partial_{k_x} A_t$$

$$A_t(k_x) = -i \langle \psi_{k_x}(t) | \partial_{t(k_x)} | \psi_{k_x}(t) \rangle$$

$$\mathcal{F}(k_x, t)$$



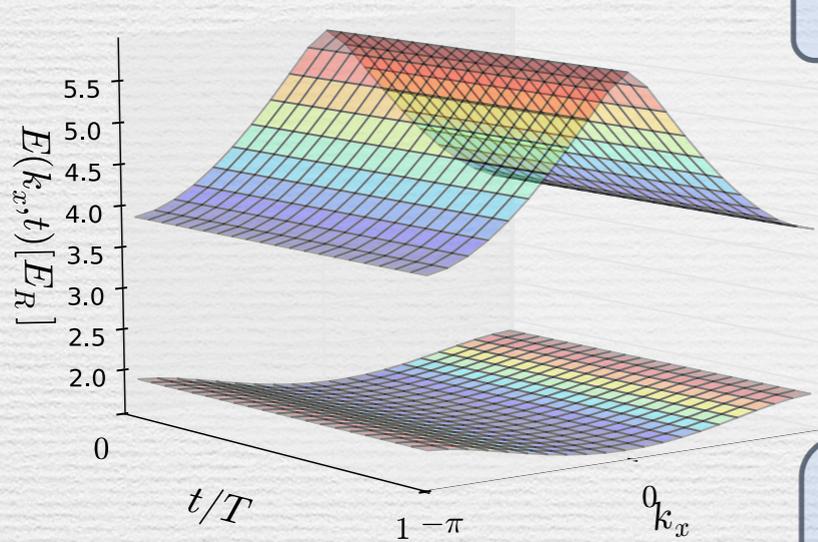
$$\mathcal{F}(k_x, t)$$



$$\Delta n = 1$$

Gap & Chern number

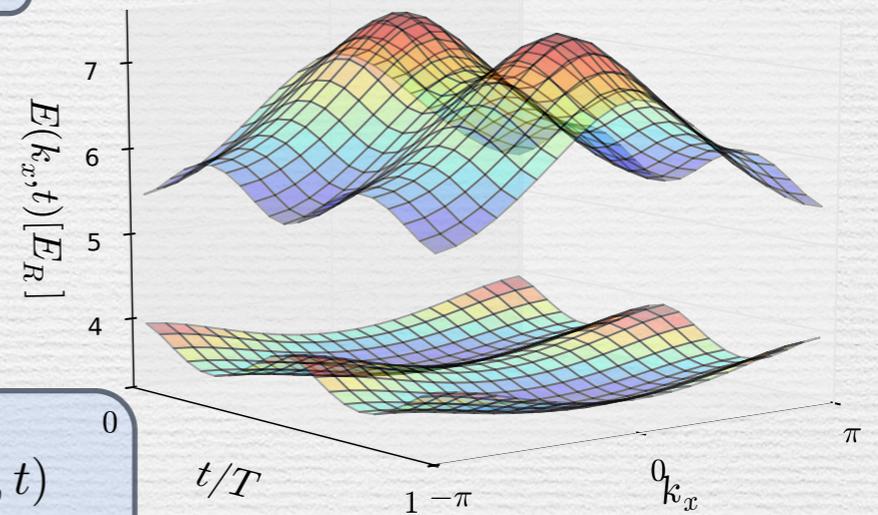
$$V_1 = 0 \quad V_2 = 4E_R$$



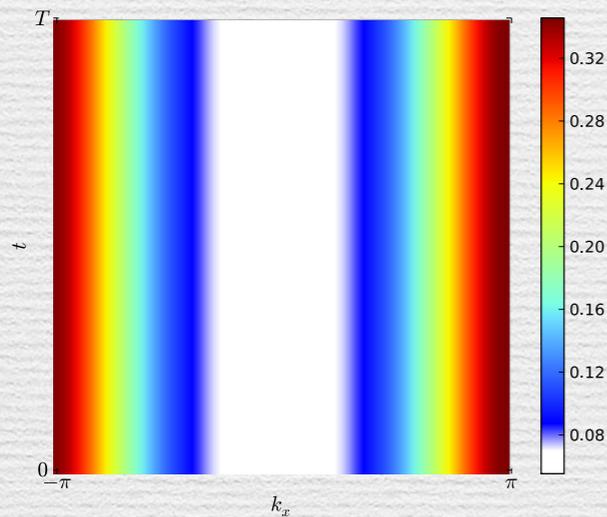
$$H(x, t) = -\frac{\hbar^2}{2m} \nabla^2 + V_{OL}(x, t)$$



$$V_1 = 4E_R \quad V_2 = 4E_R$$



$$\mathcal{F}(k_x, t)$$

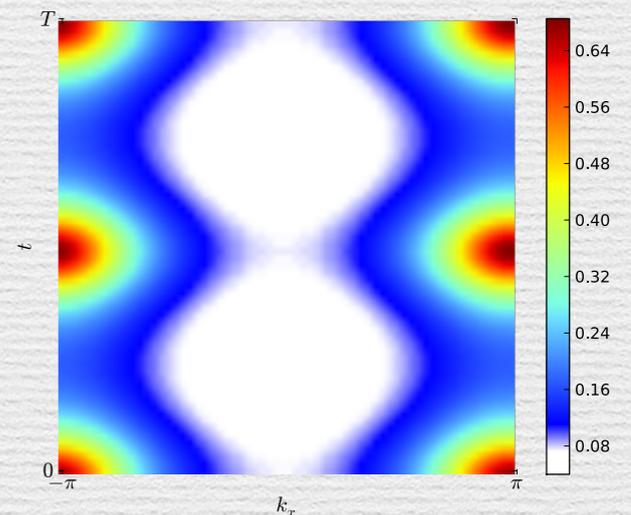


$$\Delta n = \frac{1}{2\pi} \int_0^T dt \int_0^{2\pi} dk_x \mathcal{F}(k_x, t)$$

$$\mathcal{F}(k_x, t) = \partial_t A_{k_x} - \partial_{k_x} A_t$$

$$A_t(k_x) = -i \langle \psi_{k_x}(t) | \partial_{t(k_x)} | \psi_{k_x}(t) \rangle$$

$$\mathcal{F}(k_x, t)$$

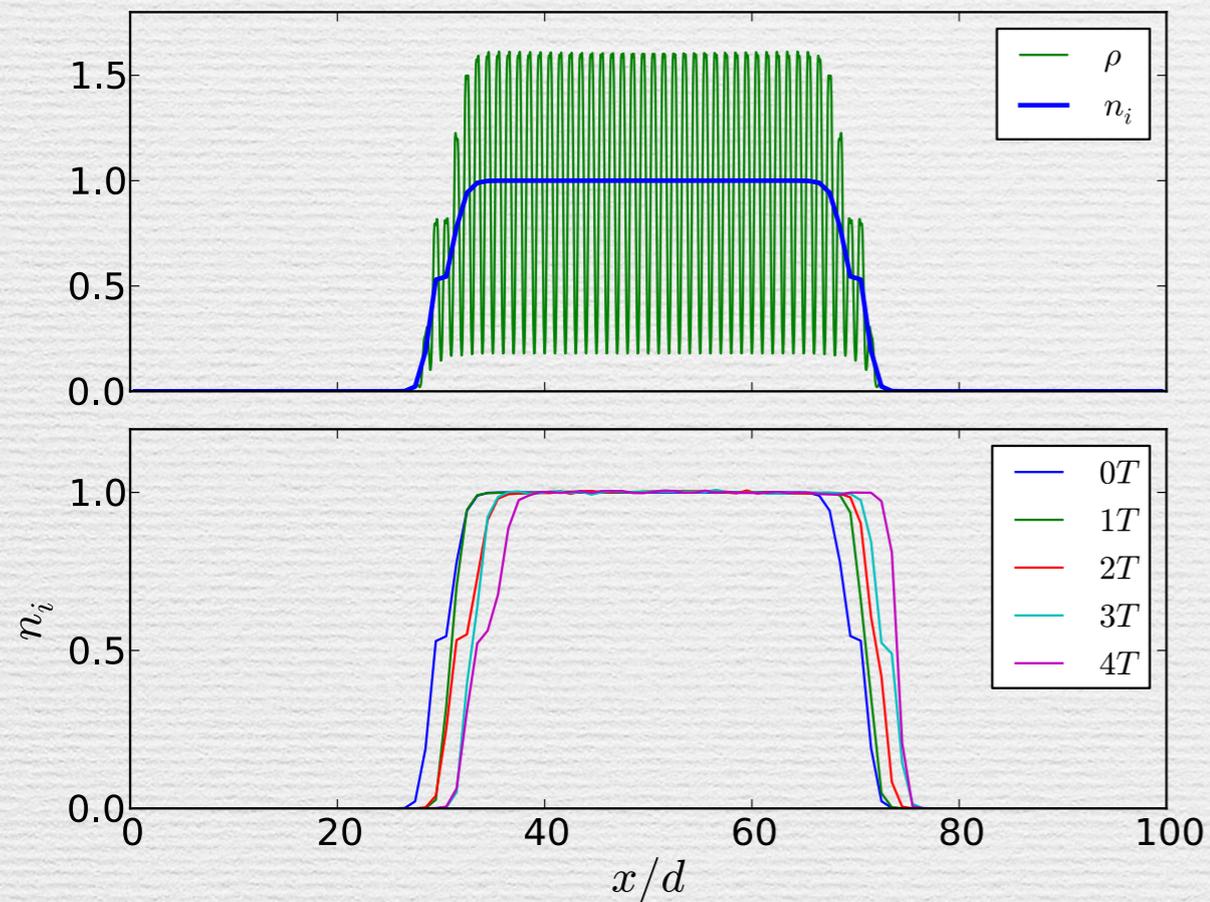


Topological pump
 $\Delta n = 1$

Practical issues

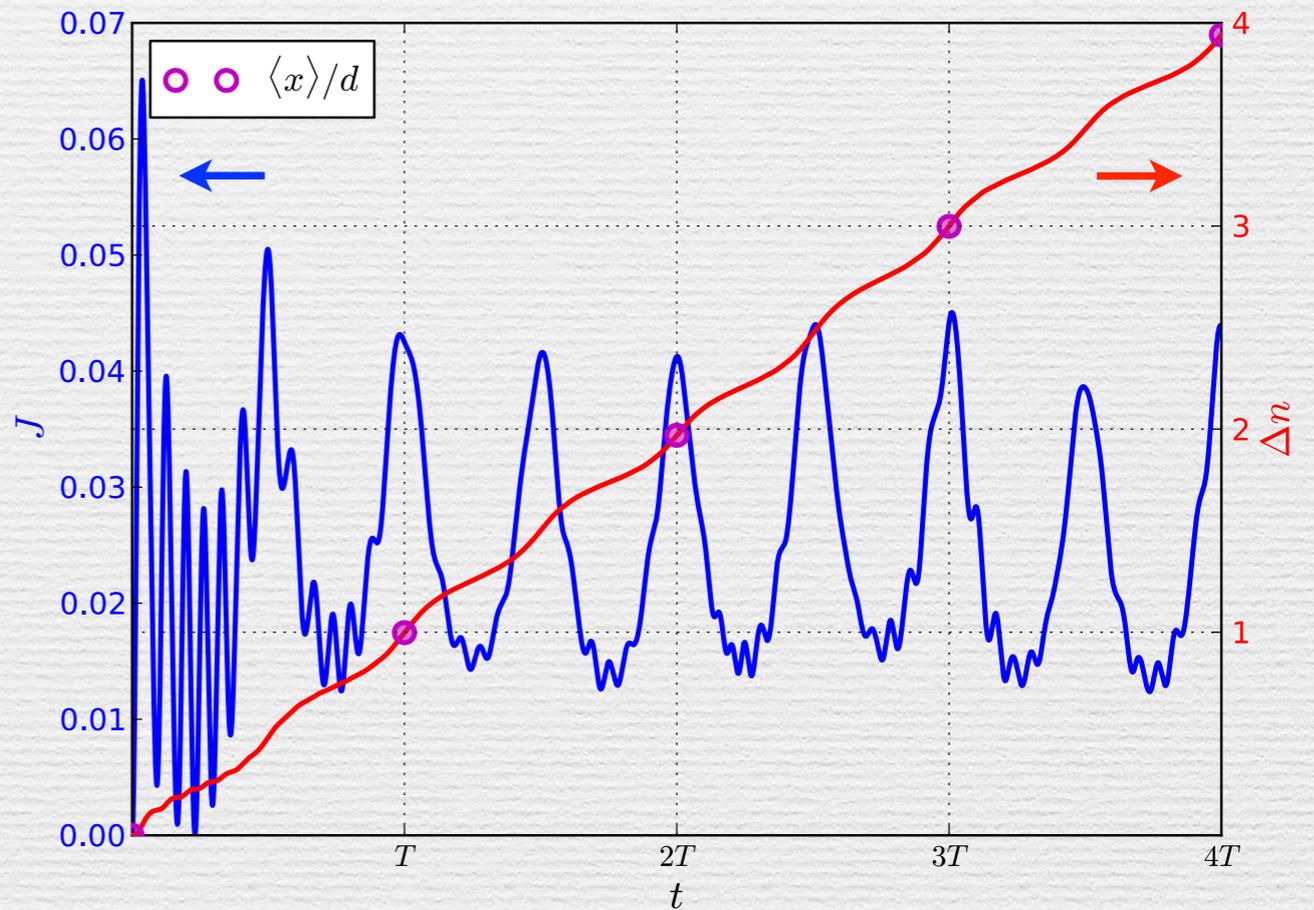
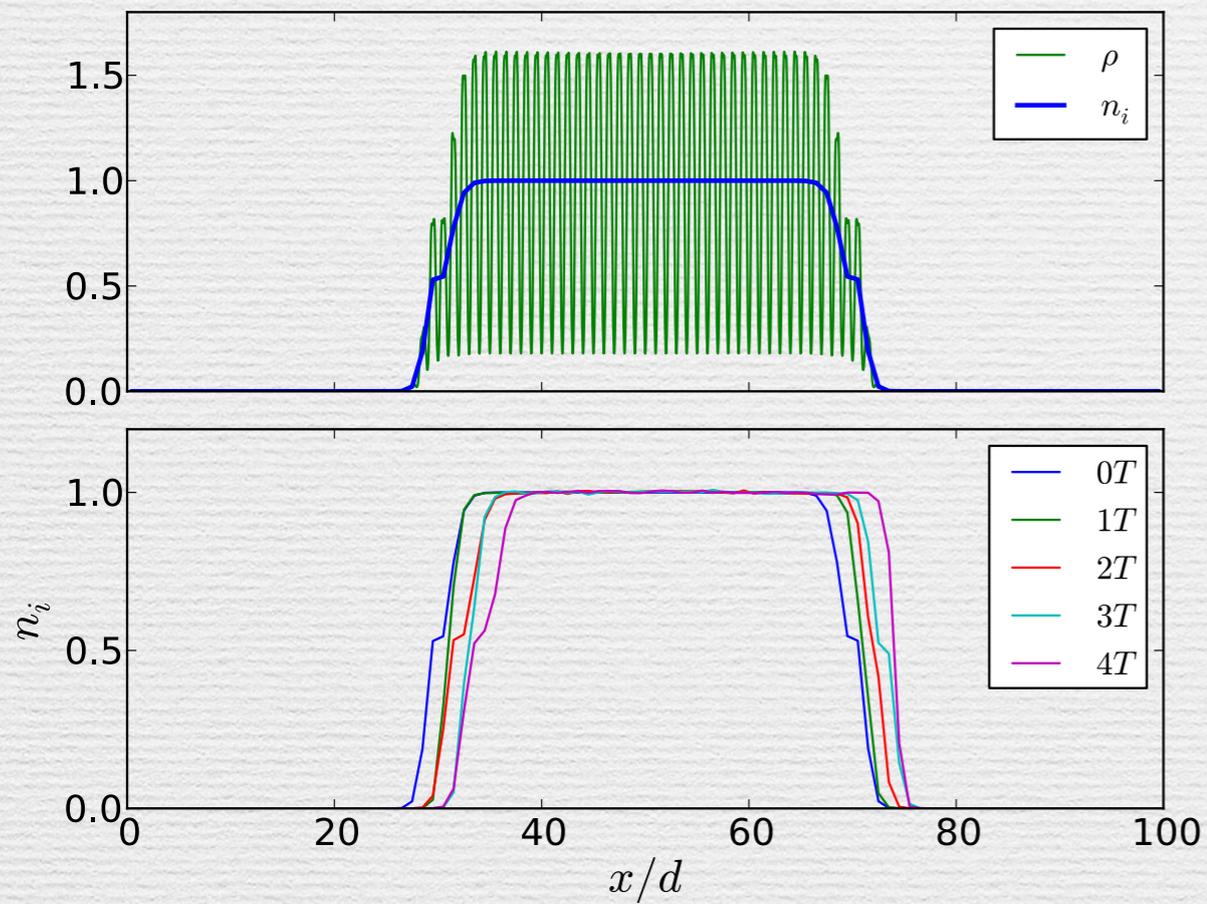
- Detection
- External trap
- Temperature effect
- Non-adiabatic effect

Trapping & Detection



$$\langle x \rangle / d = \Delta n$$

Trapping & Detection



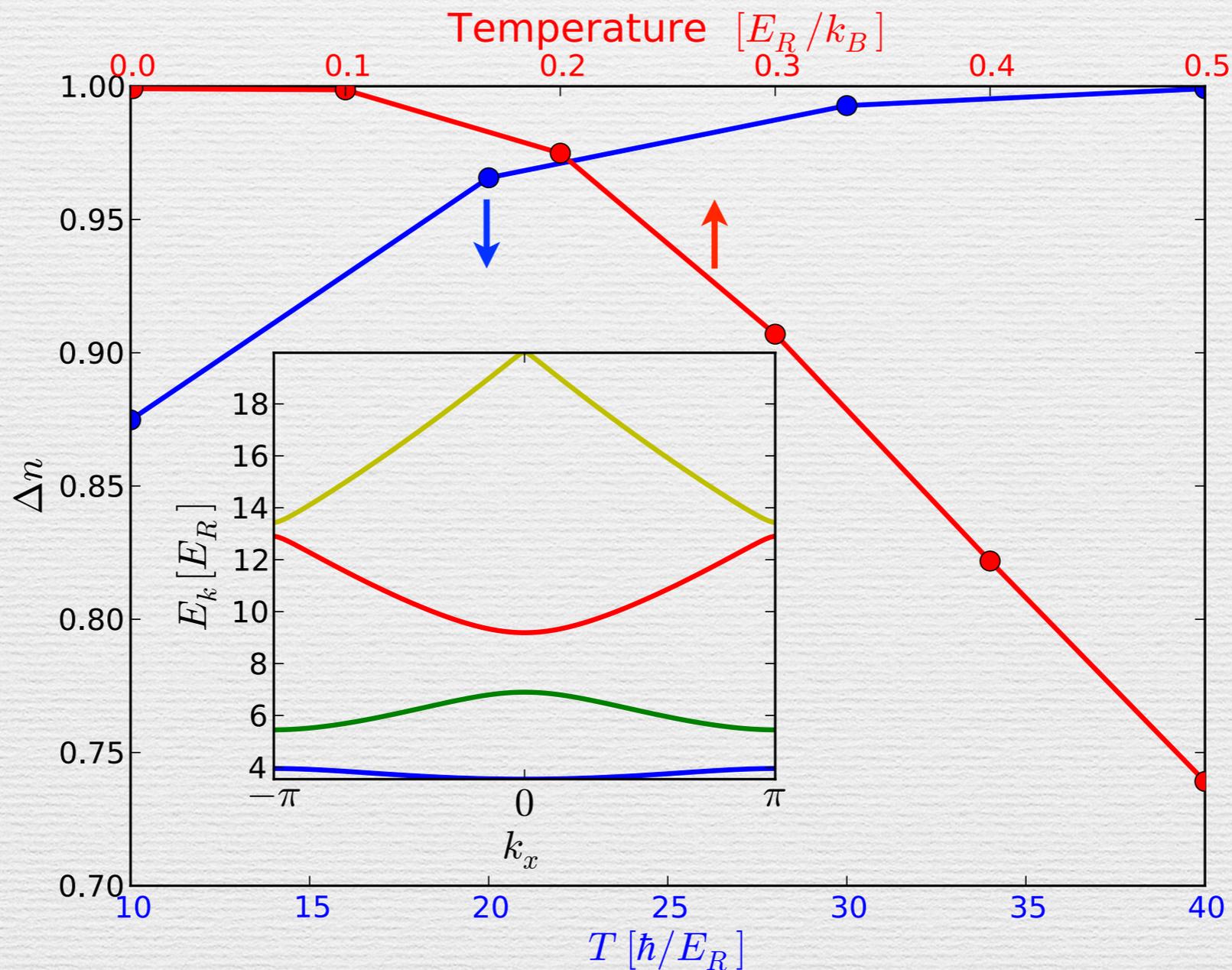
$$\langle x \rangle / d = \Delta n$$

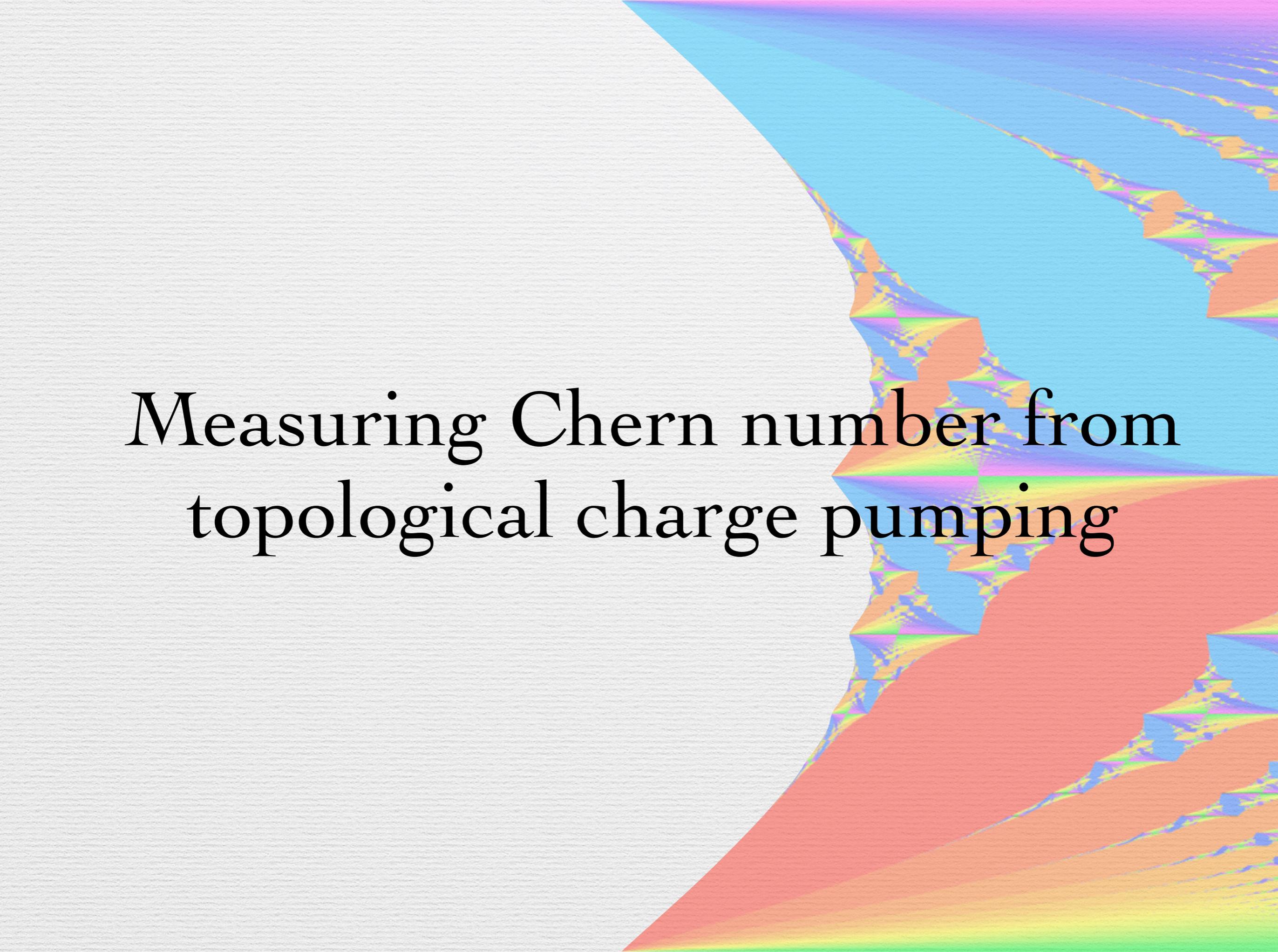
Temperature & Non-adiabatic effect

$$\text{Temperature} \ll \frac{\Delta}{k_B} \qquad T \gg \frac{\hbar}{\Delta}$$

Temperature & Non-adiabatic effect

$$\text{Temperature} \ll \frac{\Delta}{k_B} \quad T \gg \frac{\hbar}{\Delta}$$





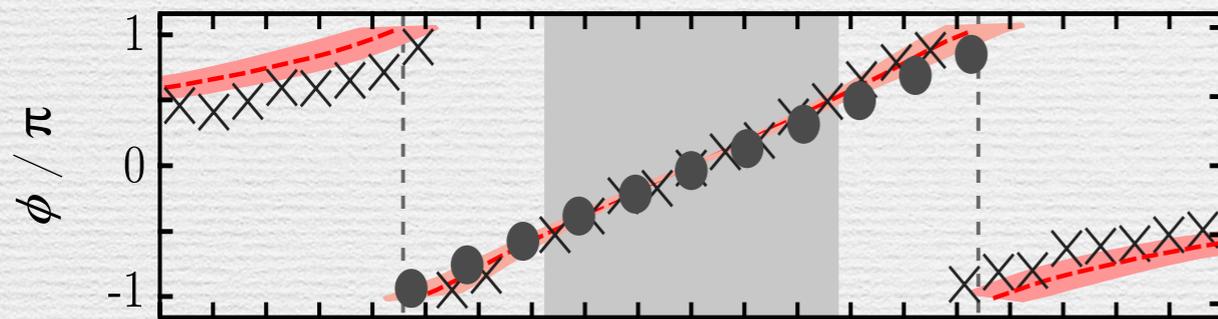
Measuring Chern number from
topological charge pumping

Synthetic gauge-field in optical lattices

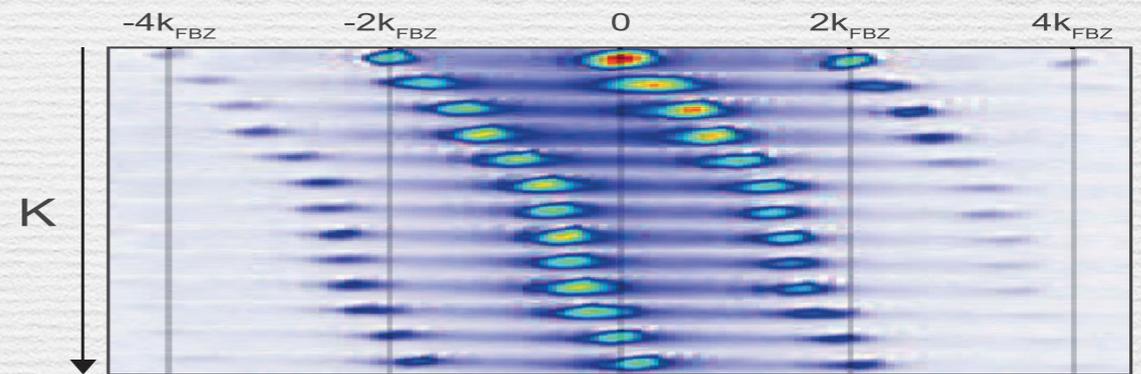
• Imprint **complex phases** to the hopping amplitude

• 1D Peierls lattice **NIST, Hamburg** $H = -J \sum_m e^{i2\pi\Phi} c_{m+1}^\dagger c_m + H.c.$

(a) Peierls tunneling phase

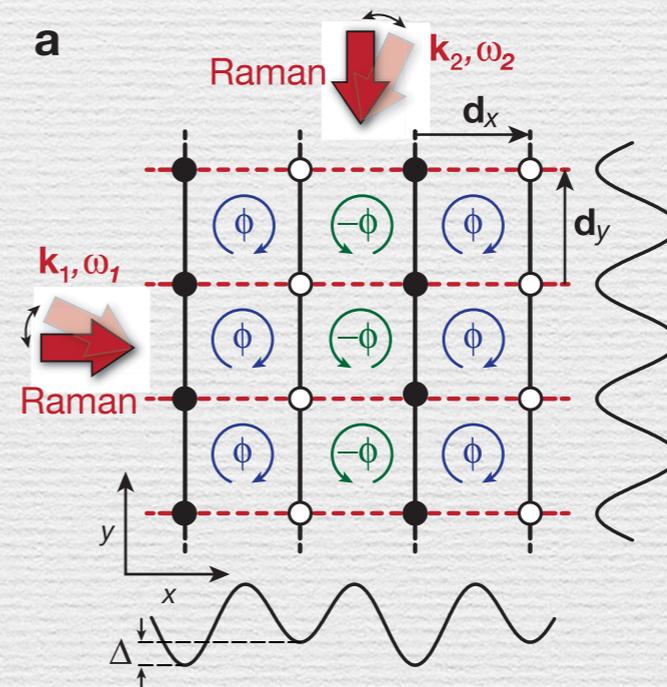


Jimenez-Garcia et al



Struck et al

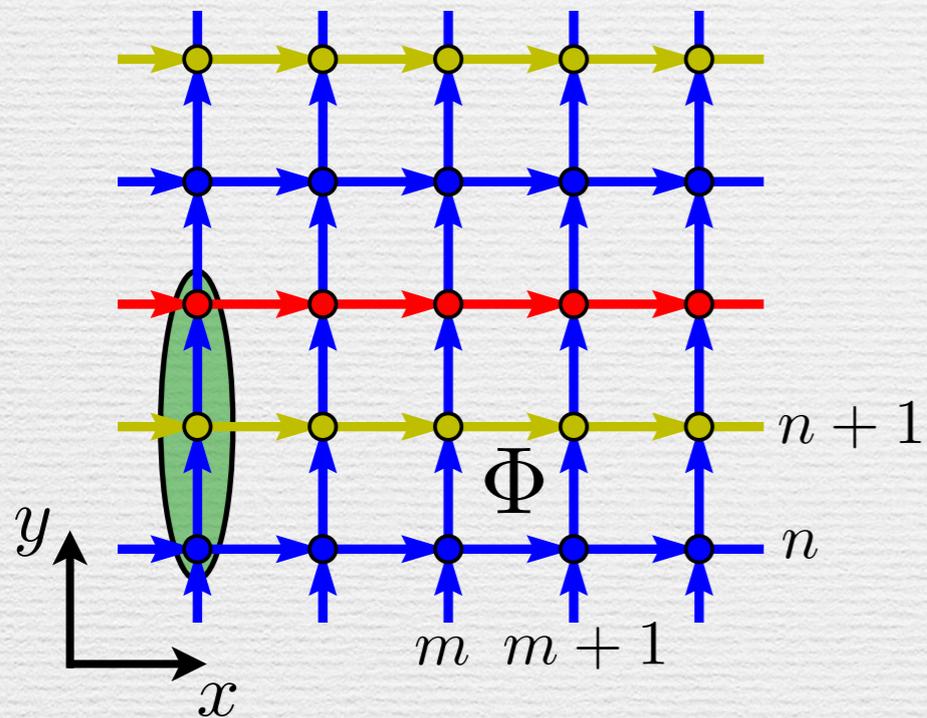
• Staggered flux lattice **Munich**



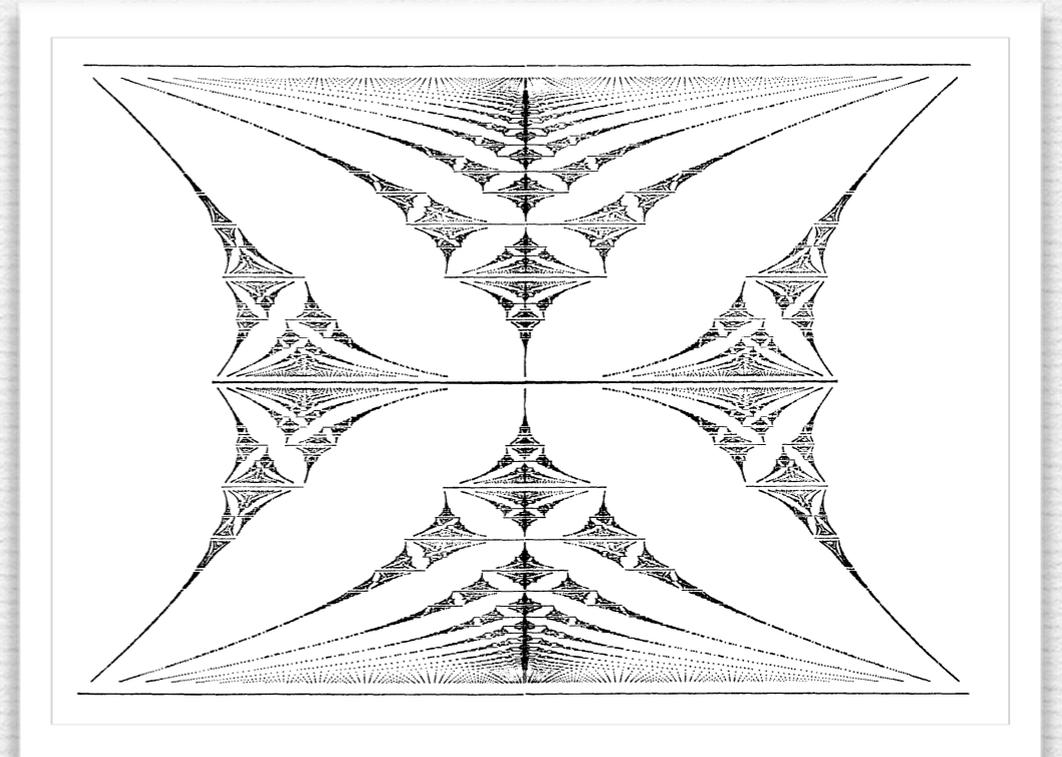
Aidelsburger et al

Hofstadter optical lattice

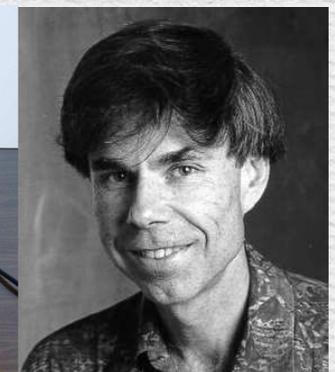
$$H = -J \sum_{m,n} e^{i2\pi n\Phi} c_{m+1,n}^\dagger c_{m,n} + c_{m,n+1}^\dagger c_{m,n} + H.c. \quad \Phi = p/q$$



Hofstadter, 1976

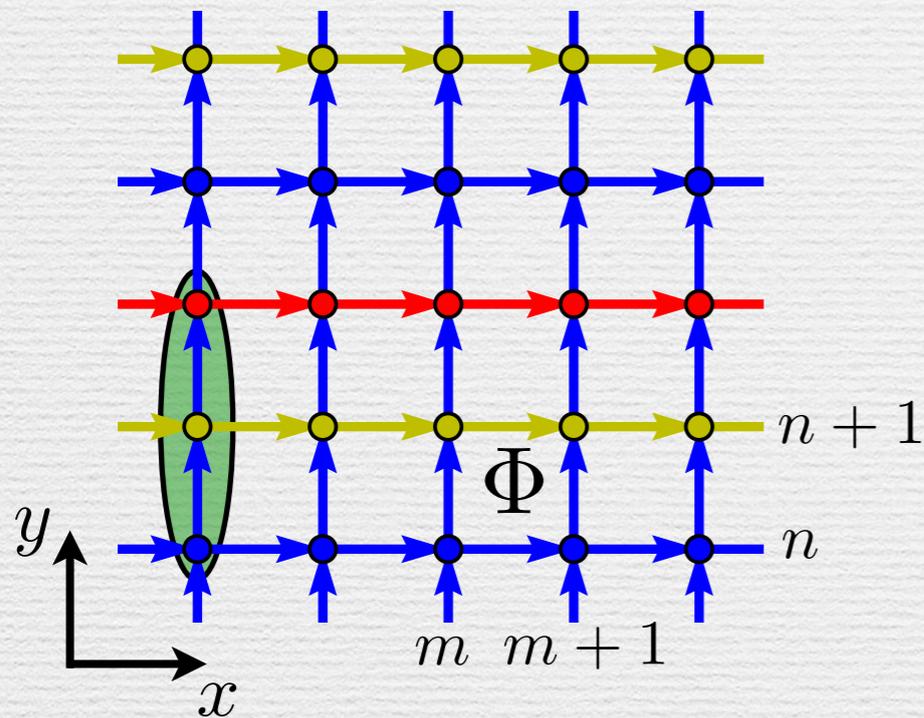


Φ
 E

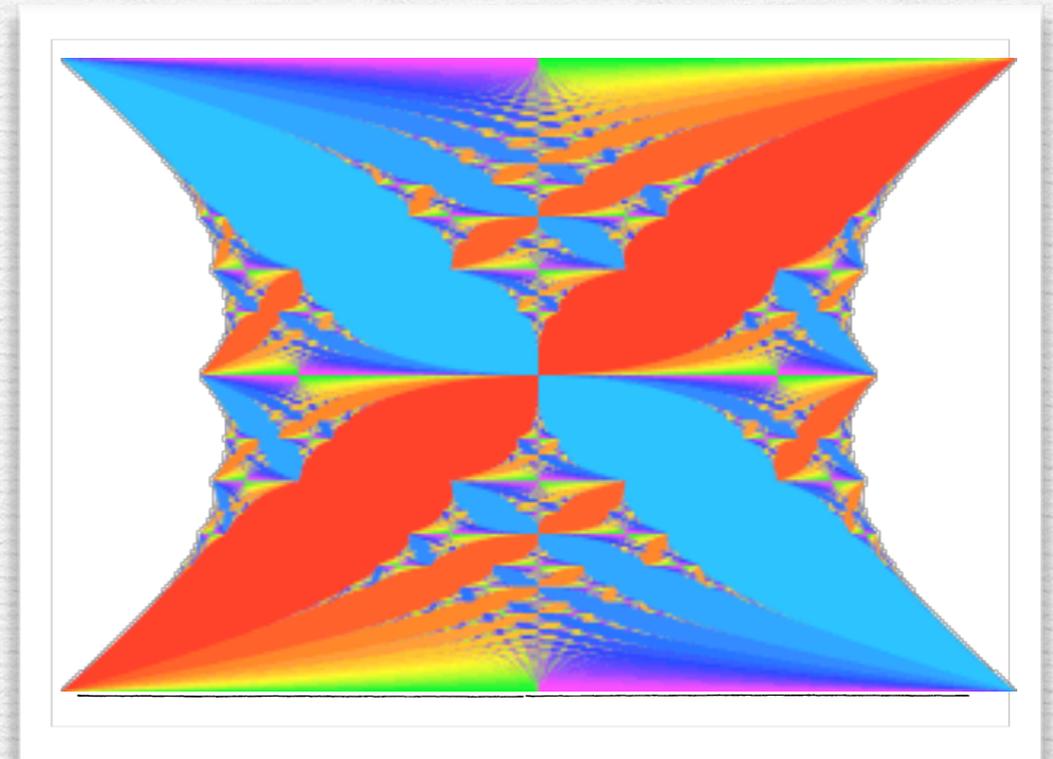


Hofstadter optical lattice

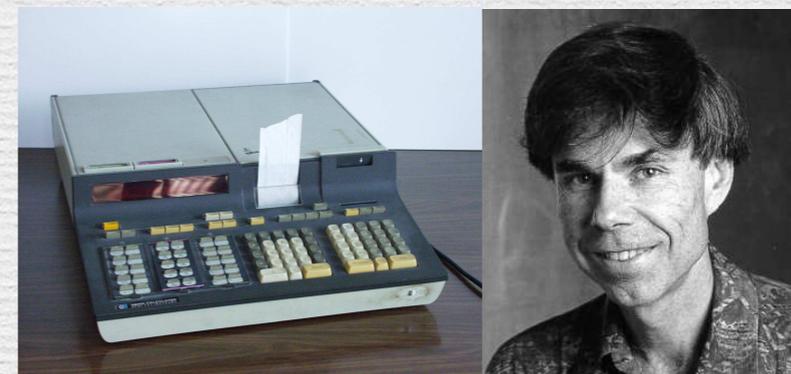
$$H = -J \sum_{m,n} e^{i2\pi n\Phi} c_{m+1,n}^\dagger c_{m,n} + c_{m,n+1}^\dagger c_{m,n} + H.c. \quad \Phi = p/q$$



Hofstadter, 1976

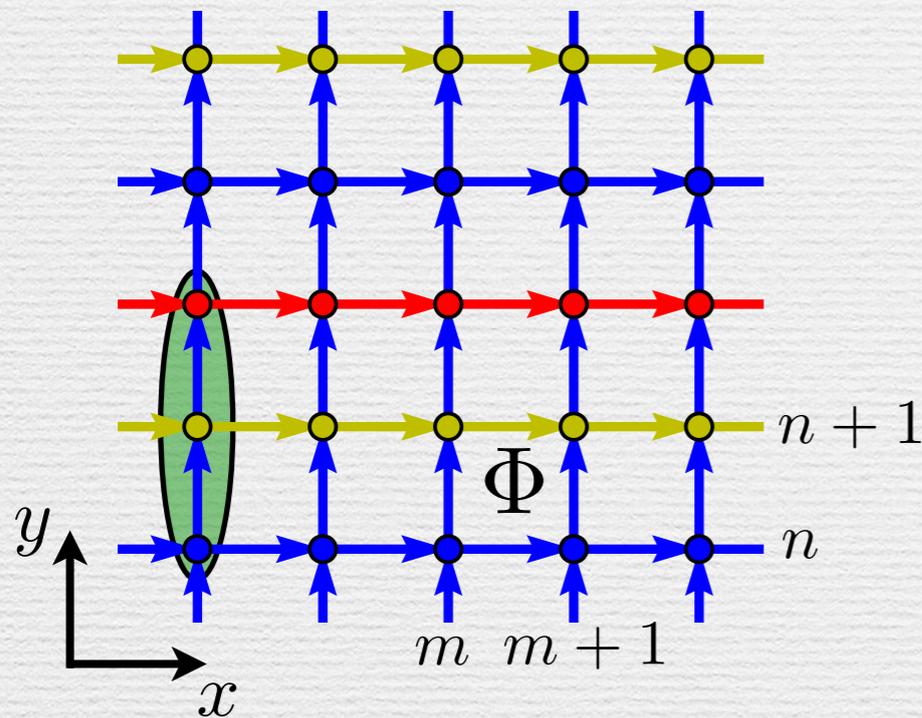


Osadchy and Avron, J. Math. Phys. 2001

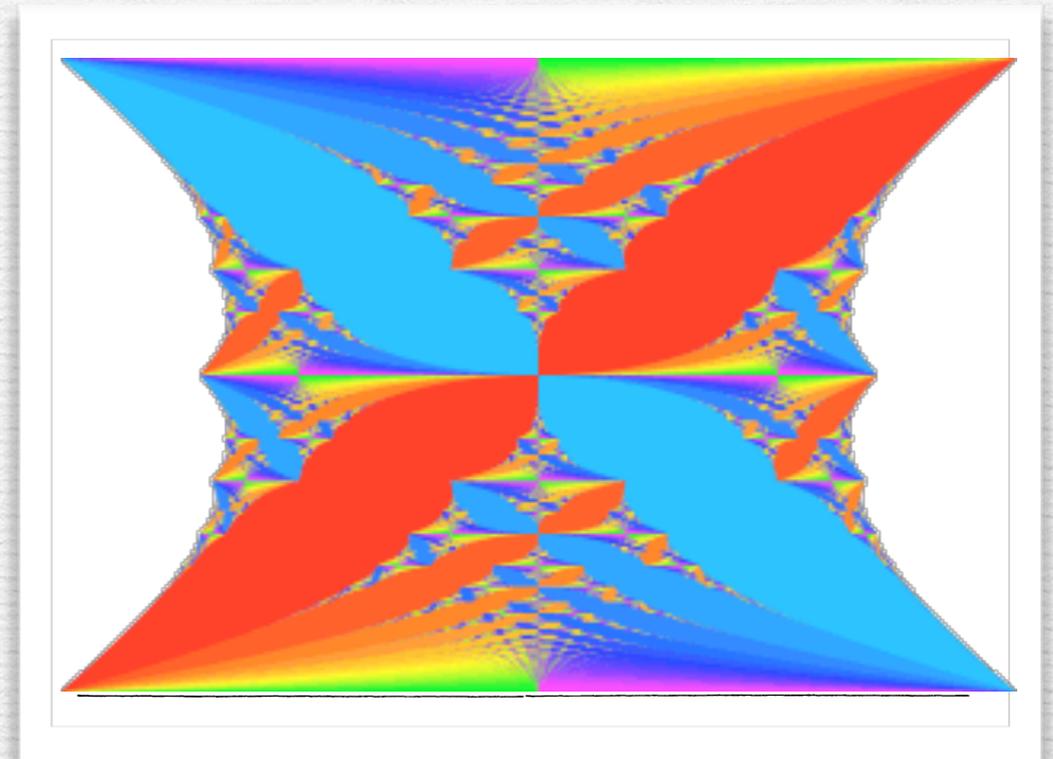


Hofstadter optical lattice

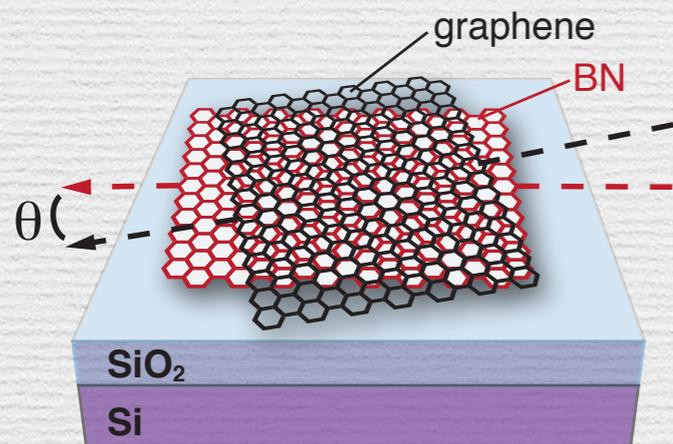
$$H = -J \sum_{m,n} e^{i2\pi n\Phi} c_{m+1,n}^\dagger c_{m,n} + c_{m,n+1}^\dagger c_{m,n} + H.c. \quad \Phi = p/q$$



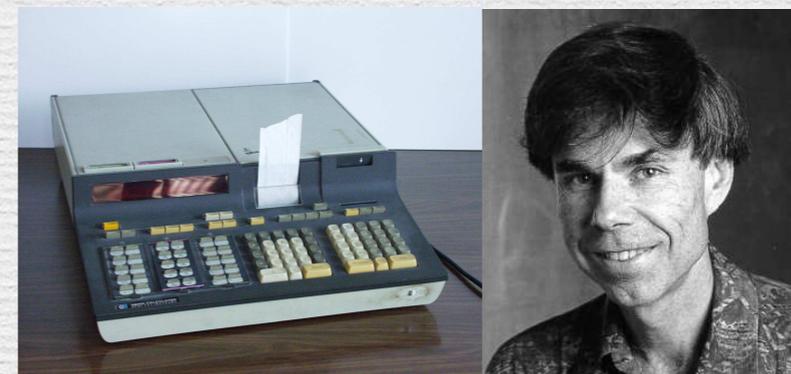
Hofstadter, 1976



Osadchy and Avron, J. Math. Phys. 2001



arXiv: 1212.4783 Hofstadter's butterfly in moiré superlattices: A fractal quantum Hall effect

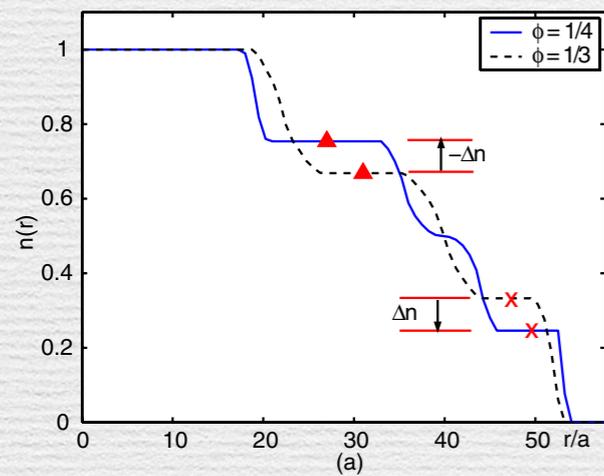


How to measure Chern # ?

How to measure Chern # ?

Density profile

Umucalilar *et al*



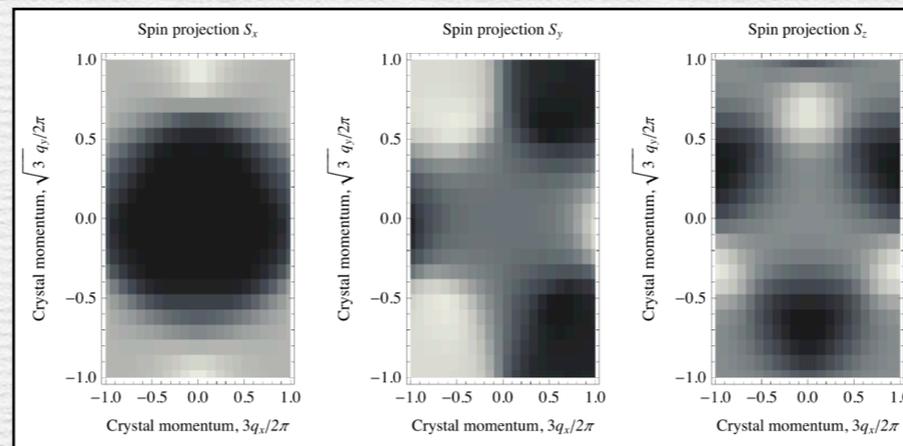
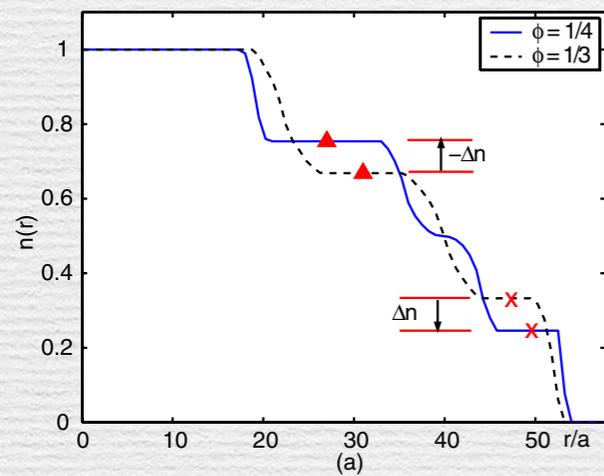
How to measure Chern # ?

Time-of-flight

Alba et al, Zhao et al

Density profile

Umucalilar et al



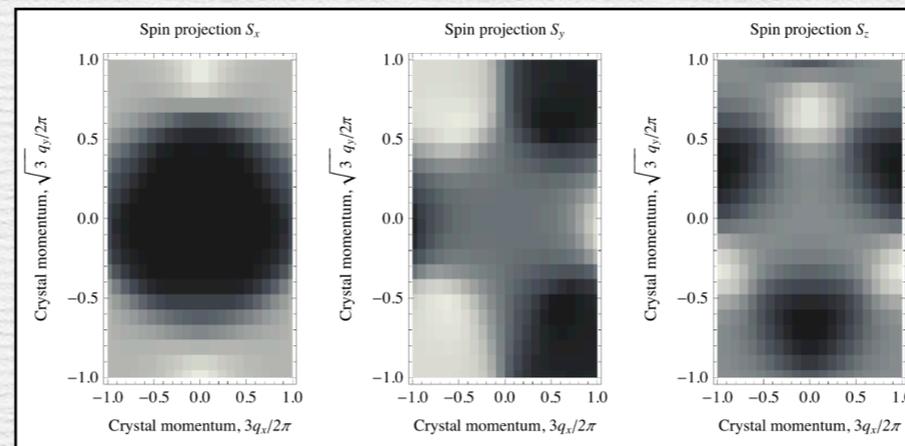
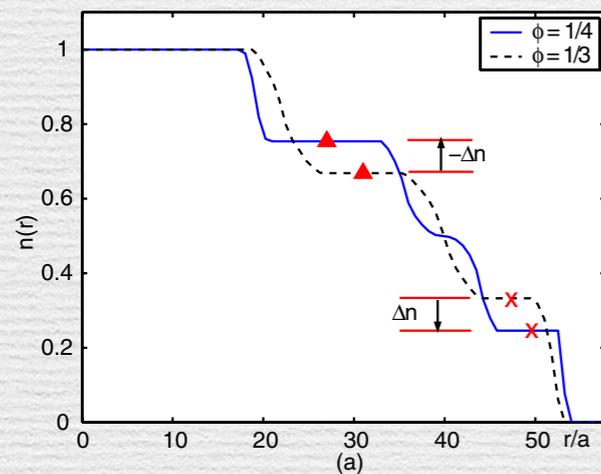
How to measure Chern # ?

Time-of-flight

Alba et al, Zhao et al

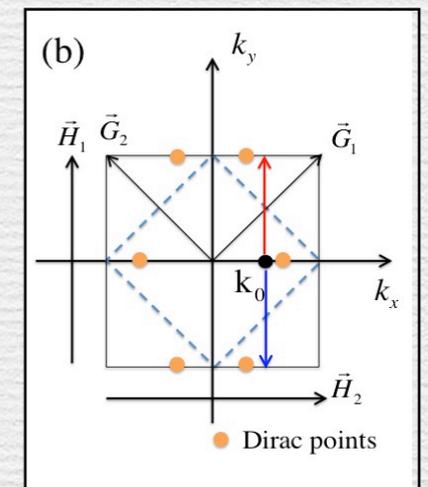
Density profile

Umucalilar et al



Zak phases

Abanin et al



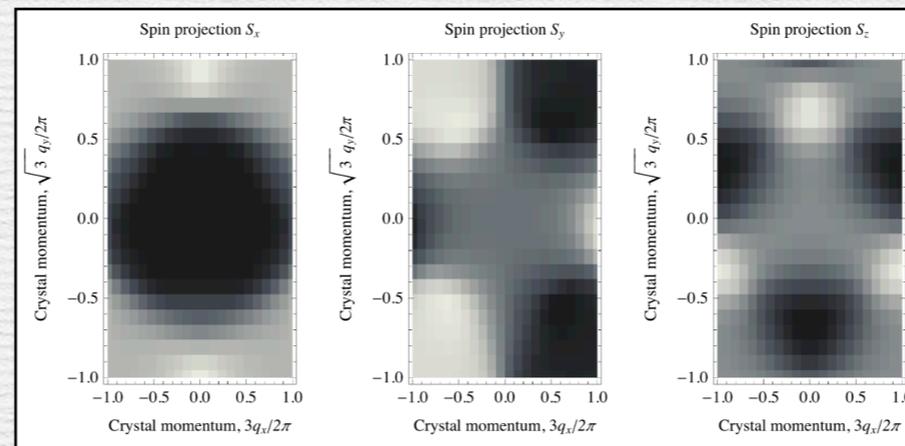
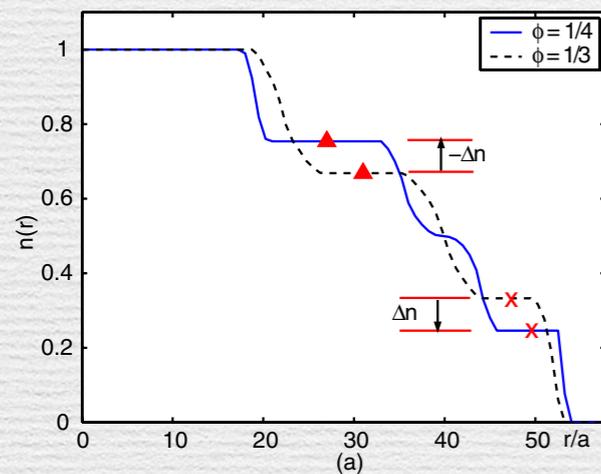
How to measure Chern # ?

Time-of-flight

Alba et al, Zhao et al

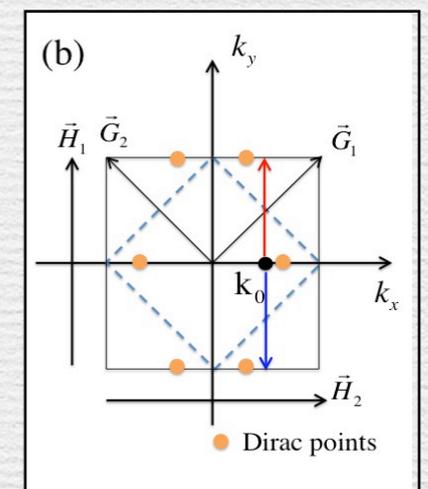
Density profile

Umucalilar et al



Zak phases

Abanin et al



Semi-classical dynamics

Price et al

$$\dot{\mathbf{r}} = \frac{1}{\hbar} \frac{\partial \varepsilon(\mathbf{k})}{\partial \mathbf{k}} - \frac{\mathbf{F}}{\hbar} \times \hat{\mathbf{z}} \Omega(\mathbf{k})$$

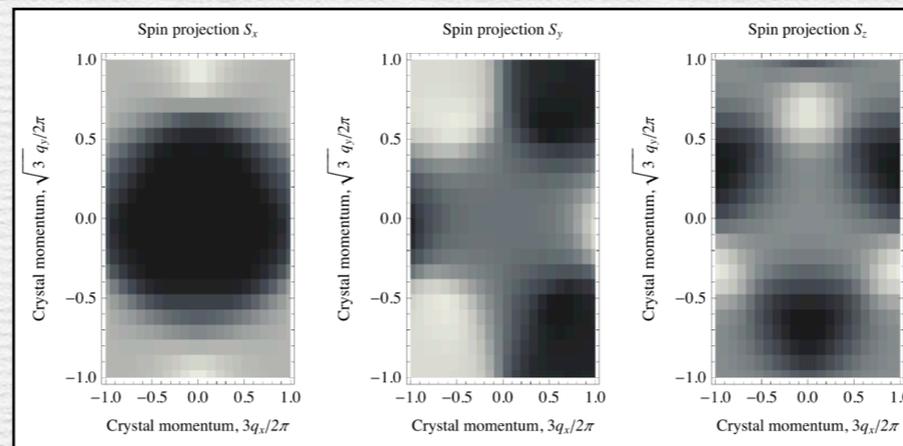
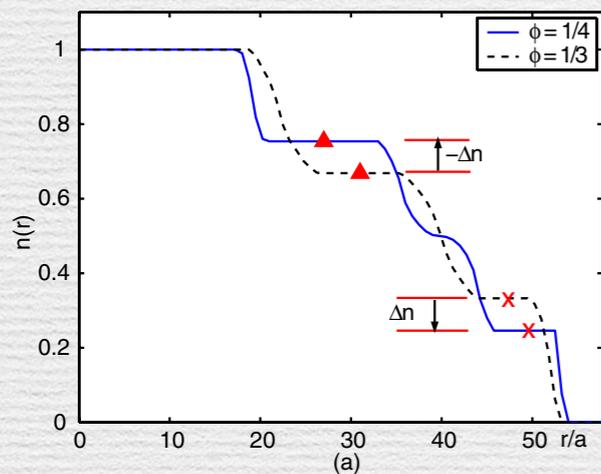
How to measure Chern # ?

Time-of-flight

Alba et al, Zhao et al

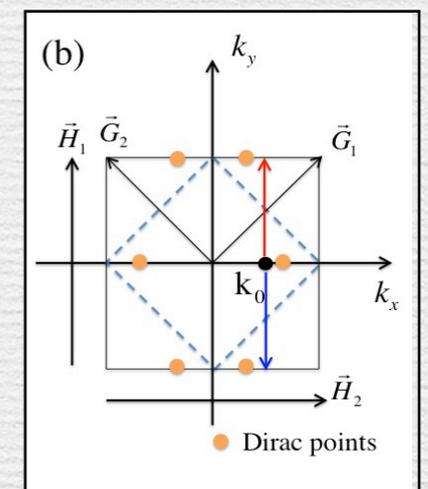
Density profile

Umucalilar et al



Zak phases

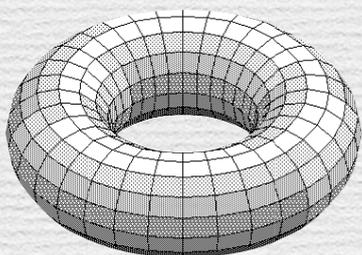
Abanin et al



Semi-classical dynamics

Price et al

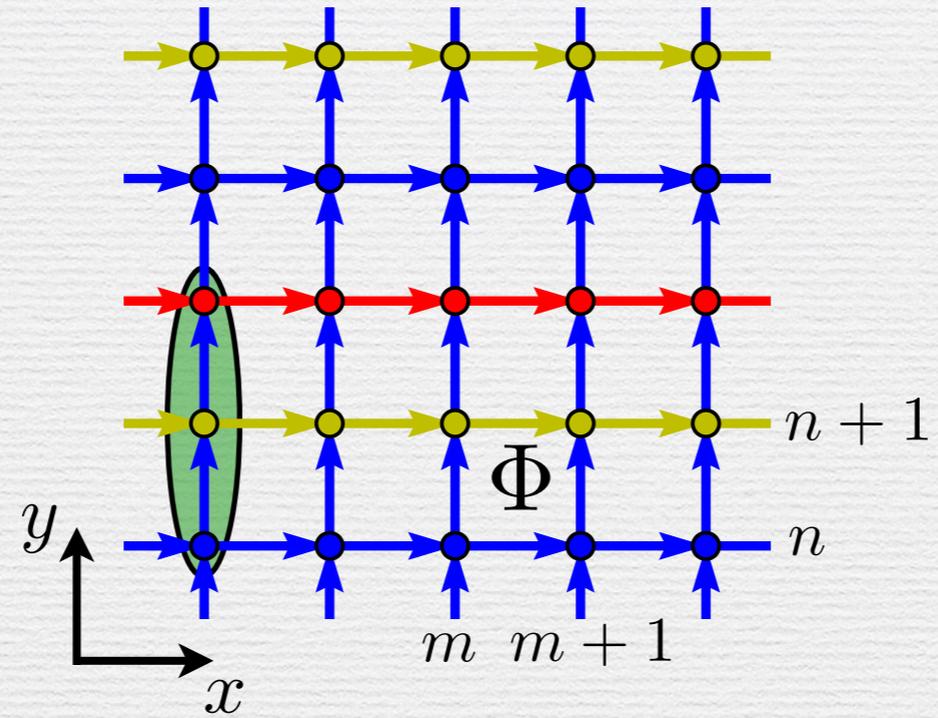
$$\dot{\mathbf{r}} = \frac{1}{\hbar} \frac{\partial \varepsilon(\mathbf{k})}{\partial \mathbf{k}} - \frac{\mathbf{F}}{\hbar} \times \hat{\mathbf{z}} \Omega(\mathbf{k})$$



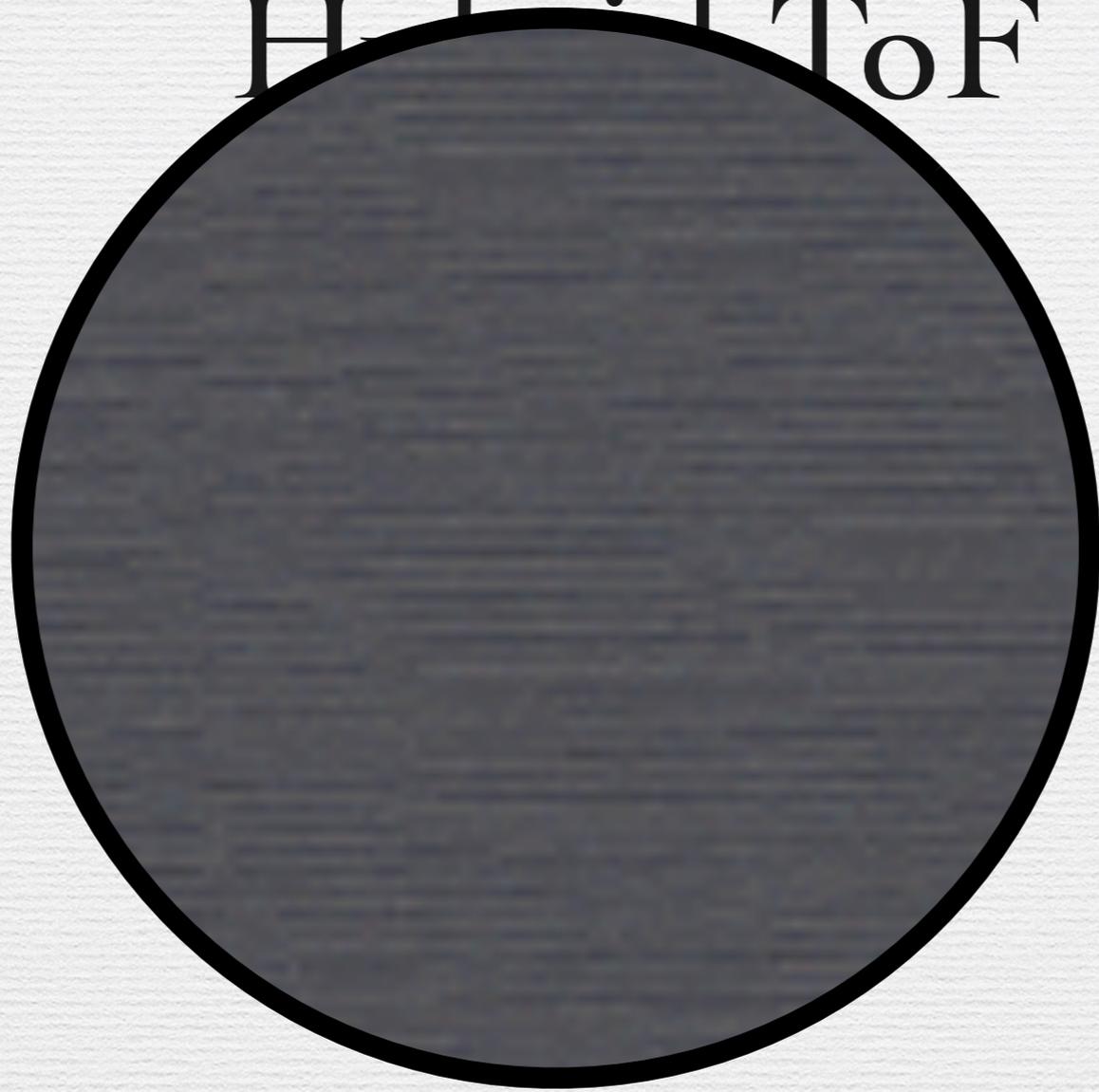
We propose a **new** probe based on
Topological Pumping Effect

$$\rho(k_x, y)$$

Hybrid ToF



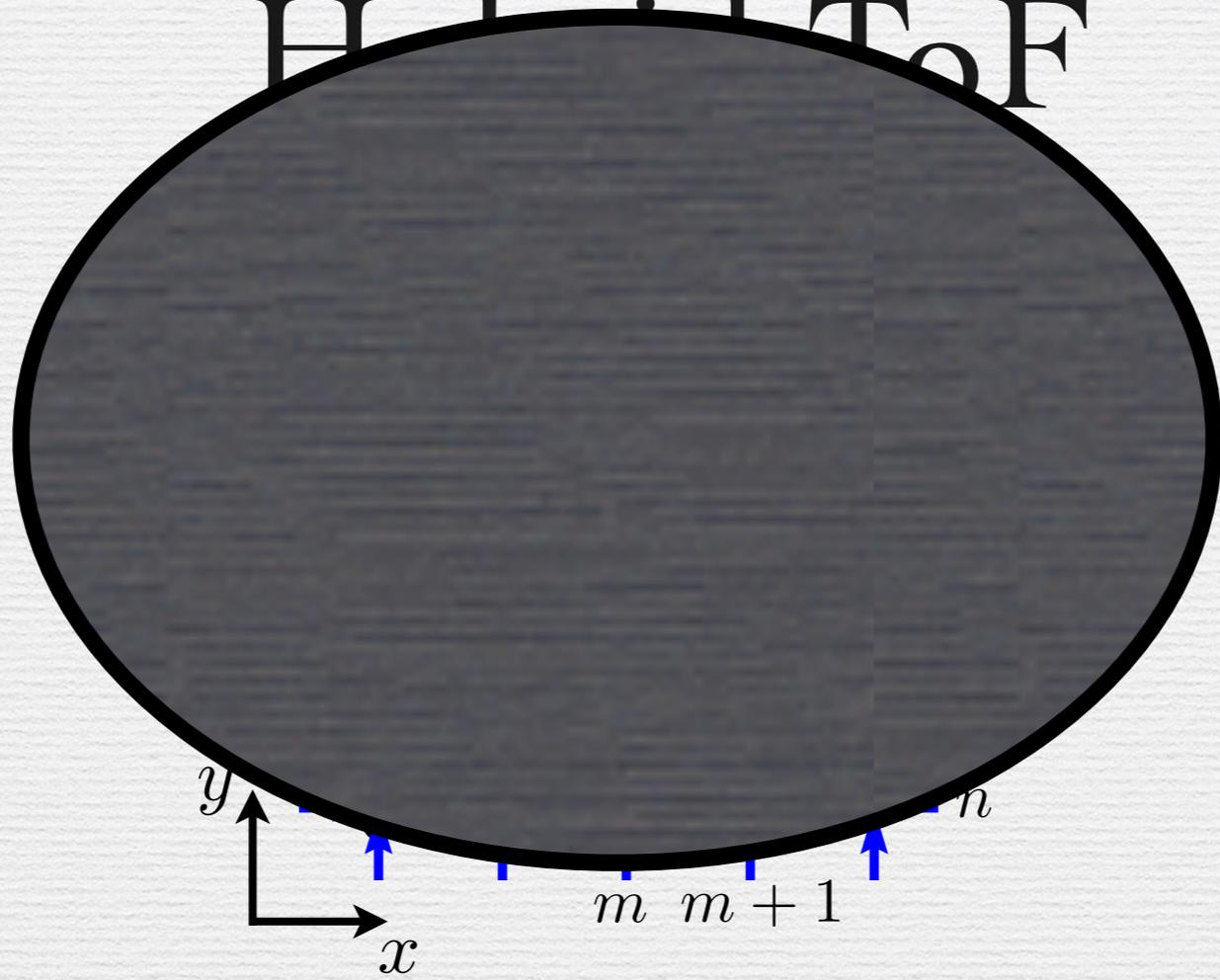
H₁ · 1 ToF



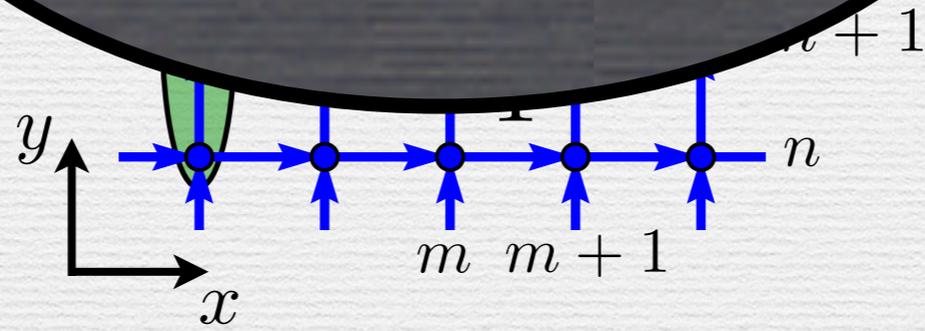
H₁ · 1 ToF



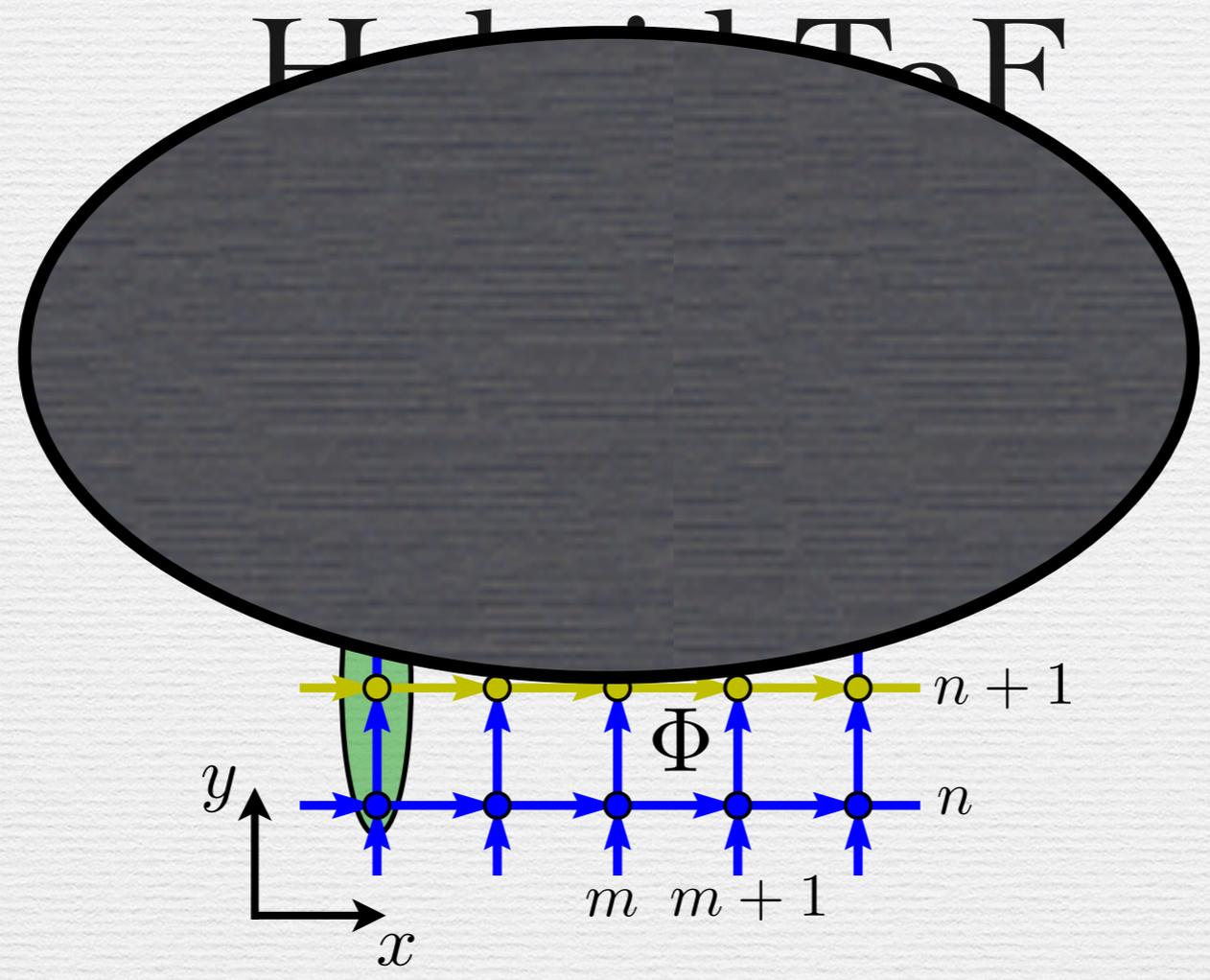
Hilbert ToF

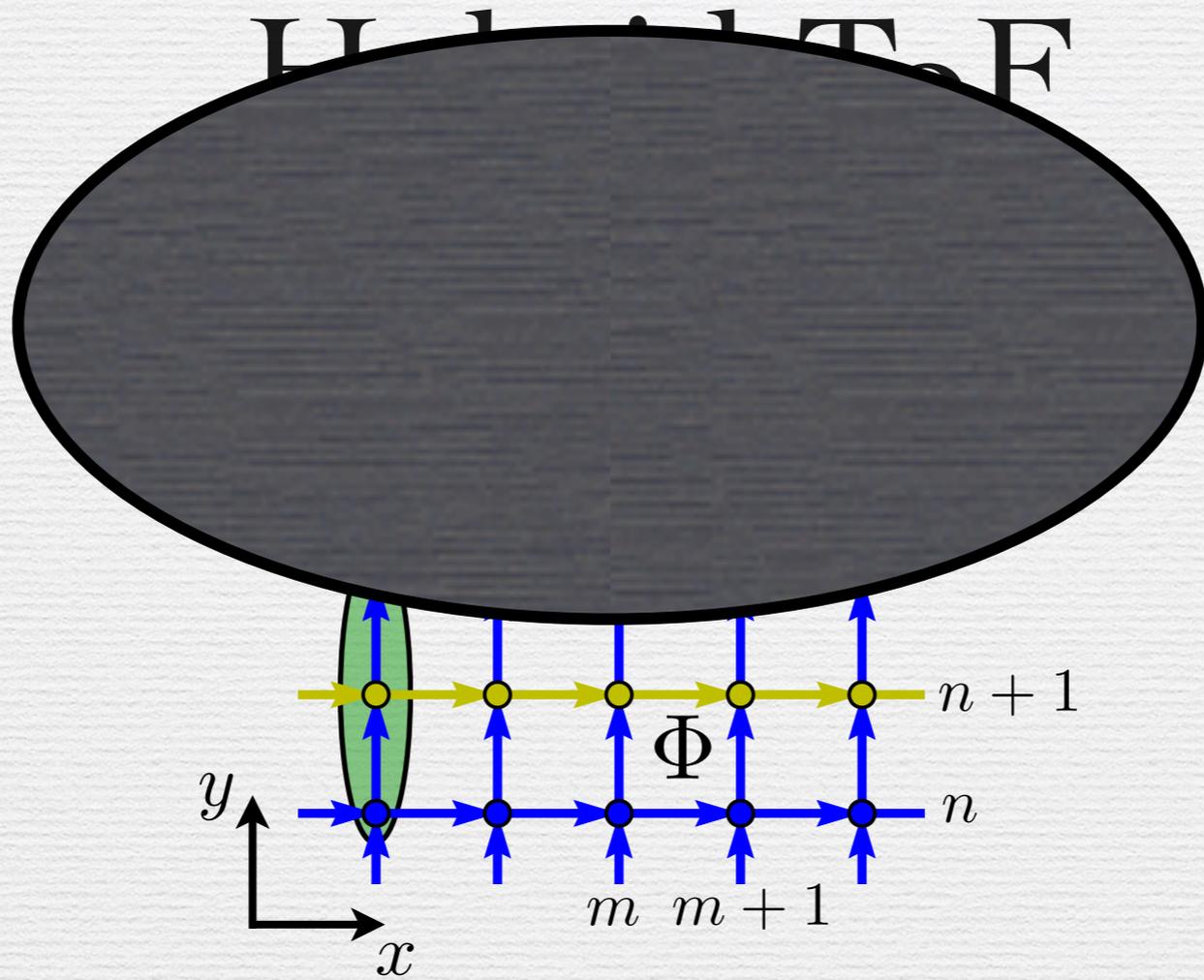


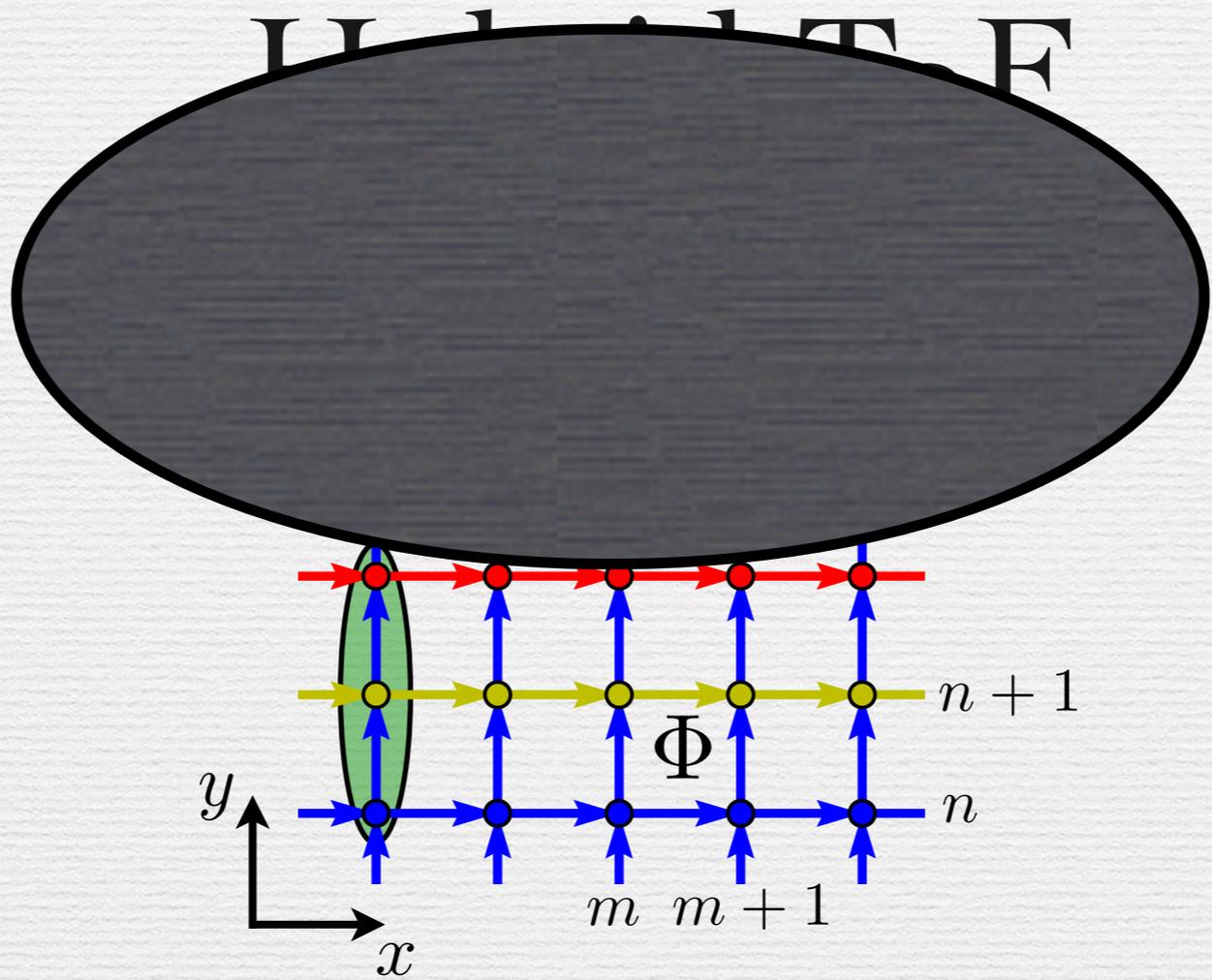
Hilbert

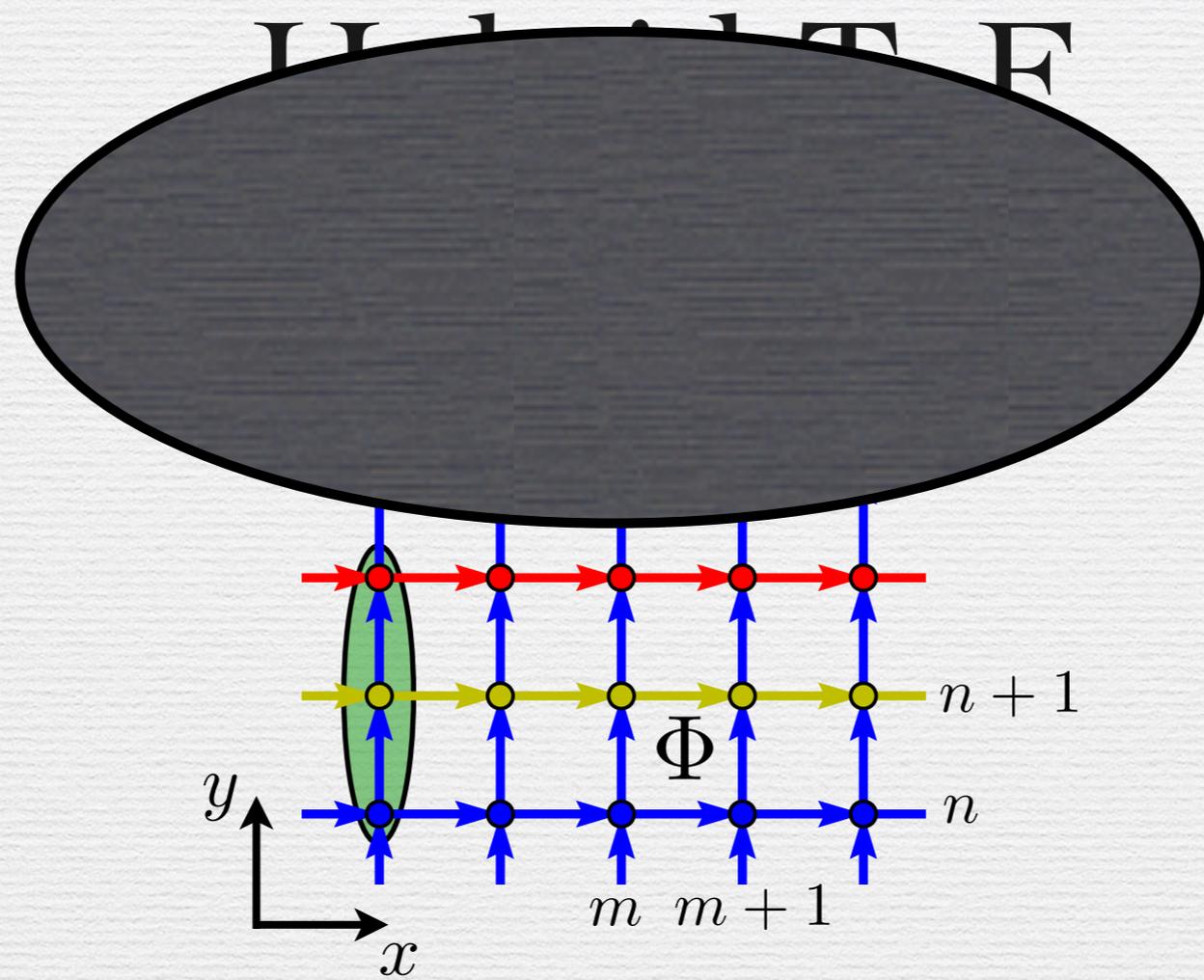


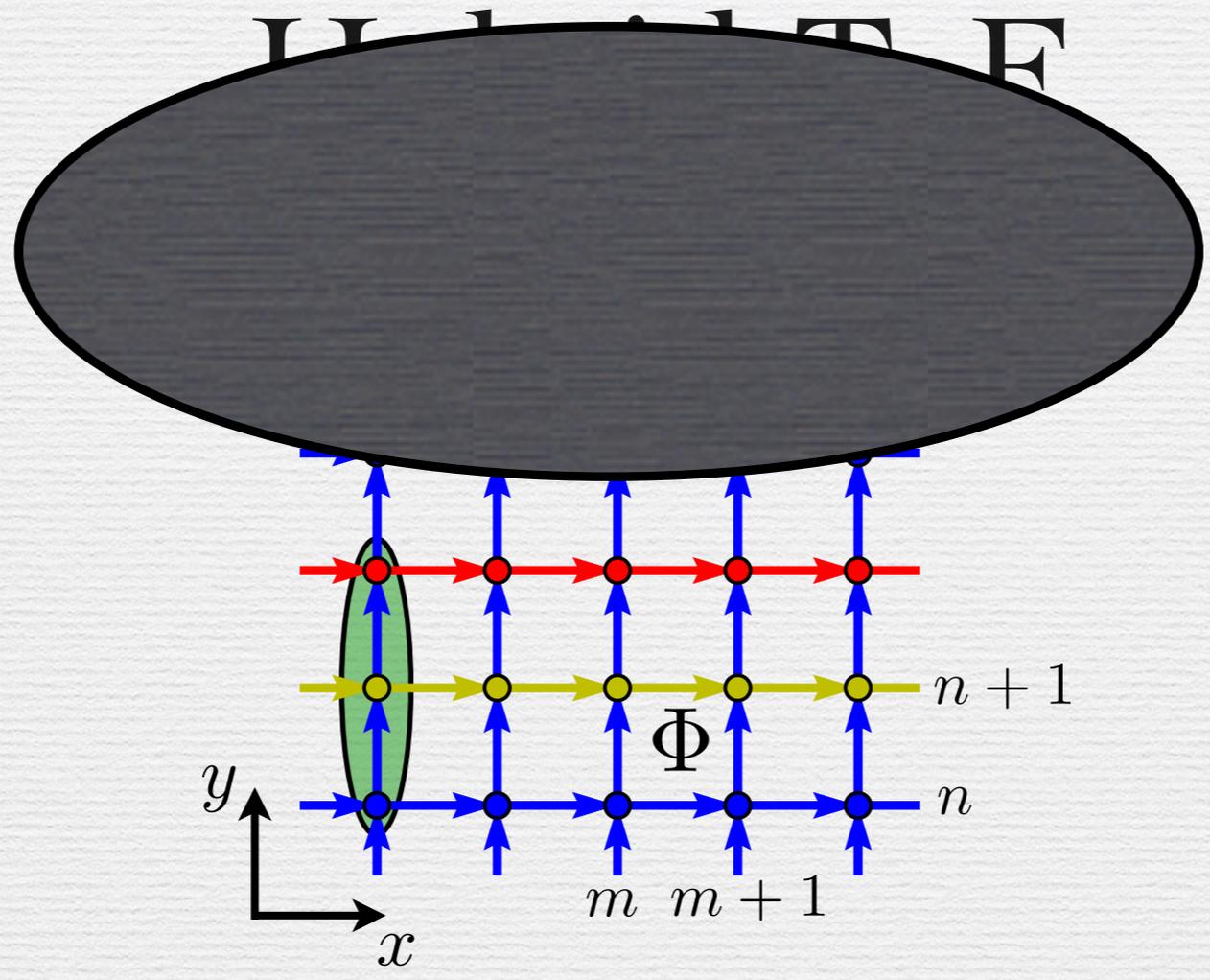
Hamiltonian

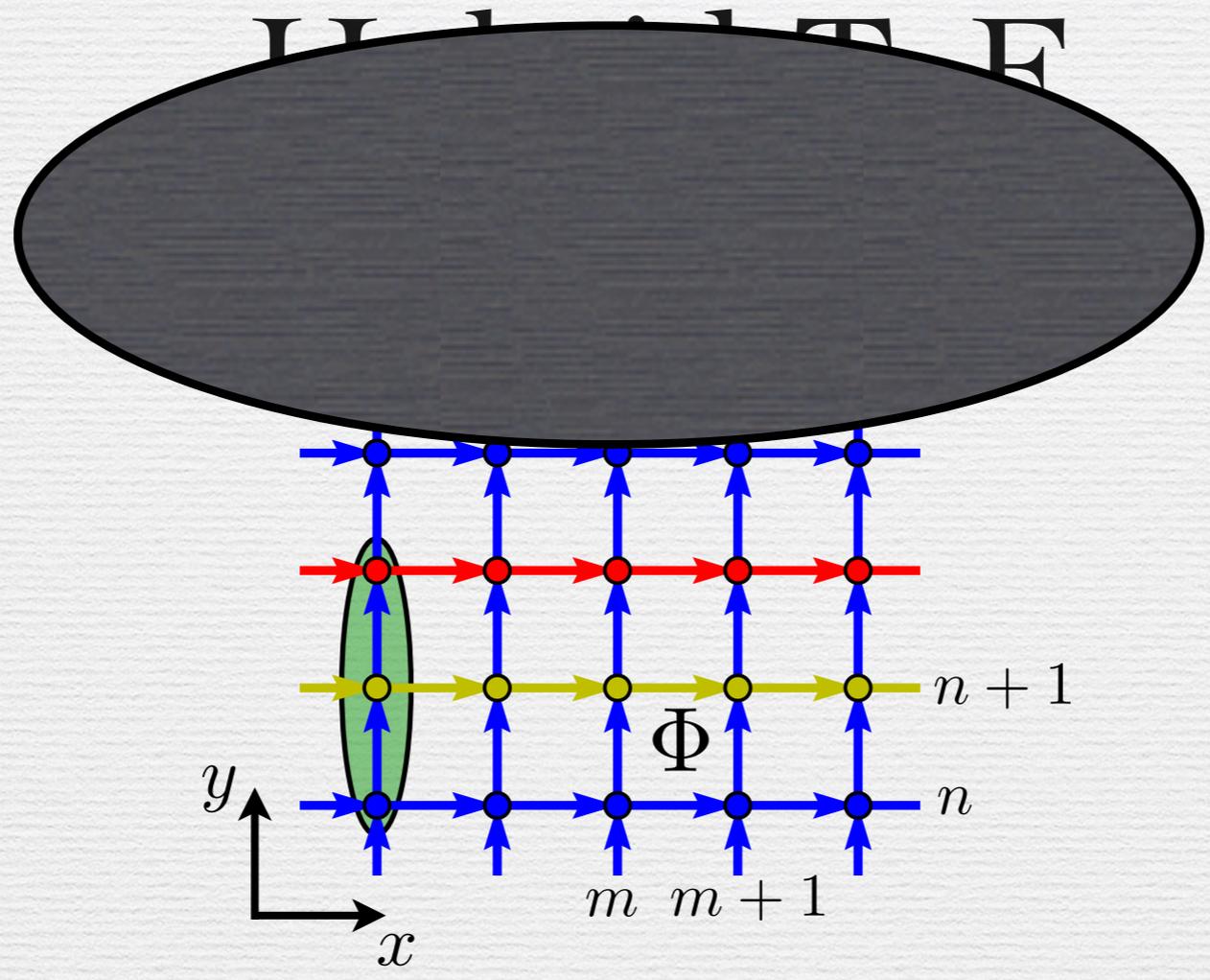


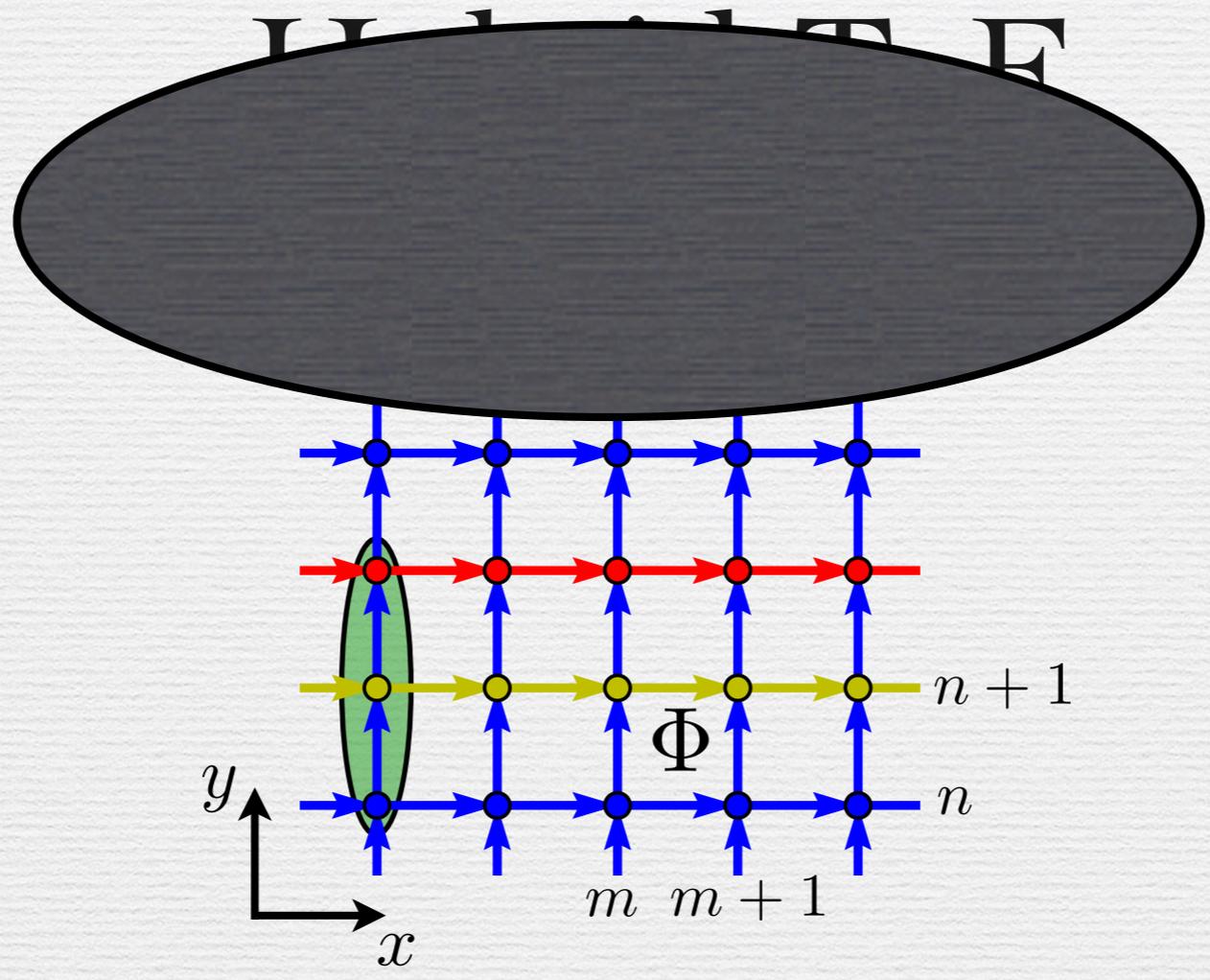


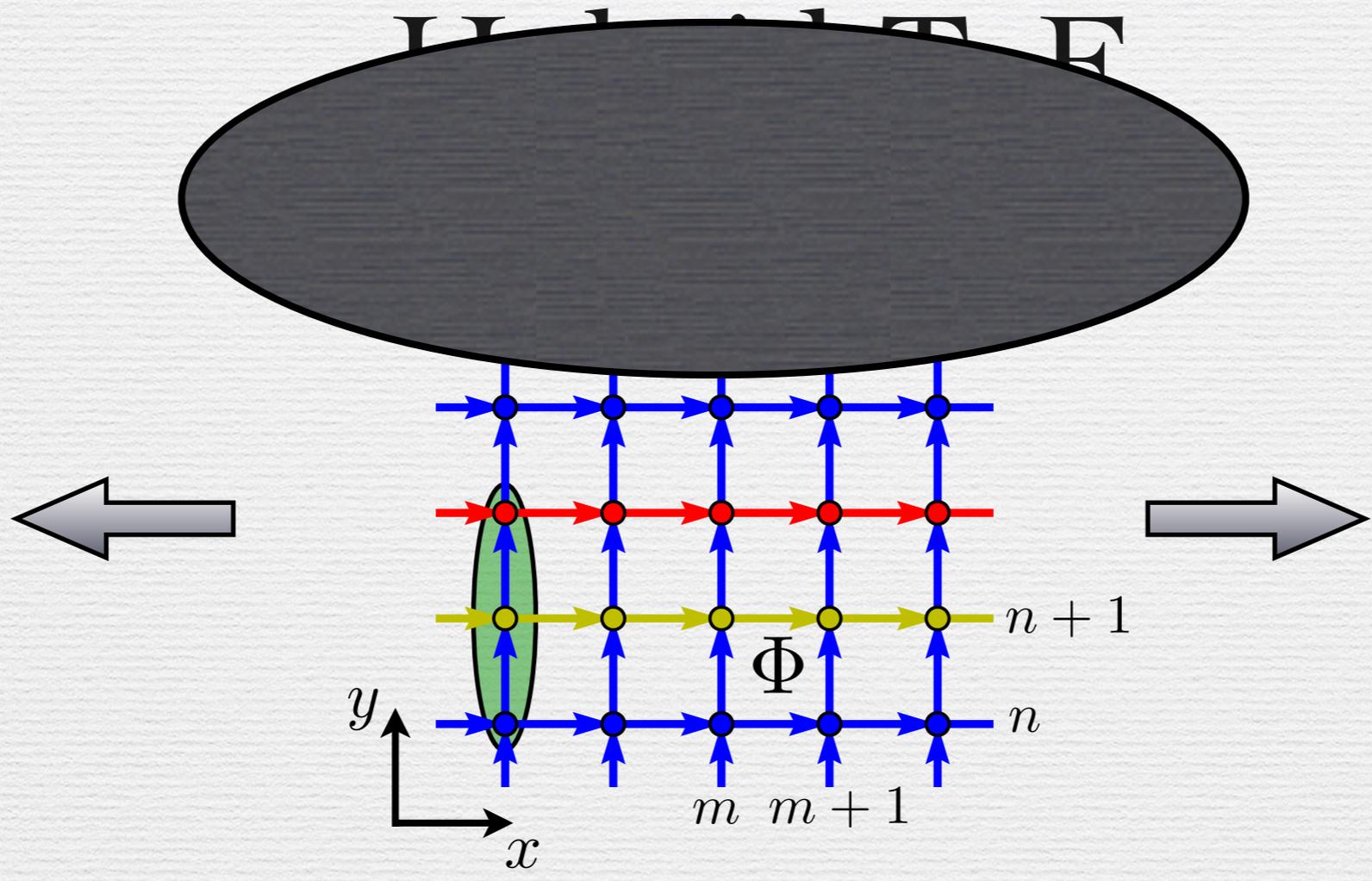


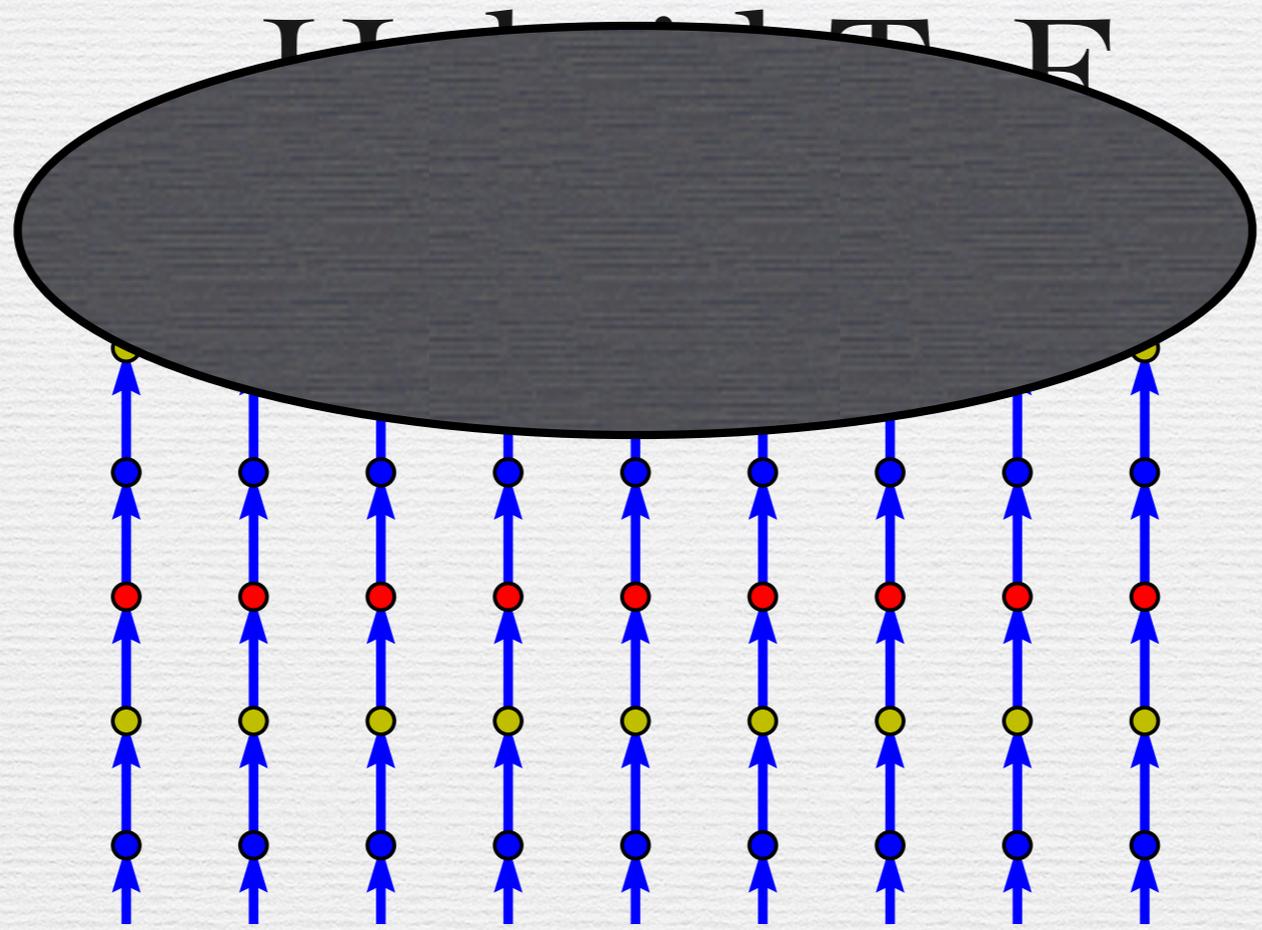




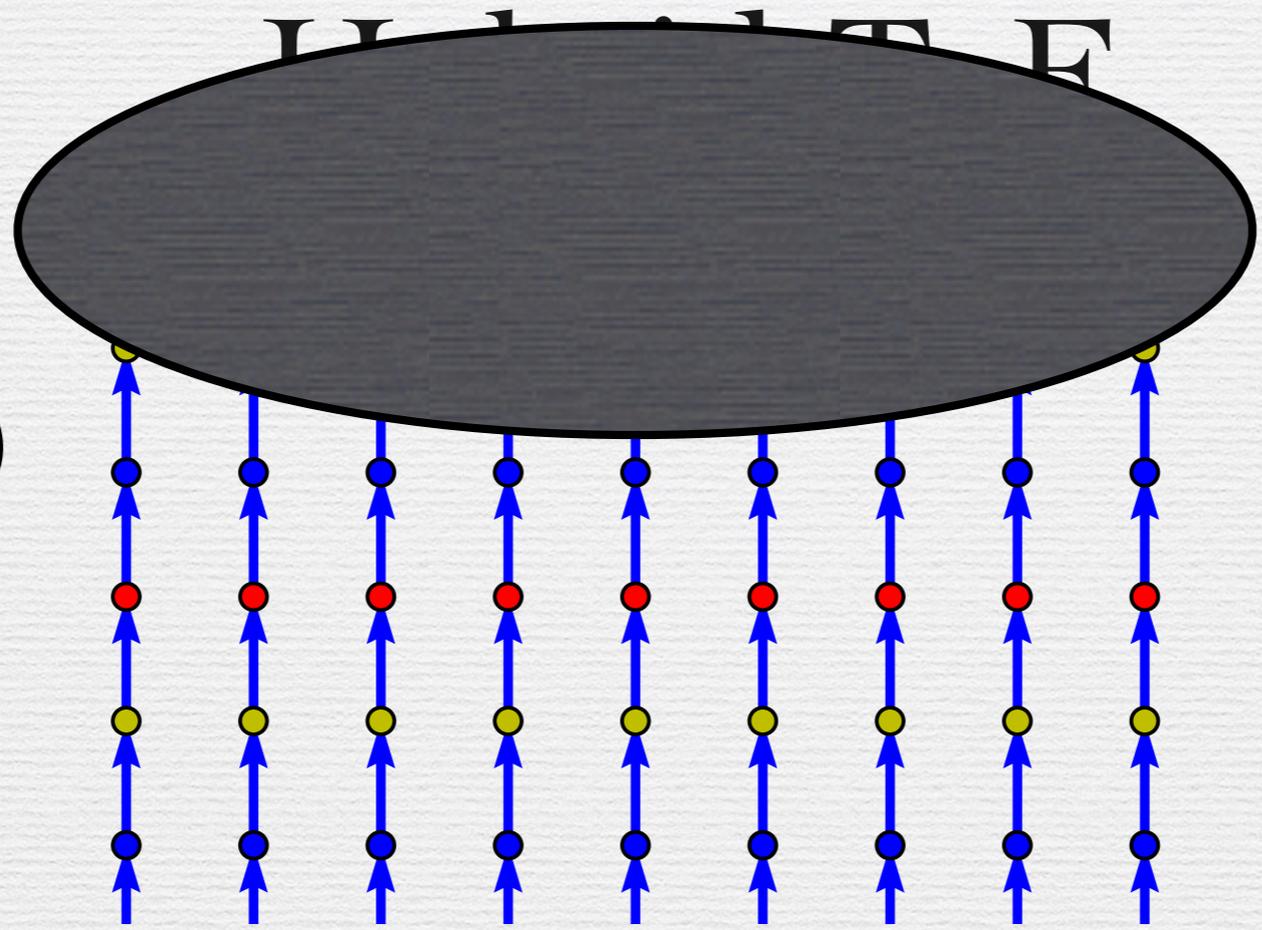
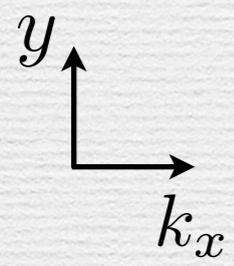




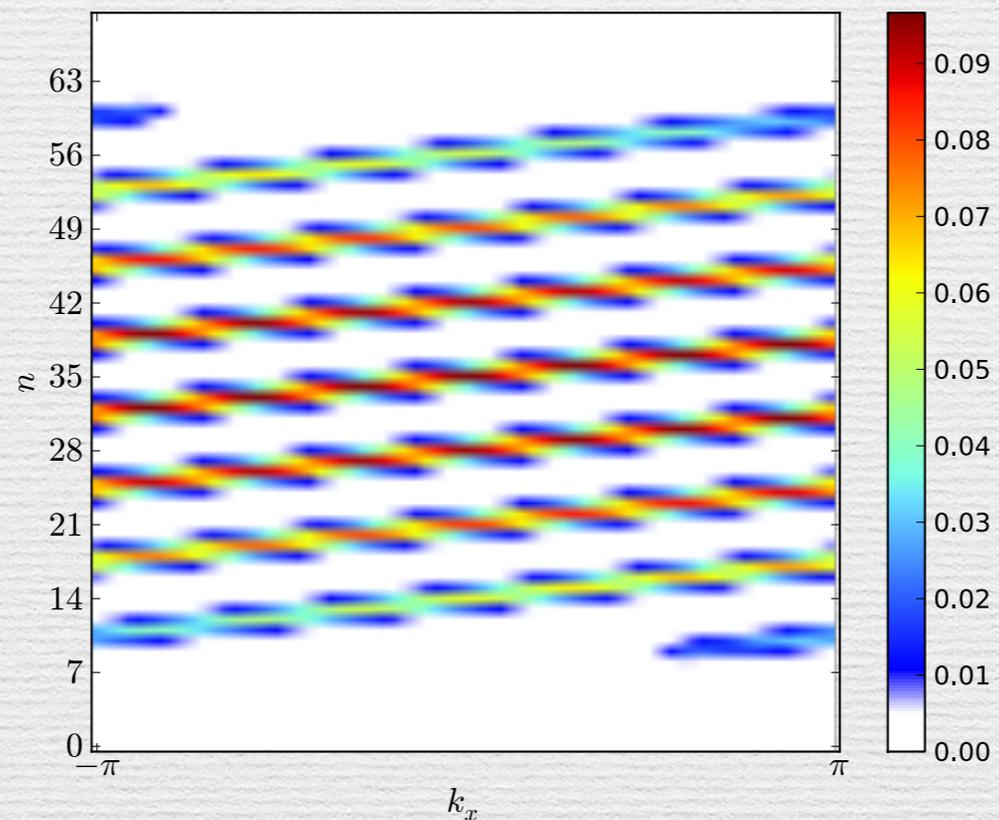
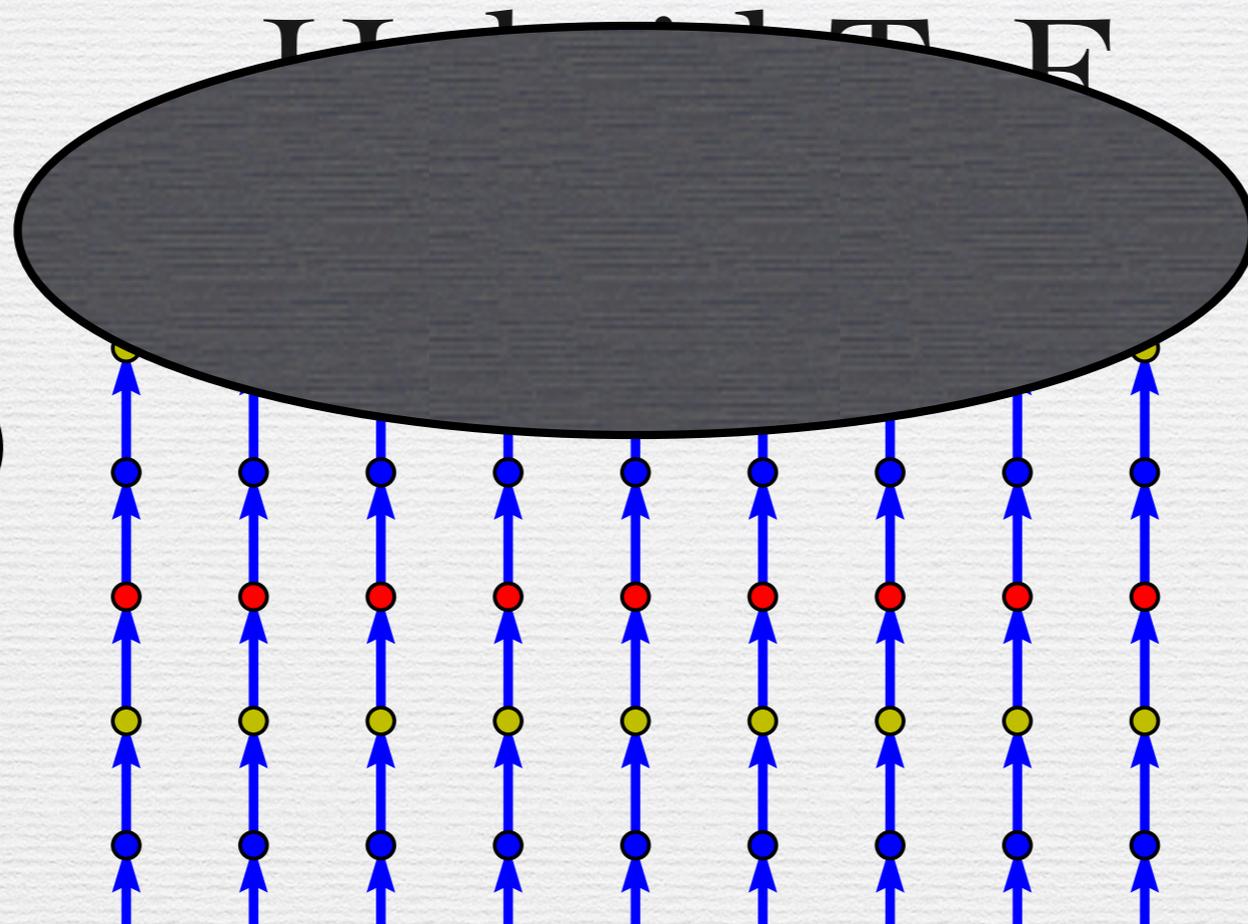
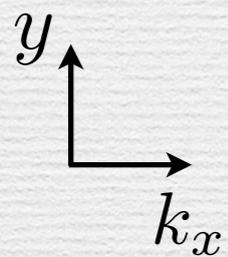




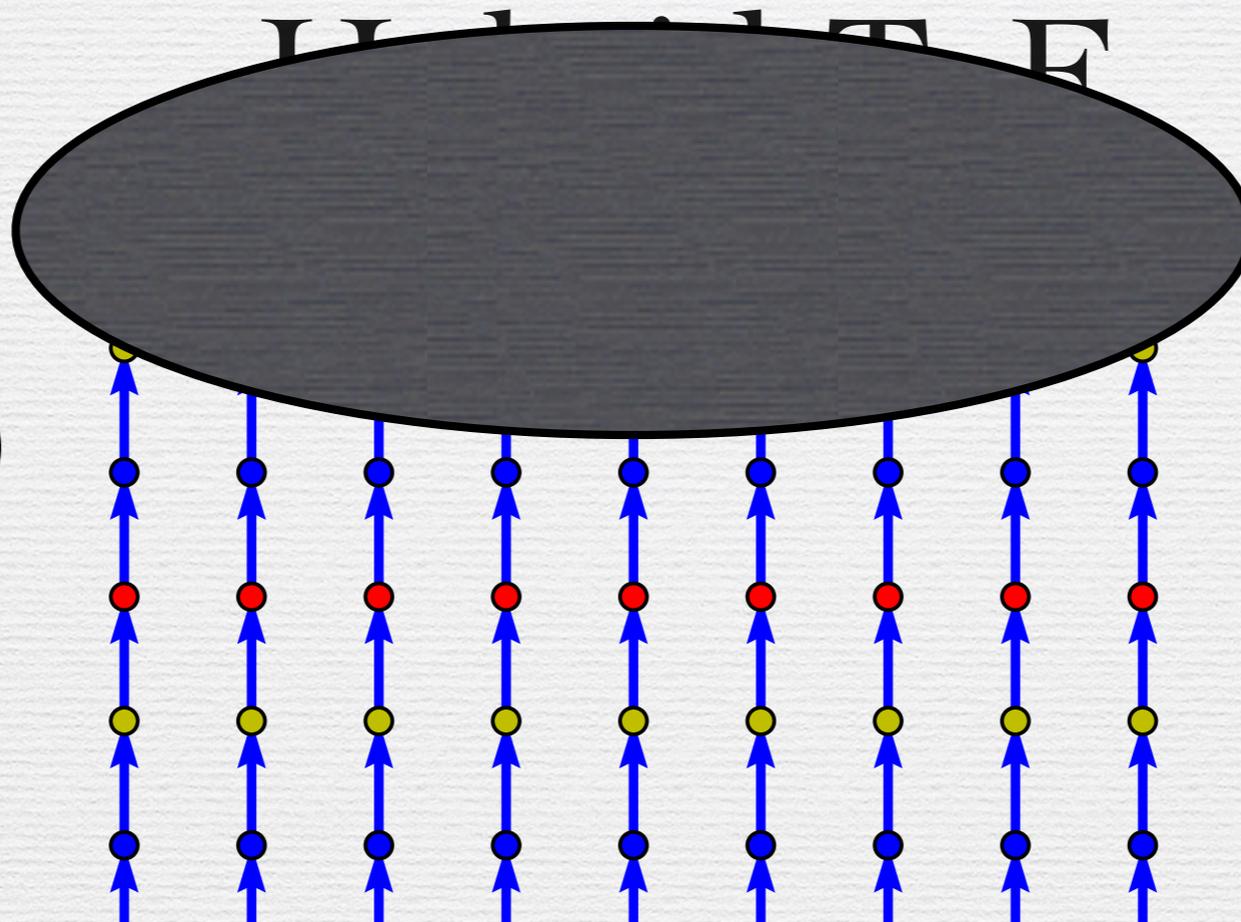
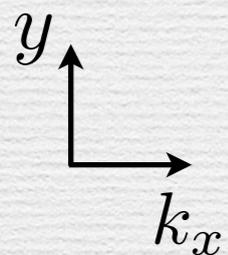
$$\rho(k_x, y)$$



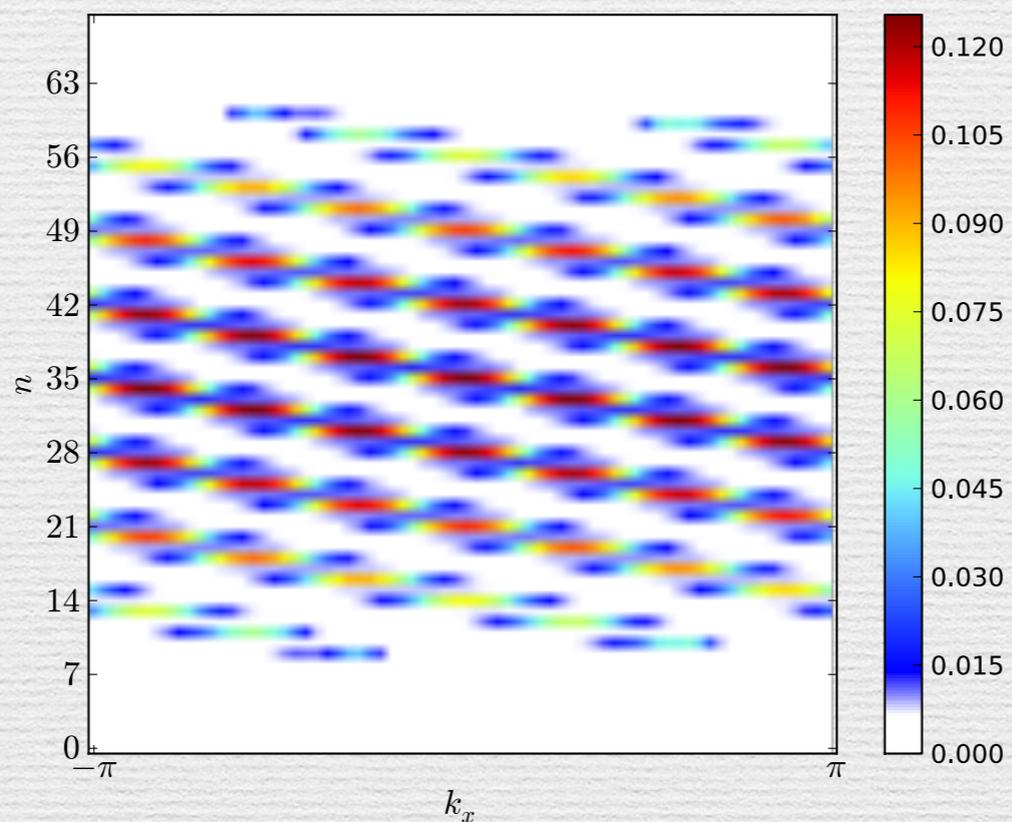
$$\rho(k_x, y)$$



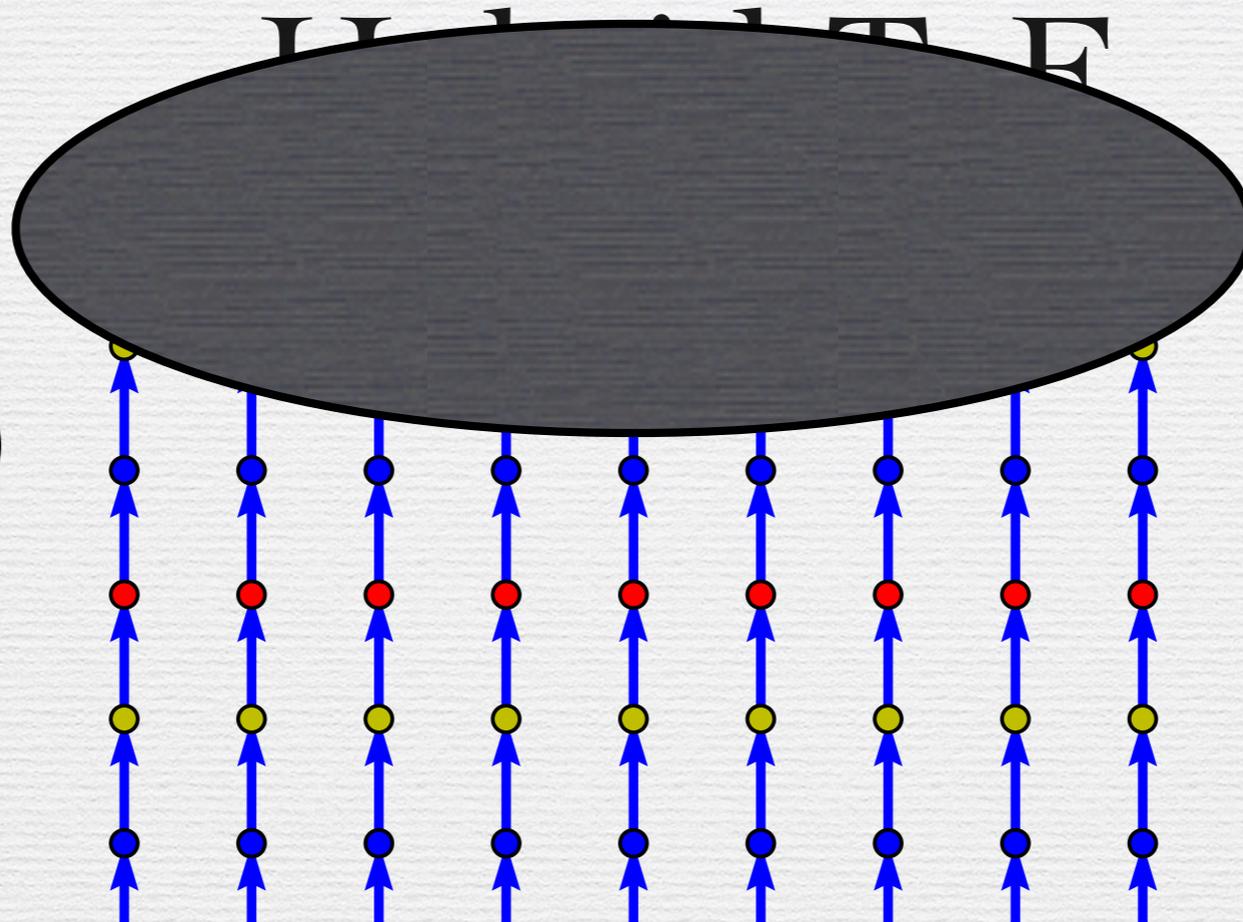
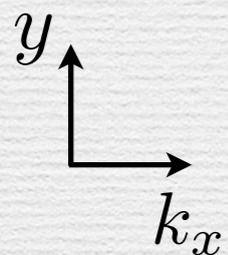
$\rho(k_x, y)$



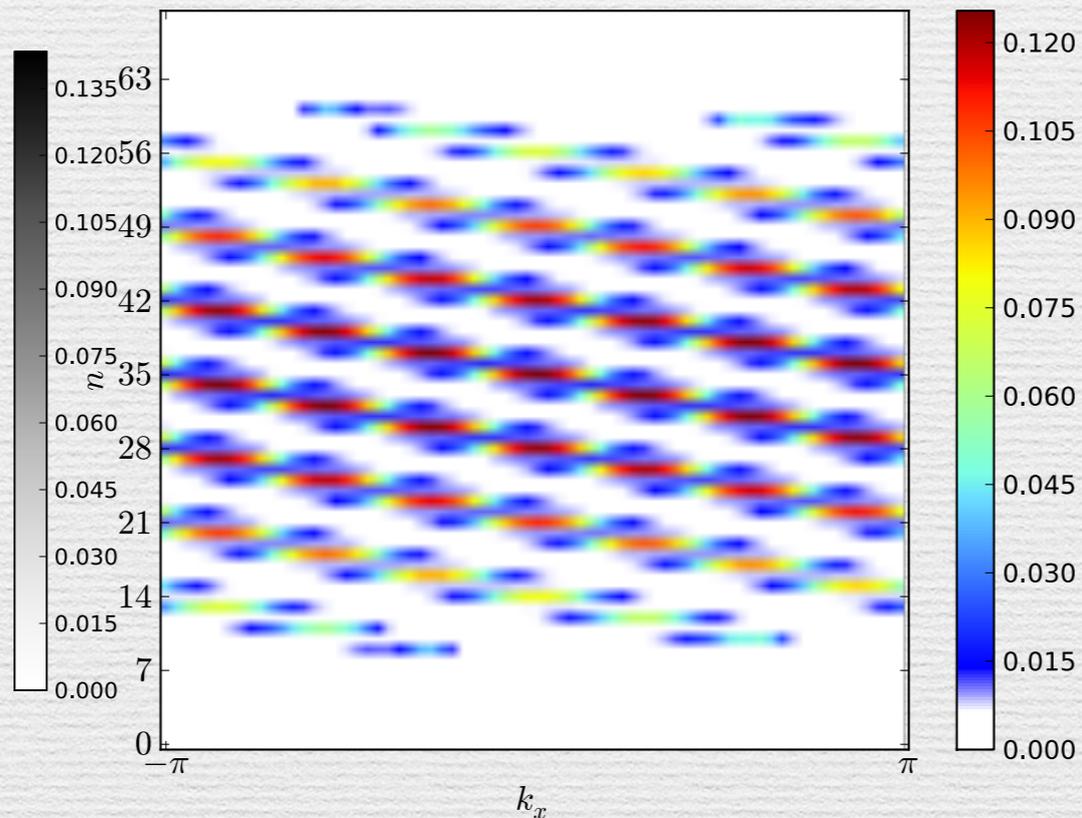
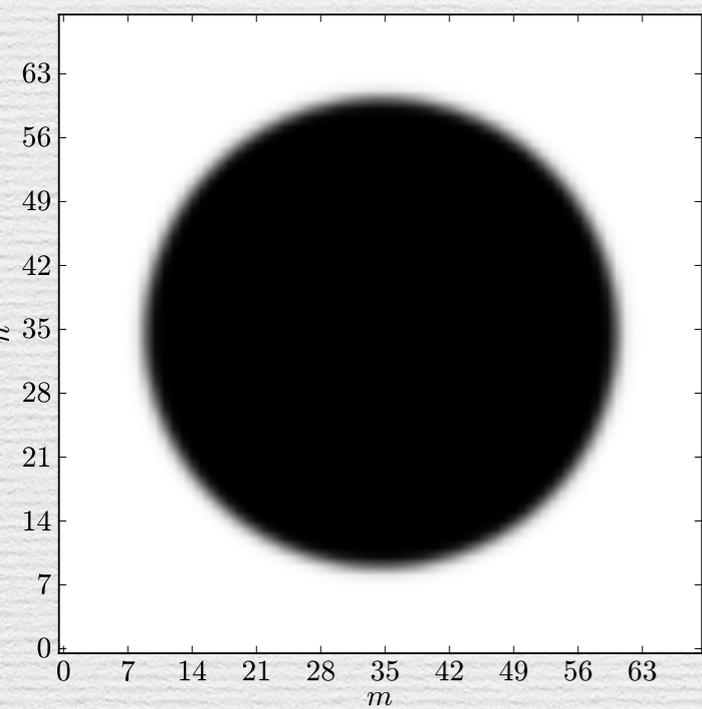
$$\Phi = 3/7 \quad C = -2$$



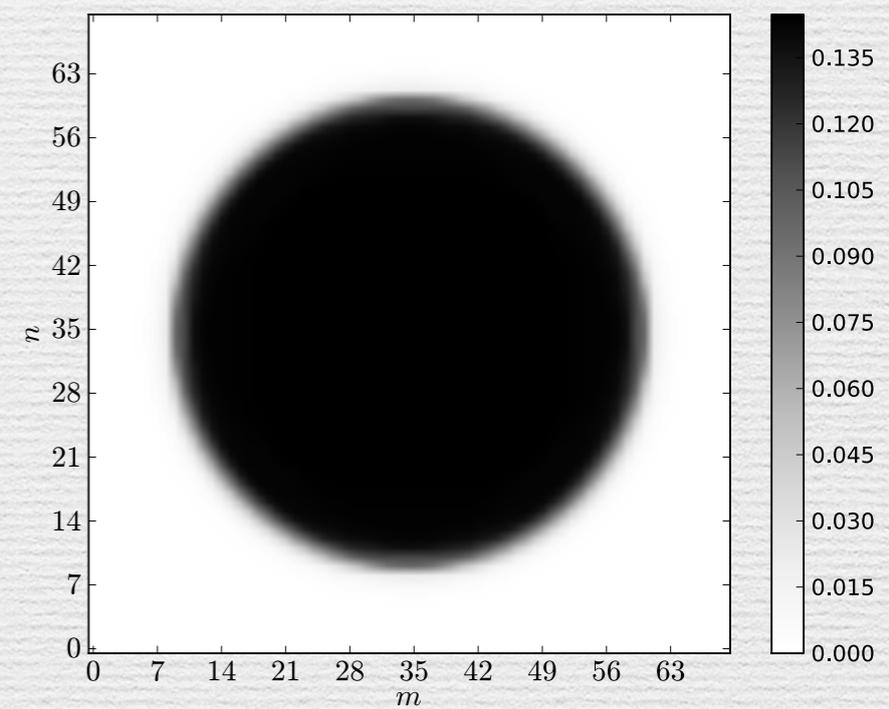
$\rho(k_x, y)$



$\Phi = 1/7 \quad C = 1$

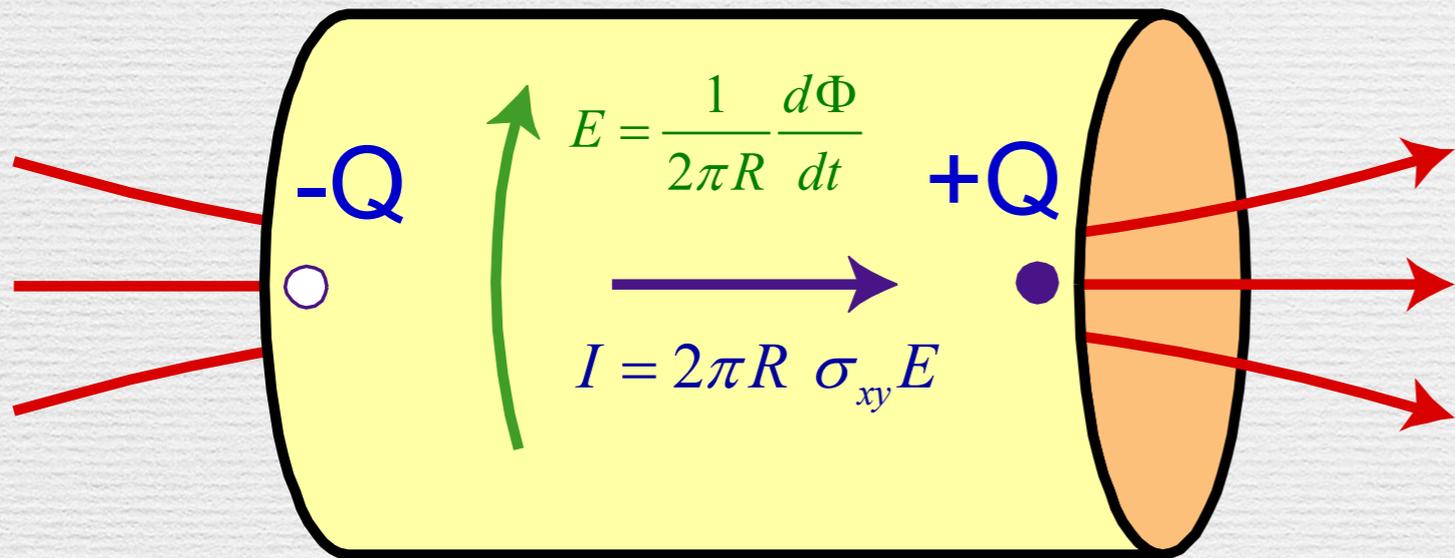


$\Phi = 3/7 \quad C = -2$



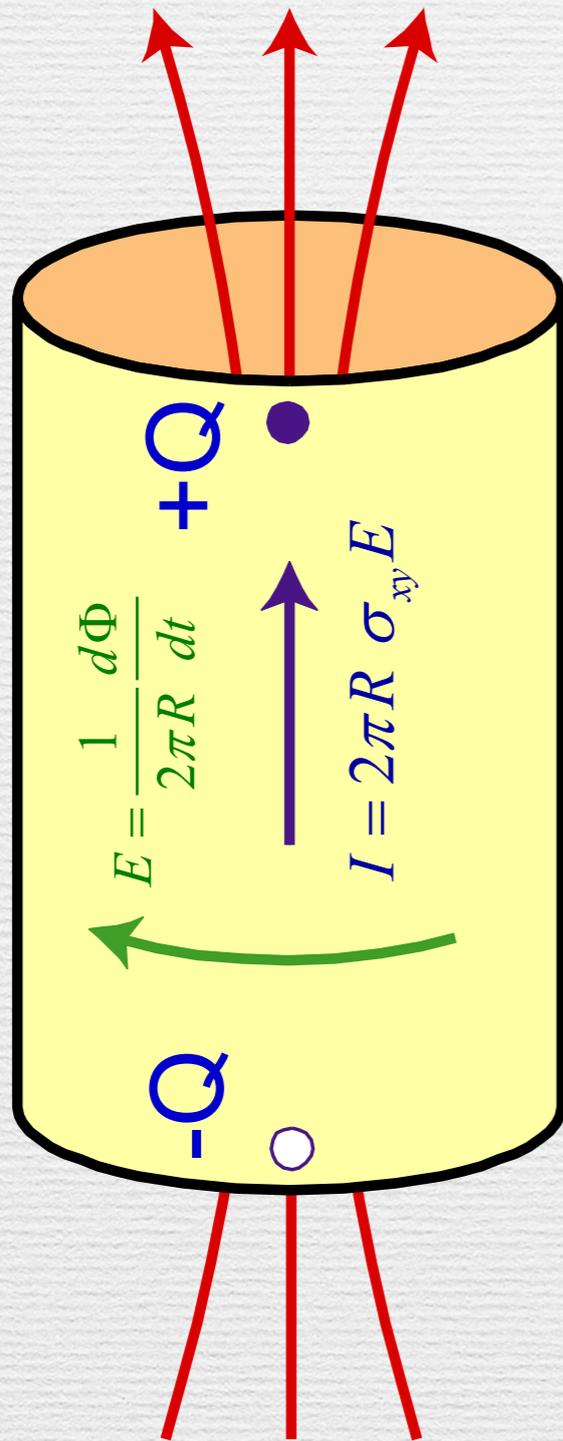
Why it works?

Topological charge pumping



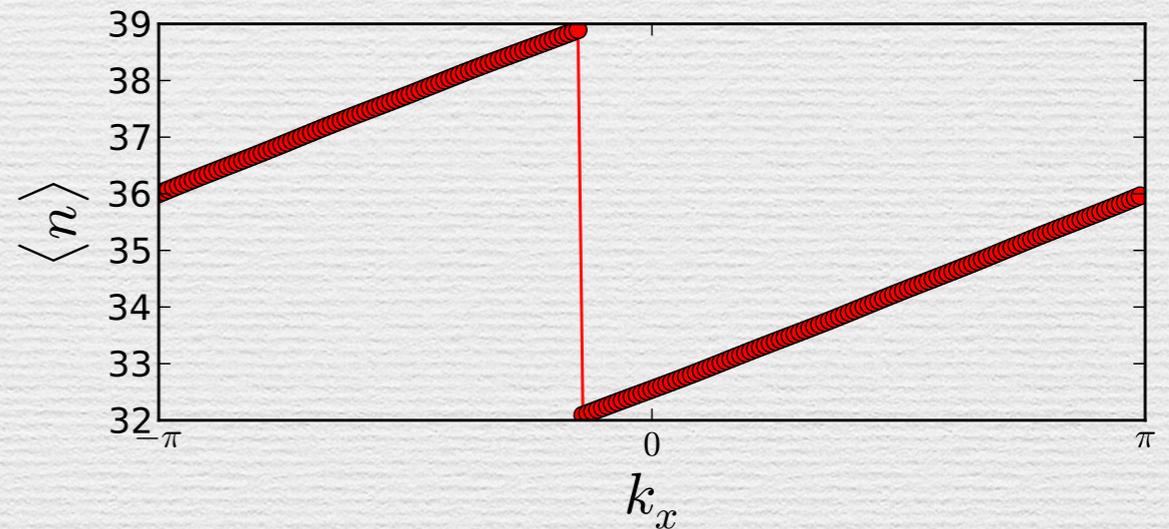
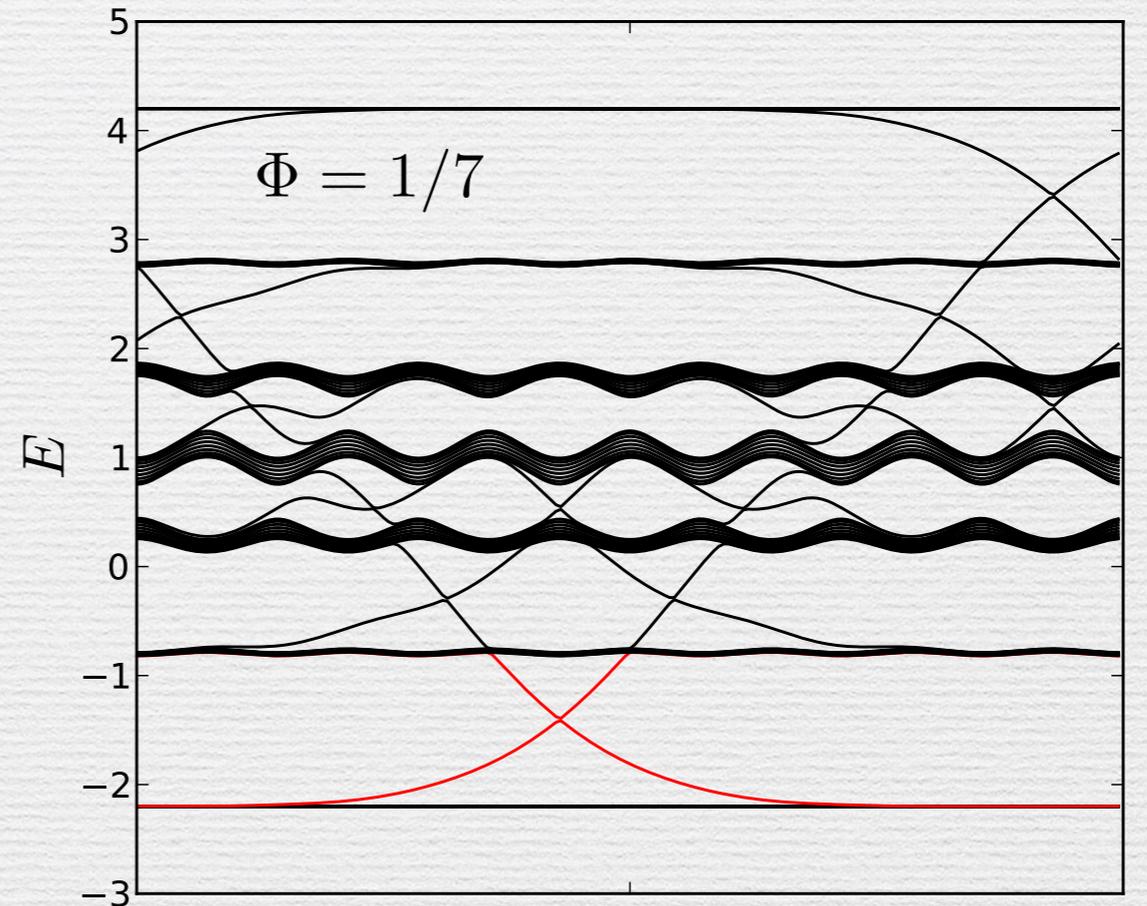
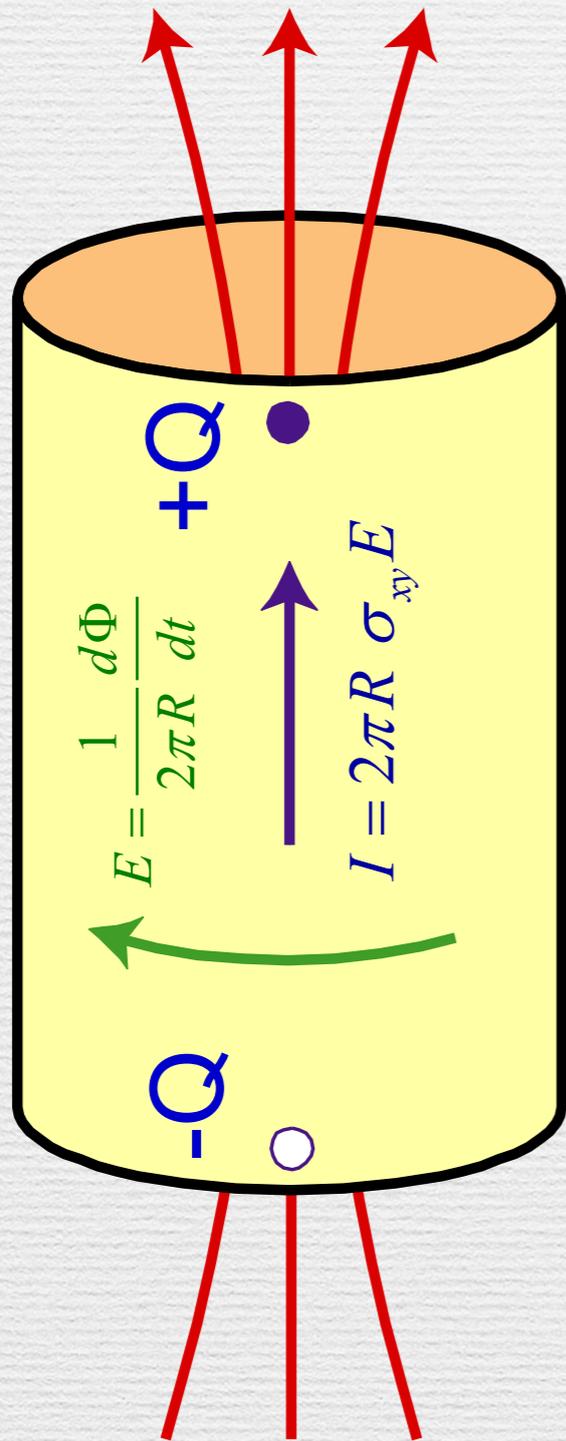
Why it works?

Topological charge pumping



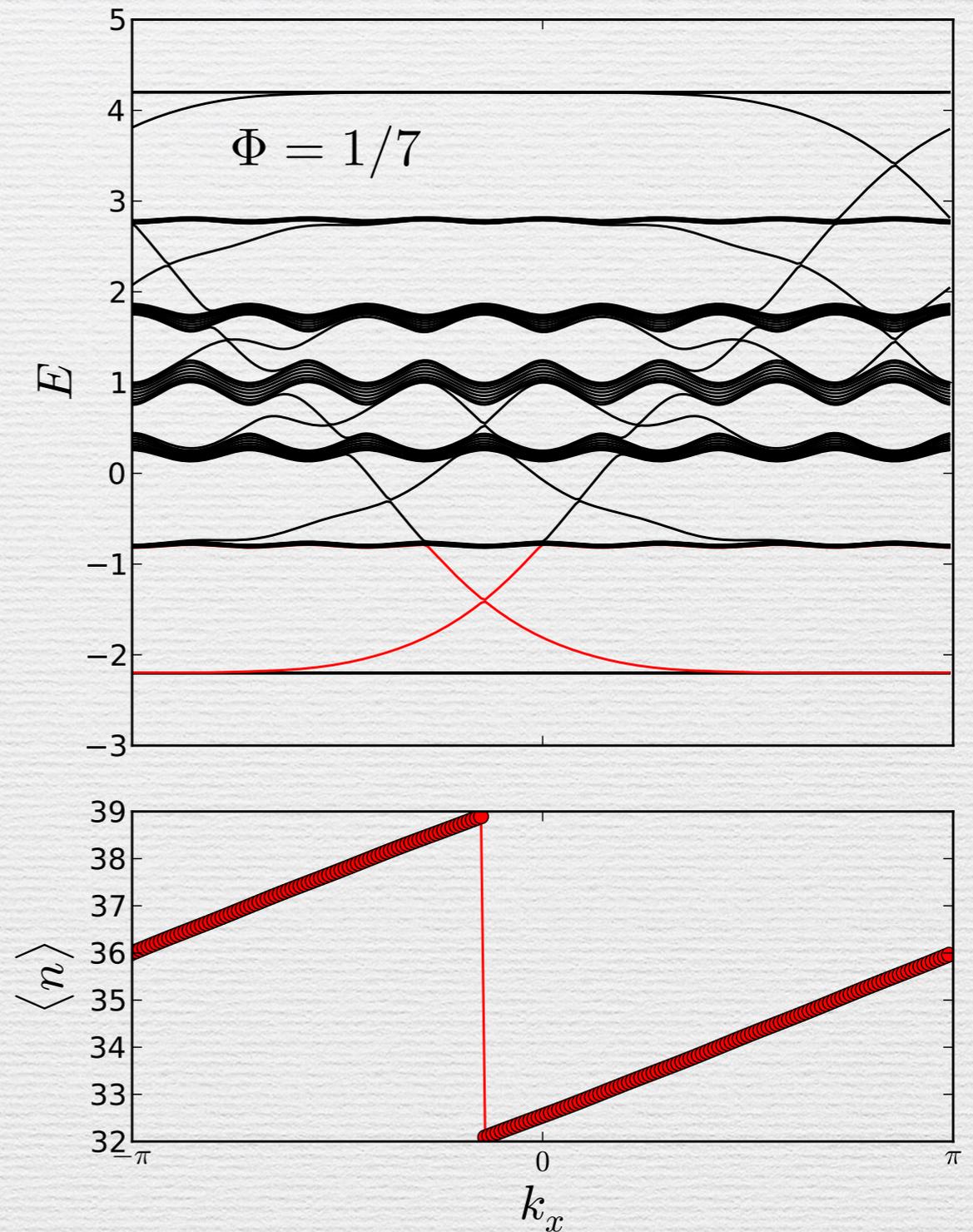
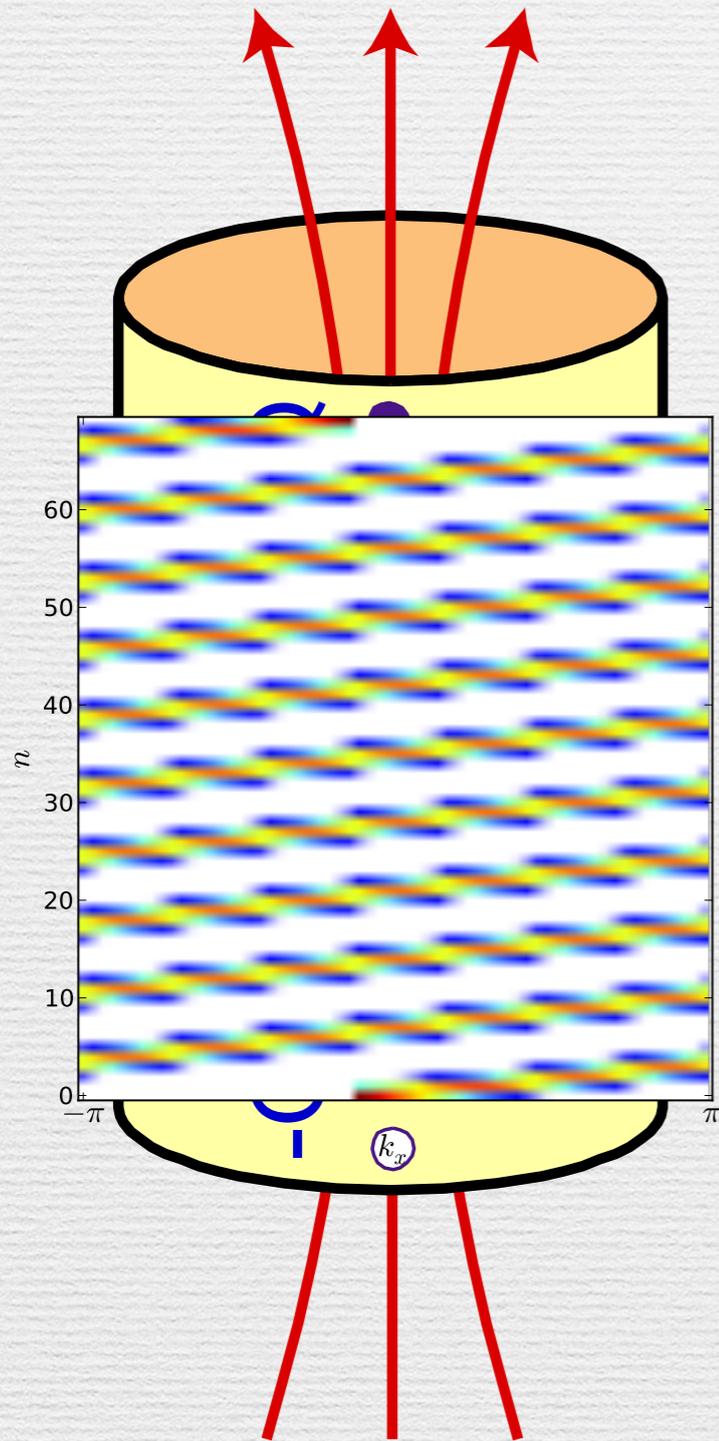
Why it works?

Topological charge pumping

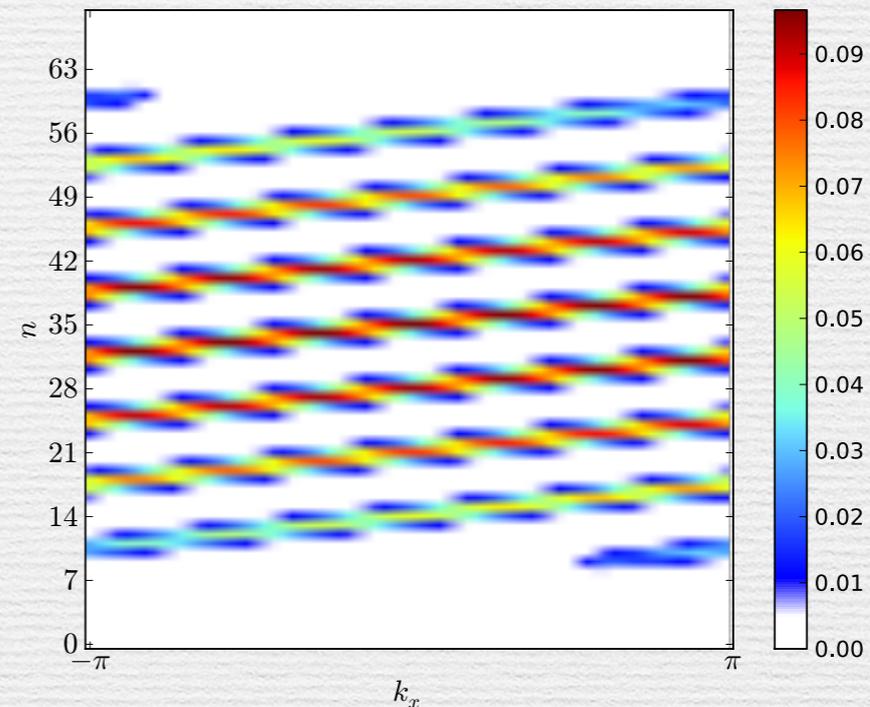
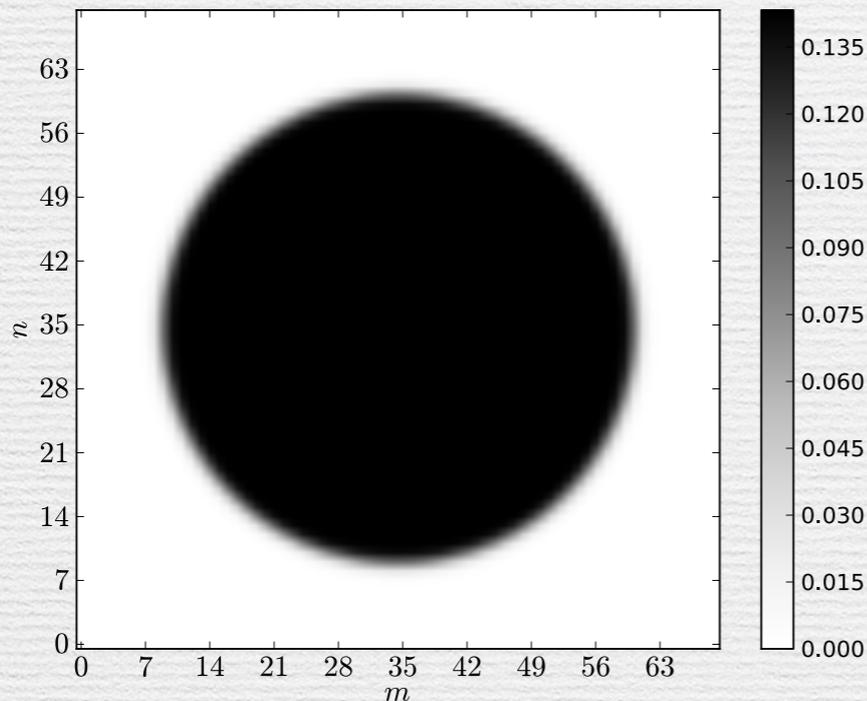


Why it works?

Topological charge pumping



Summary



Topological charge pumping is a common thread unifies many features of topological states

Guideline for design and detection of topological phases in cold atom systems

Thank you!

FAQ

Tight binding limit? Do not need

Edge state modes, fractionalized charge ? Do not need

Is sliding topological ? Yes